



Electronic structure theory and magnetism

Olle Eriksson
Uppsala University
and

WISE-Wallenberg Initiative Materials Science



A sustainable future through materials science

The Wallenberg Initiative Materials Science for Sustainability (WISE), is a Swedish research program initiated and financed by the Knut and Alice Wallenberg Foundation.

M€ 270 spanning 2022–2033

186 PhD students and 186 postdocs

Graduate school

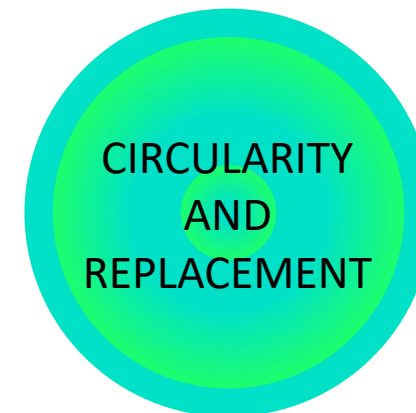
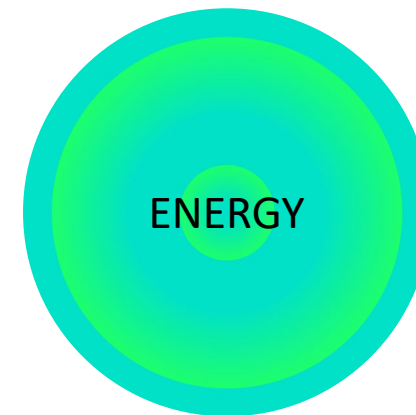
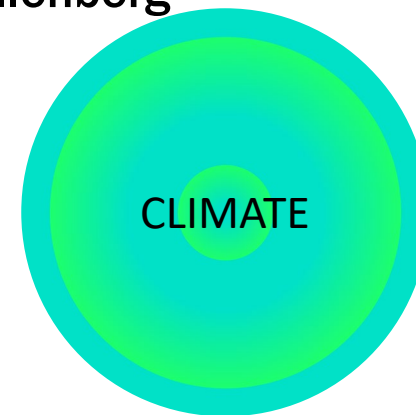
Research and technology platforms

