

Stabilizing skyrmions in multiferroics materials

Bertrand Dupé

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www.tom.uliege.be

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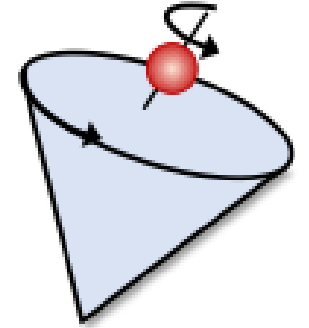
Lobbying and networking – STF, magnetometry...



European School on Magnetism 2026

Spin-orbit driven magnetic phenomena

Ångström laboratory, Uppsala University, Sweden
August 17-28, 2026



Enjoy ten days of fascinating topics around magnetism at the oldest university of Scandinavia. Lectures and hands-on sessions will cover fundamental aspects, theoretical modelling as well as experimental characterization of a variety of topics:

- Nanomagnetism
- Ultrafast magnetization dynamics
- Spin-orbit torques
- Spin-dependent transport
- Spin and orbital Hall effect
- Skyrmions
- Permanent magnets
- 2D magnets
- AFM spintronics
- Altermagnetism...



Registration fee (including accommodation and lunch)

Onsite: 700 €, Online: 300 €

Information: [WEBSITE](#)

Organization:

ESM chair: Assoc. Prof. Biplab Sanyal

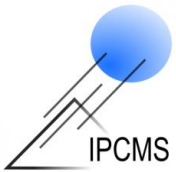
ESM co-chair: Dr. Heike Herper



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Charles Paillard
Markus Hoffmann
Stephan von Malottki



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Marie Hervé
Loic Mougél



Geoffrey Beach
Felix Bütner
Ivan Lemesh



Stefan Blügel
Gustav Bihlmayer
Yuriy Mokrousov
Samir Lounis



Michel Viret
Jean-Yves Chateau
Vincent Garcia
Aurore Finco
Vincent Jacques



Bin Xu



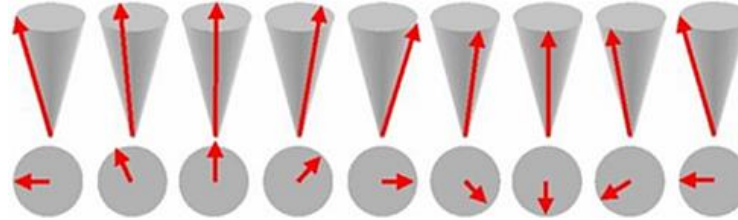
Laurent Bellaïche

Topological orders in Matter (TOM – tom.uliege.be)

**Magnonics and
thermal driven long
range excitations**

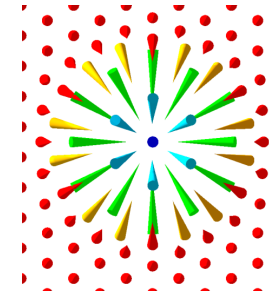
Mexx Regout

In collaboration with
Matthieu Verstraete



**Topological magnetic
textures in real space
(skyrmions, merons...)**

Sebastian Meyer
Olivier Gonnet

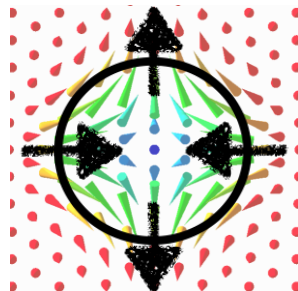


Topological orders in matter

[Github.com/bertdupe/Matjes](https://github.com/bertdupe/Matjes)

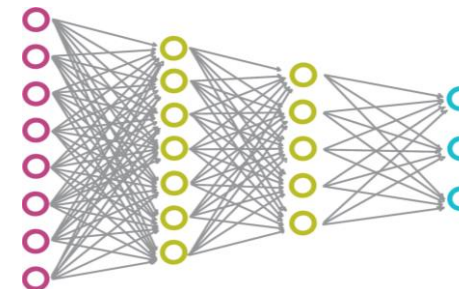
**Topology and transport
properties (Majorana
states...)**