Research arenas with Swedish industry

Strategic recruitment of 28 asst/assoc professors

International guest professor program

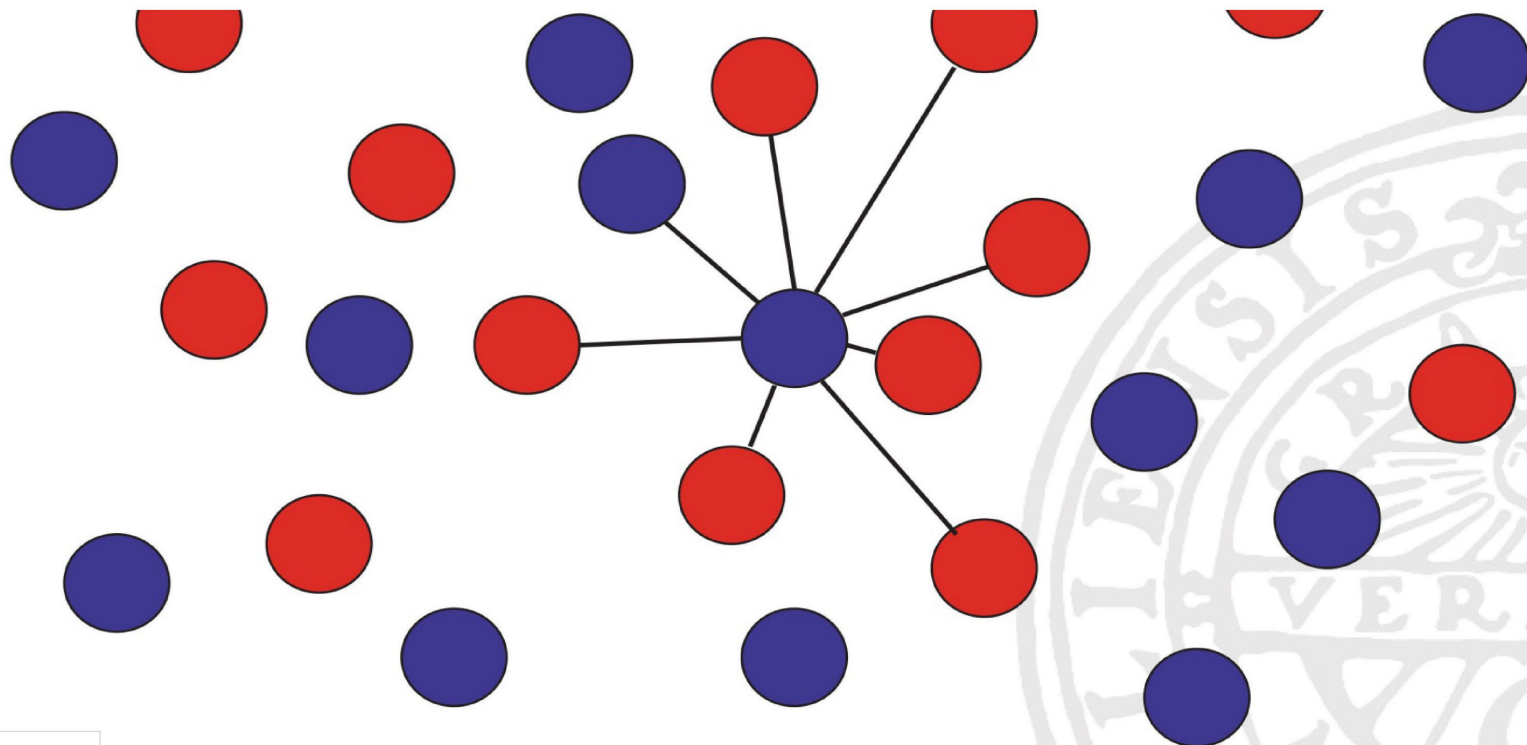
wise-materials.org





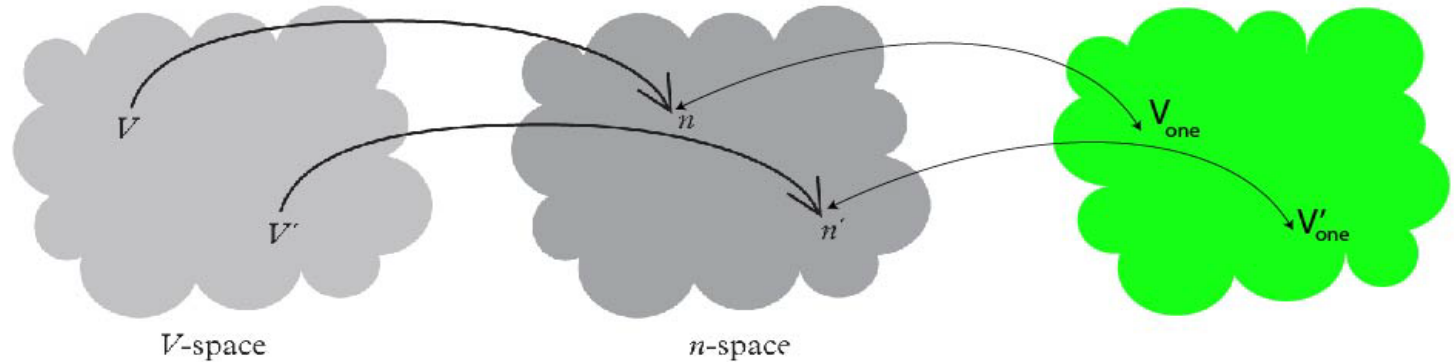
Full electron Hamiltonian

$$H = -\frac{1}{2} \sum_i \nabla_i^2 + \frac{1}{2} \sum_{i \neq j} \frac{e^2}{|\mathbf{r}_i - \mathbf{r}_j|} - \sum_{i,I} \frac{Z_I e^2}{|\mathbf{r}_i - \mathbf{R}_I|}$$





Walter Kohns ideas



Hamiltonian of real systems

Hamiltonian of one-electron system

- i) A functional exists that couples total energy to ground state electron density
- ii) Find best approximation of $E[n(r)]$ and minimize w.r.t. $n(r) \rightarrow$ Dirac eller Schrödinger eqn. for one-electron system



Density Functional Theory

$$\hat{H}\Psi_{\text{gs}} = E_{\text{gs}} \Psi_{\text{gs}}$$

$$n_{\text{gs}}(\mathbf{r}) = \sum_{i=1}^N \int d^3r_i |\Psi_{\text{gs}}(\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_N)|^2 \delta(\mathbf{r} - \mathbf{r}_i) \text{ electron density}$$

$$\hat{H}'\Psi'_{\text{gs}} = E'_{\text{gs}} \Psi'_{\text{gs}}$$

$$E_{\text{gs}} = \langle \Psi_{\text{gs}} | \hat{H} | \Psi_{\text{gs}} \rangle < \langle \Psi'_{\text{gs}} | \hat{H} | \Psi'_{\text{gs}} \rangle \quad \text{variational principle}$$

$$\langle \Psi'_{\text{gs}} | \hat{H} | \Psi'_{\text{gs}} \rangle = \langle \Psi'_{\text{gs}} | \hat{H} + V'_{\text{ext}} - V'_{\text{ext}} | \Psi'_{\text{gs}} \rangle = \langle \Psi'_{\text{gs}} | \hat{H}' + V_{\text{ext}} - V'_{\text{ext}} | \Psi'_{\text{gs}} \rangle$$



Density Functional Theory

$$\langle \Psi'_{\text{gs}} | \hat{H} | \Psi'_{\text{gs}} \rangle = E'_{\text{gs}} + \int n'_{\text{gs}} (V_{\text{ext}} - V'_{\text{ext}}) d^3 r$$

$$E_{\text{gs}} < E'_{\text{gs}} + \int n'_{\text{gs}} (V_{\text{ext}} - V'_{\text{ext}}) d^3 r$$

$$E'_{\text{gs}} < E_{\text{gs}} + \int n_{\text{gs}} (V'_{\text{ext}} - V_{\text{ext}}) d^3 r$$

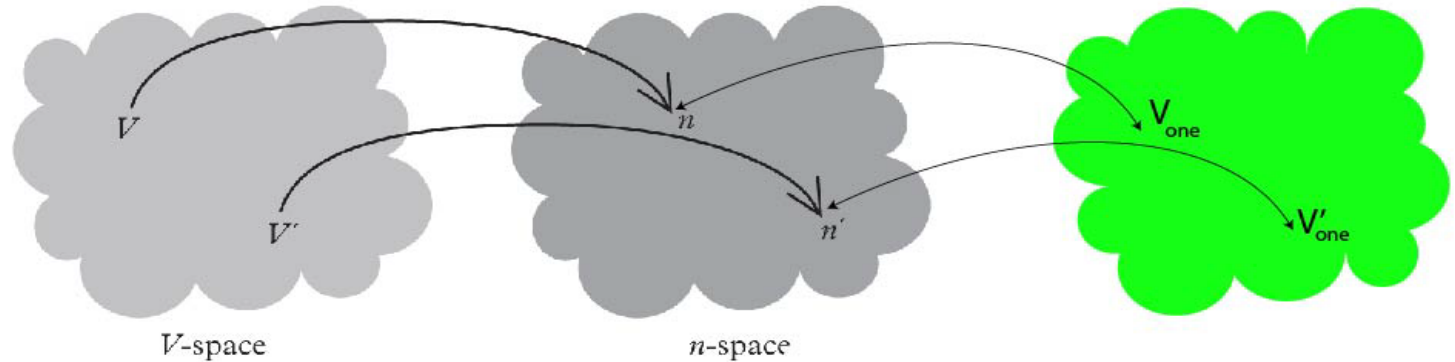
Assume densities are the same

$$E_{\text{gs}} + E'_{\text{gs}} < E'_{\text{gs}} + E_{\text{gs}} + \int n_{\text{gs}} (V_{\text{ext}} - V'_{\text{ext}}) d^3 r + \int n_{\text{gs}} (V'_{\text{ext}} - V_{\text{ext}}) d^3 r,$$

$$E_{\text{gs}} + E'_{\text{gs}} < E'_{\text{gs}} + E_{\text{gs}}$$



Walter Kohns ideas



Hamiltonian of real systems

Hamiltonian of one-electron system

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One-electron Hamiltonian

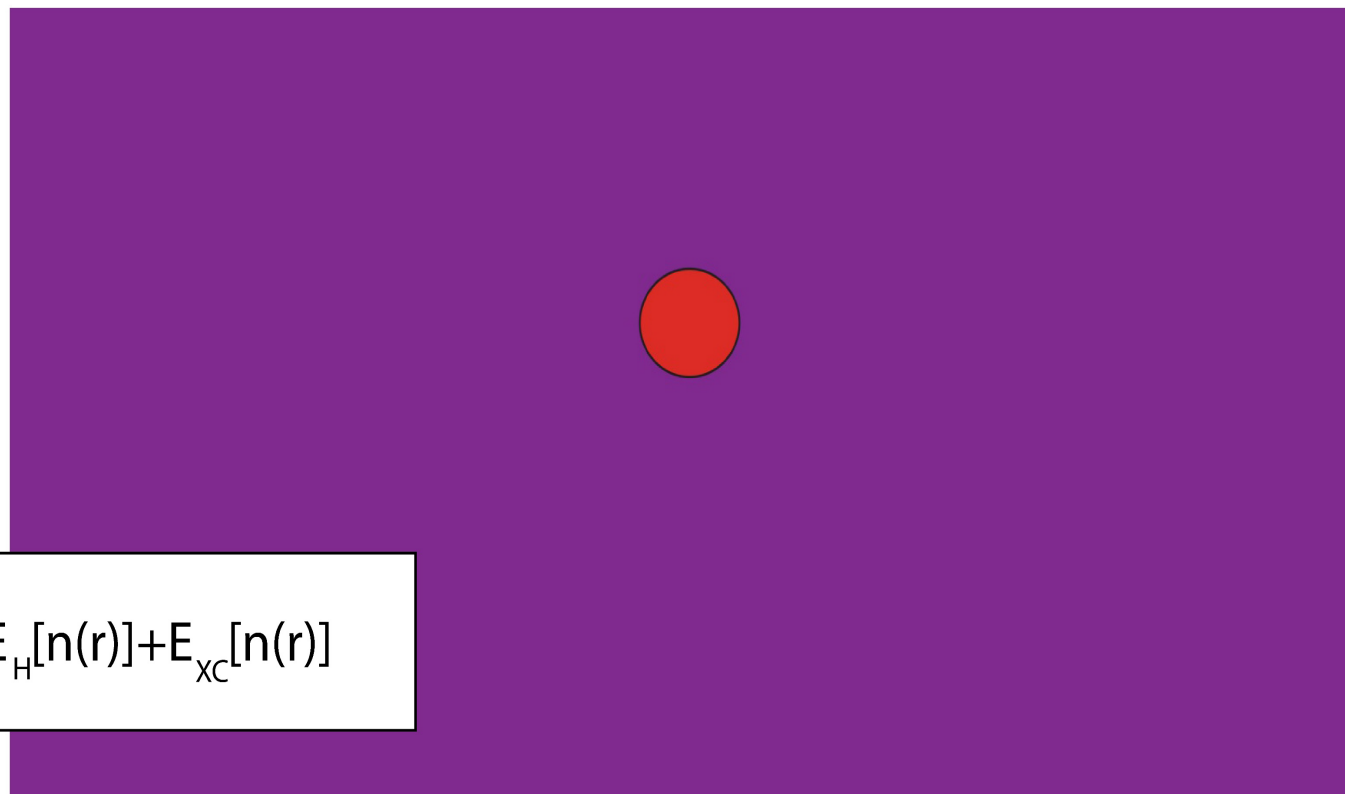
$$\left(-\frac{\nabla^2}{2} + V_{\text{eff}}\right) \psi = \varepsilon \psi$$

$$V_{\text{eff}}(n(r))$$

$$n(r) = \sum |\psi(r)|^2$$

$$dE[n(r)]/dn=0$$

$$E[n(r)] = E_{\text{kin}}[n(r)] + E_{\text{H}}[n(r)] + E_{\text{XC}}[n(r)]$$





The density functional

$$E[n_{\text{op}}(\mathbf{r})] = T_{\text{op}}[n_{\text{op}}(\mathbf{r})] + \int n_{\text{op}}(\mathbf{r}) V_{\text{ext}}(\mathbf{r}) d^3 r \\ + \iint \frac{n_{\text{op}}(\mathbf{r}) \cdot n_{\text{op}}(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} d^3 r d^3 r' + E_{\text{xc}}[n_{\text{op}}(\mathbf{r})]$$

$$E_{\text{xc}}[n_{\text{op}}(\mathbf{r})] = \int \epsilon_{\text{xc}}[n_{\text{op}}(\mathbf{r})] n_{\text{op}}(\mathbf{r}) d^3 r$$

$$V_{\text{eff}}(\mathbf{r}) = V_{\text{ext}}(\mathbf{r}) + \int \frac{n_{\text{op}}(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} d^3 r' + \mu_{\text{xc}}[n_{\text{op}}(\mathbf{r})]$$

$$\mu_{\text{xc}}[n_{\text{op}}(\mathbf{r})] = \frac{\partial \{ \epsilon_{\text{xc}}[n_{\text{op}}(\mathbf{r})] n_{\text{op}}(\mathbf{r}) \}}{\partial n_{\text{op}}(\mathbf{r})} = \epsilon_{\text{xc}}[n_{\text{op}}(\mathbf{r})] + n_{\text{op}}(\mathbf{r}) \frac{\partial \{ \epsilon_{\text{xc}}[n_{\text{op}}(\mathbf{r})] \}}{\partial n_{\text{op}}(\mathbf{r})}$$



For a crystalline material

$$H_{one} = \frac{-\nabla^2}{2} + V_{eff}$$

$$H_{one}\psi(\mathbf{r}) = \epsilon\psi(\mathbf{r})$$

$$\psi(\mathbf{r}) = e^{i\mathbf{k}\cdot\mathbf{r}}u(\mathbf{r})$$

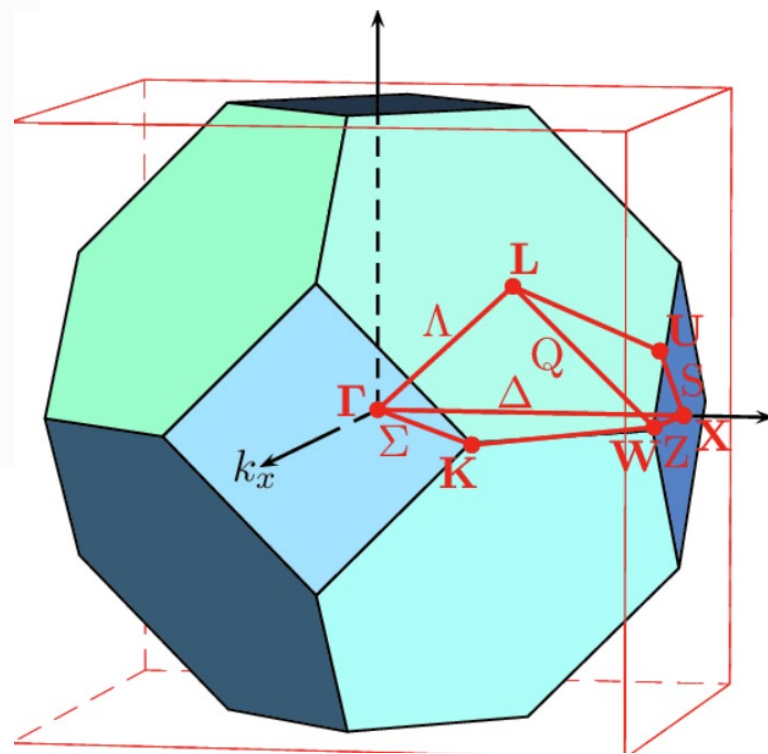
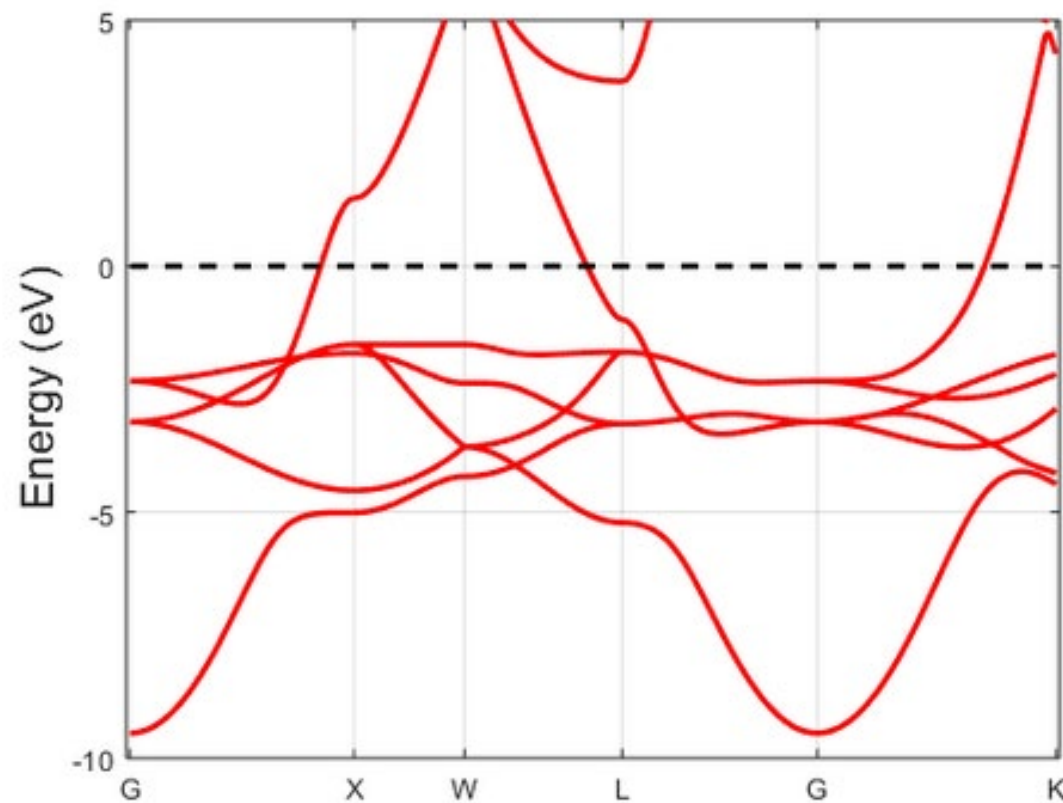
with $u(\mathbf{r})$ having the same periodicity as the lattice

$$\psi(\mathbf{r}) \rightarrow \psi_{\mathbf{k}}(\mathbf{r})$$

$$\epsilon \rightarrow \epsilon_{\mathbf{k}}$$

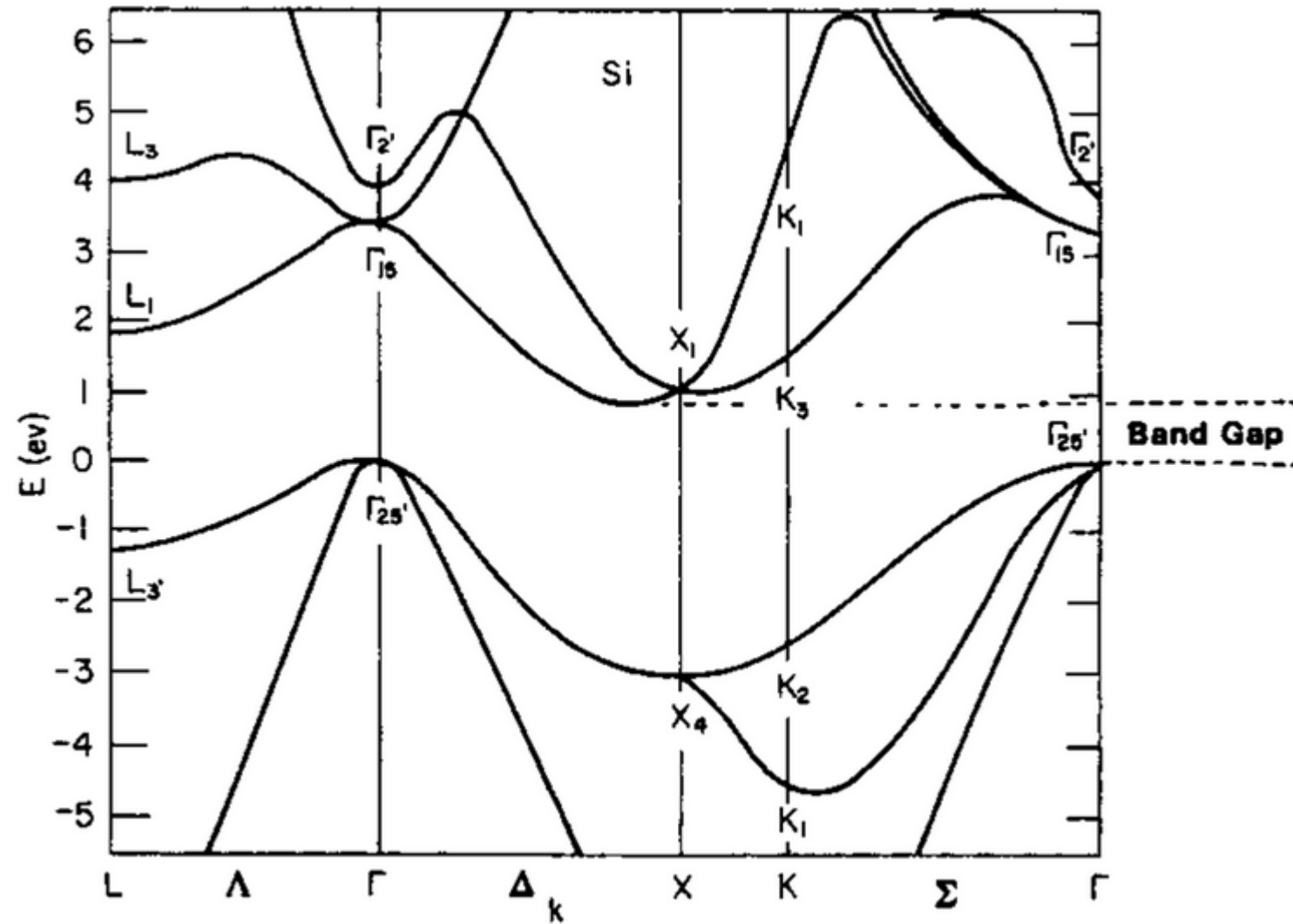


Bands of fcc Cu



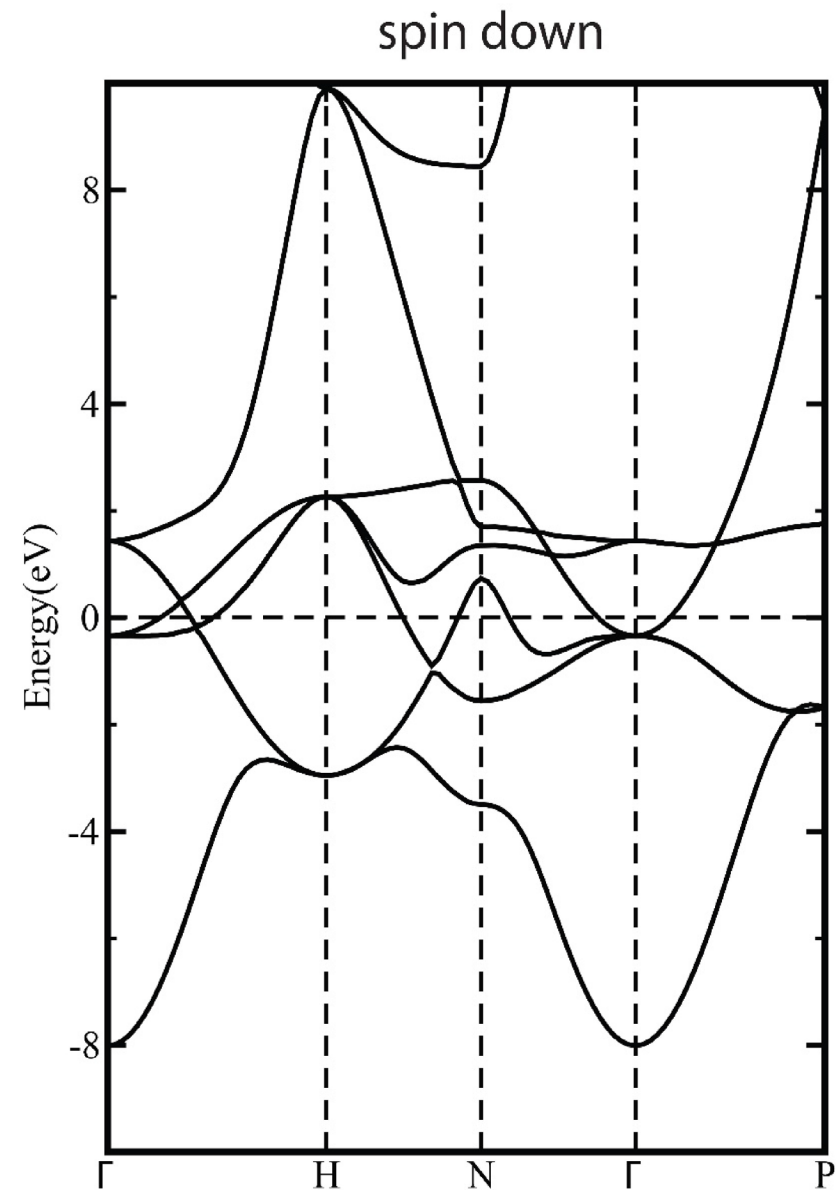
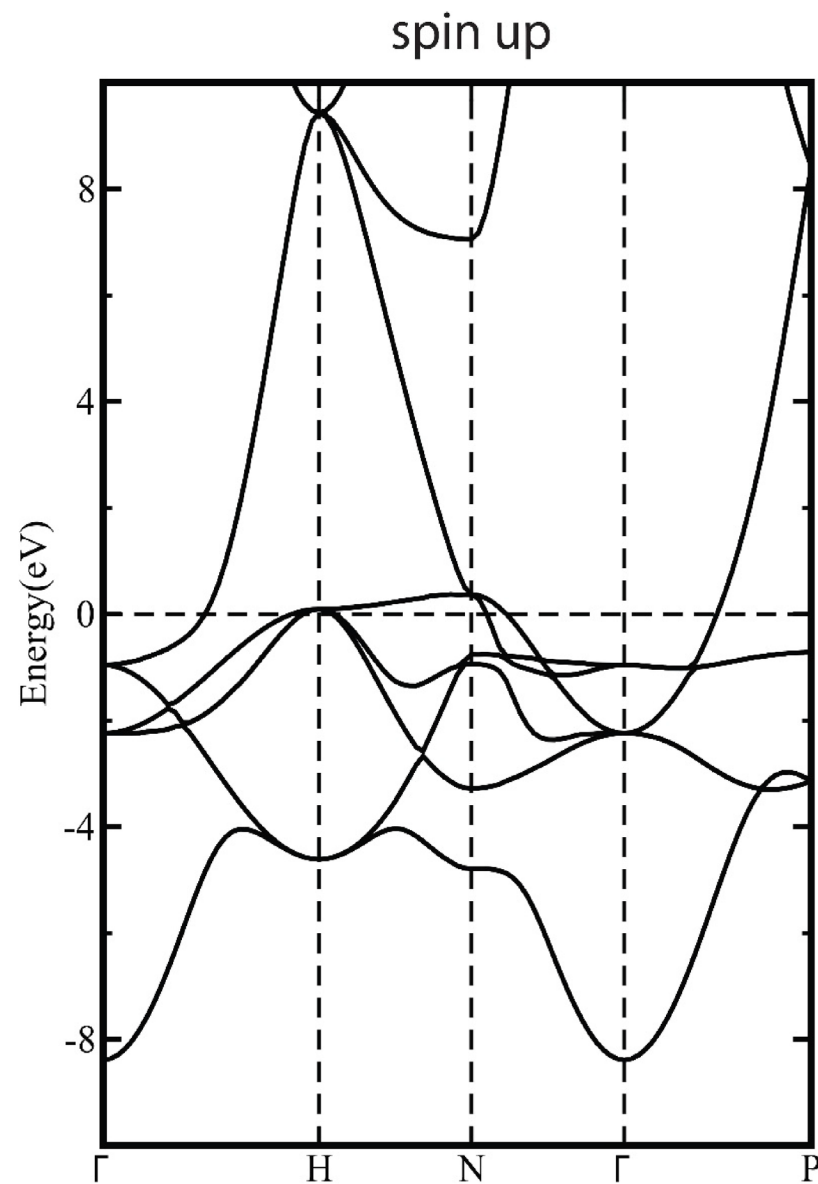


Bands of Si



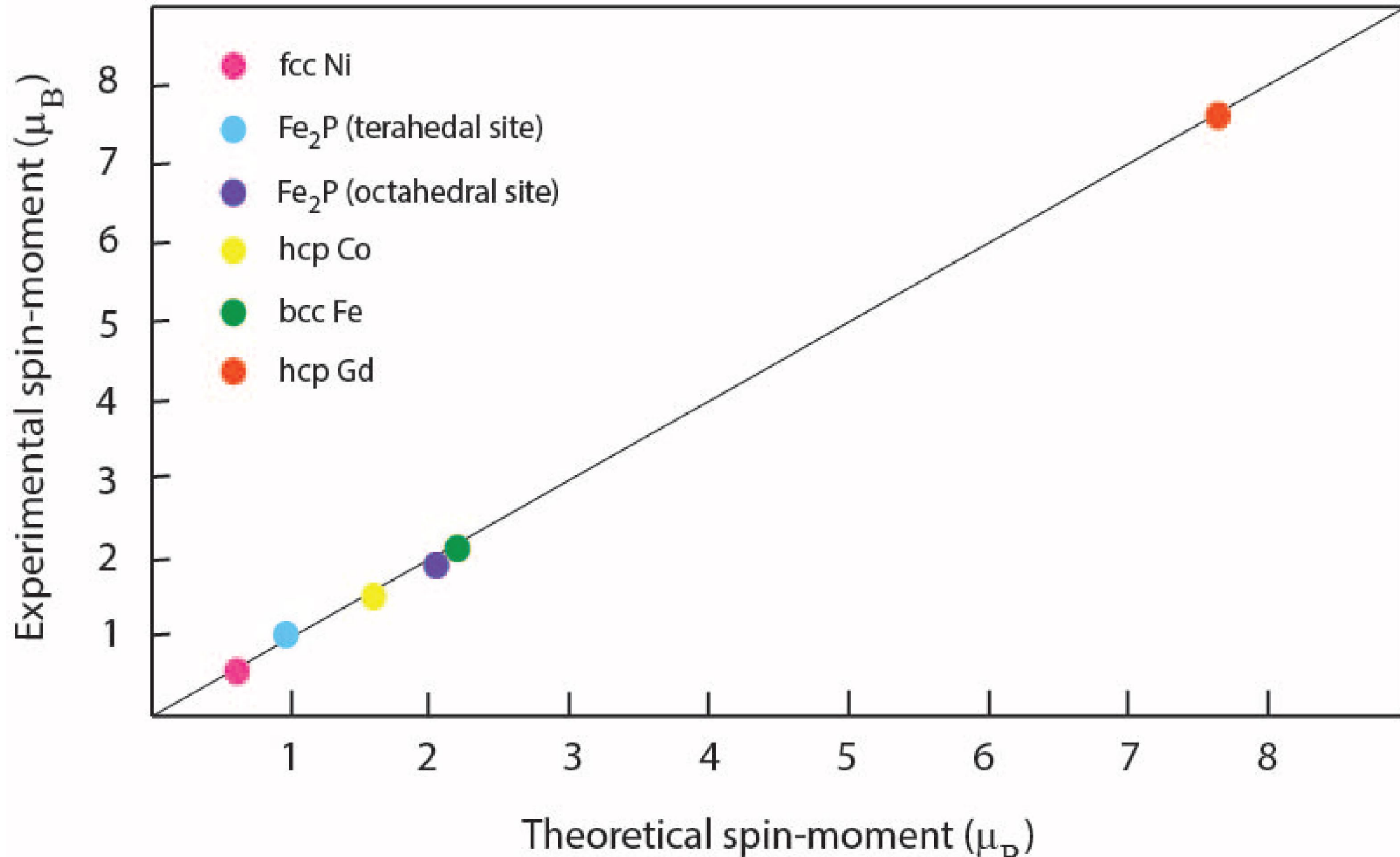


Bands of bcc Fe



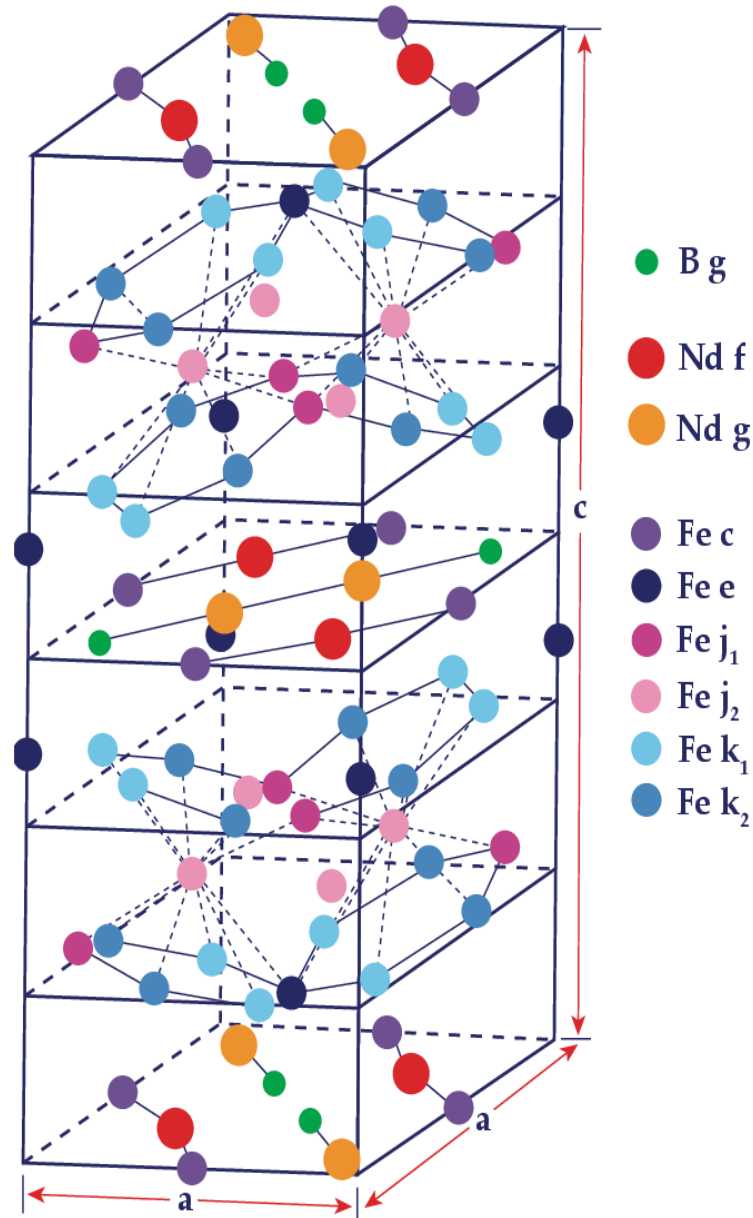


Calculations of magnetic moments



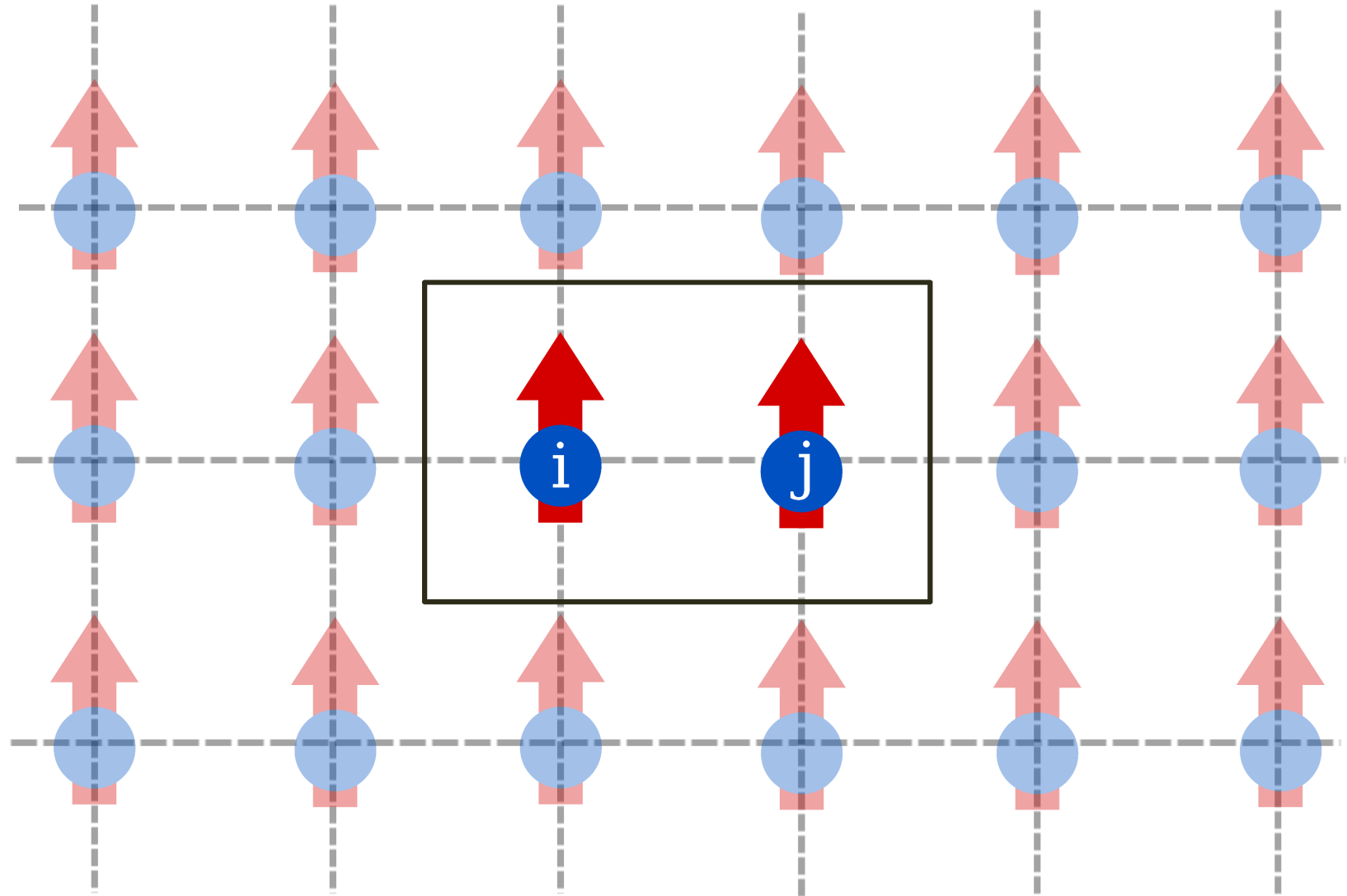


Permanent magnet applications



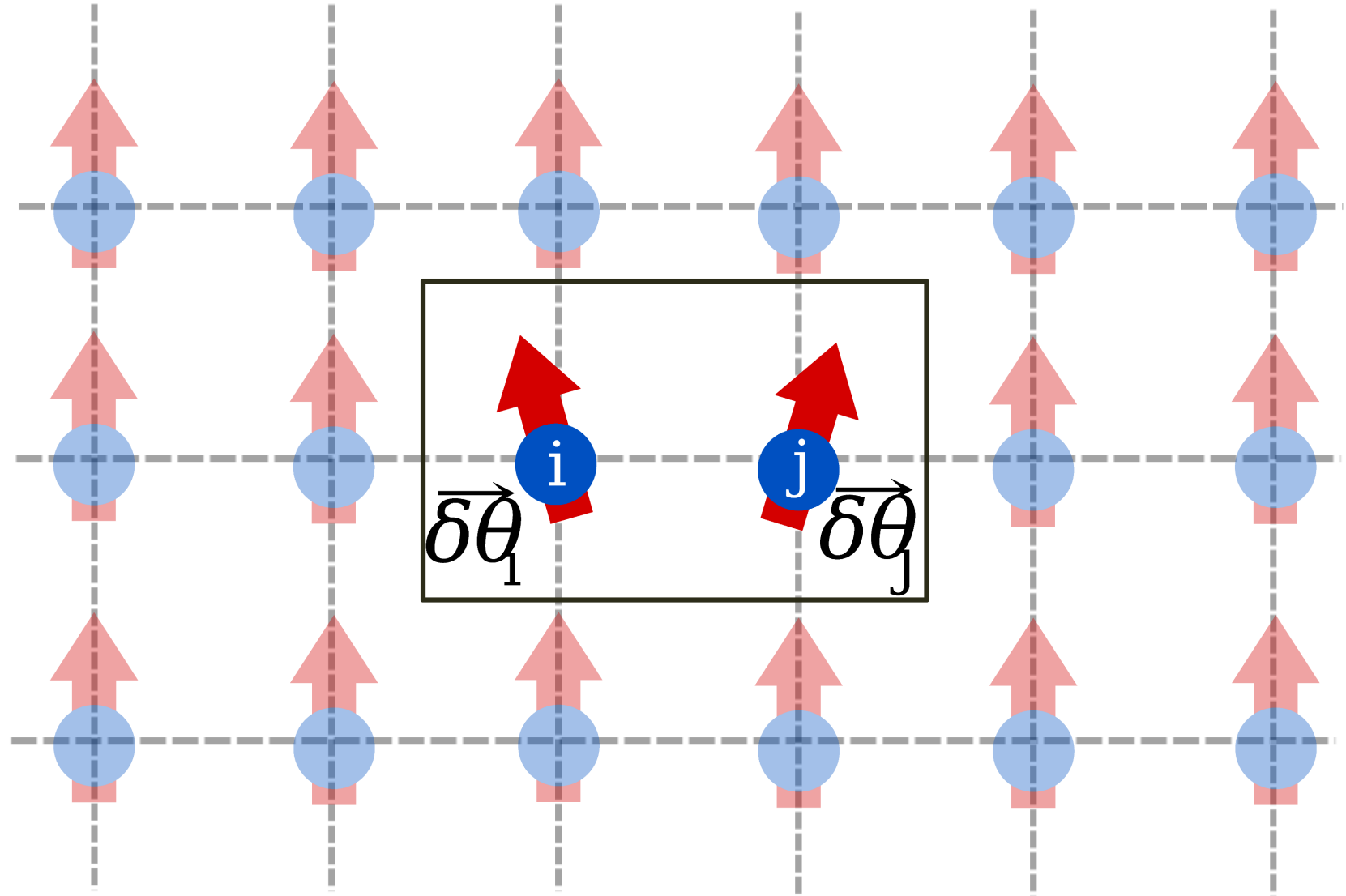


Inter-atomic exchange





Inter-atomic exchange





Heisenberg Hamiltonian

$$\Delta E = -J_{ij} \mathbf{m}_i \mathbf{m}_j$$



$$H = -J_{ij} \mathbf{m}_i \mathbf{m}_j$$

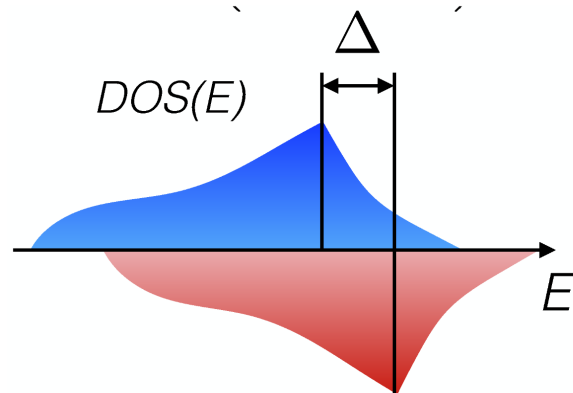


Expansion of the Hamiltonian in $\delta\vec{\theta}_i$ and $\delta\vec{\theta}_j$ gives

$$J_{ij} = \frac{-1}{4\pi} \int_{-\infty}^{E_F} \delta\epsilon \text{Tr}_m \left[\Delta_i \cdot G_{ij}^{\uparrow}(\epsilon) \cdot \Delta_j \cdot G_{ij}^{\downarrow}(\epsilon) \right]$$

Local exchange field

$$\Delta_i = (\hat{H}_i^{\uparrow} - \hat{H}_i^{\downarrow})$$



Inter-site Greens function

$$G_{ij}^{\sigma} = \langle i | \hat{G}(z) | j \rangle = \left\langle i \left| \frac{1}{z - \hat{H}^{\sigma}} \right| j \right\rangle$$

Lichtenstein *et al* JMMM **67** 65 (1987),
Katsnelson *et al* PRB **61** 8906 (2000),
Kvashnin *et al* PRB **91** 125133 (2015)
A.Szilva *et al*. RMP 95, 035004 (2023)



hcp Gd

Ferromagnetic metal ($T_c=292$ K)