In collaboration with
Marie Hervé
Tristan Cren
Wulf Wulfhekel



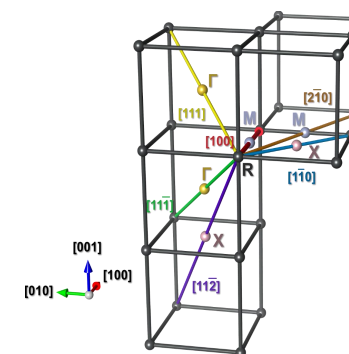
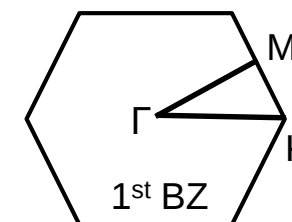
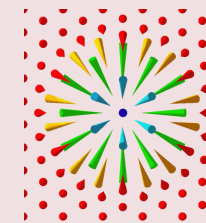
**Artificial intelligence
and topology**

Xavier Bosch



Outline

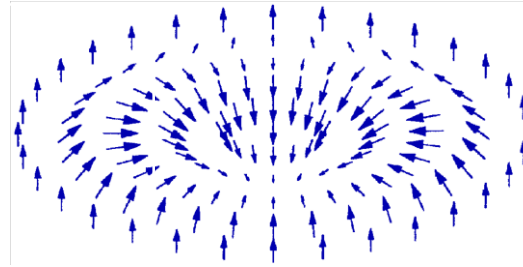
- Introduction on magnetic skyrmions
 - The discovery
 - The racetrack memory
 - Topology and stability
- The different magnetic interactions
 - Origin of the DM interaction
 - Obtaining magnetic interactions from density functional theory
 - Some examples of use
- Isolated skyrmions in Multiferroics
 - BiFeO_3 (BFO) as a good playground
 - Determining the magnetic interactions in BFO
 - Using strain to tune the interactions in BFO



Skyrmions in bulk magnetic materials

Micromagnetic model prediction

Existence of magnetic skyrmions
on a micrometer length scale

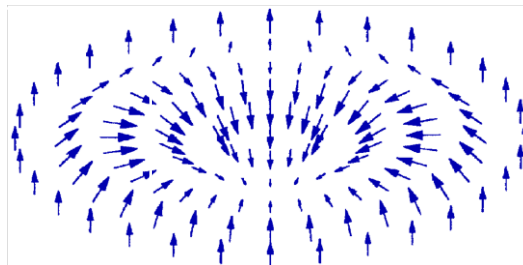


A. N. Bogdanov & D. A. Yablonskii, *Sov. Phys. JETP* **68**, 101 (1989)

Skyrmions in bulk magnetic materials

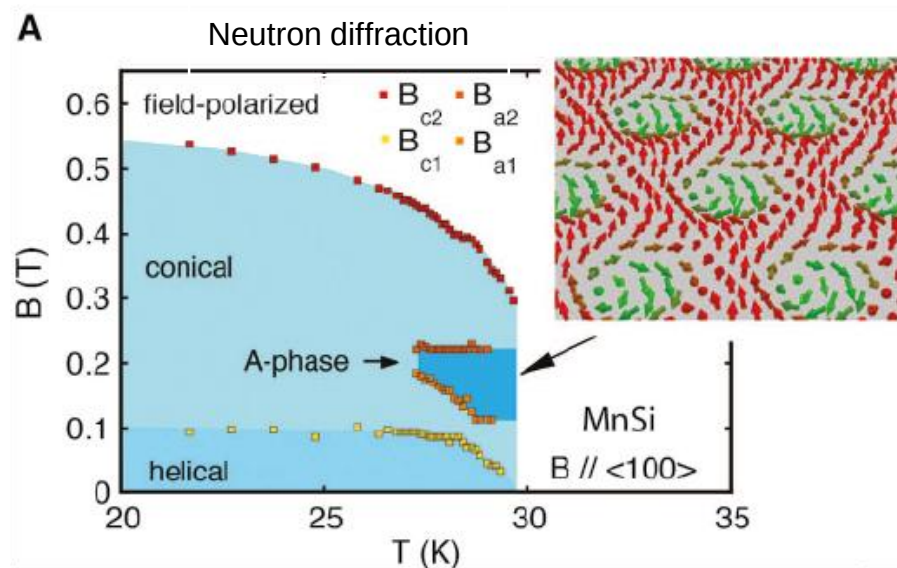
Micromagnetic model prediction

Existence of magnetic skyrmions
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A. N. Bogdanov & D. A. Yablonskii, *Sov. Phys. JETP* **68**, 101 (1989)

Experimental discovery



S. Mühlbauer *et al.*, *Science* **323**, 915 (2009)

A. Neubauer *et al.*, *PRL* **102**, 186602 (2009)

X. Z. Yu *et al.*, *Nature* **465**, 901 (2010)

M. Lee *et al.*, *PRL* **102**, 186601 (2009)

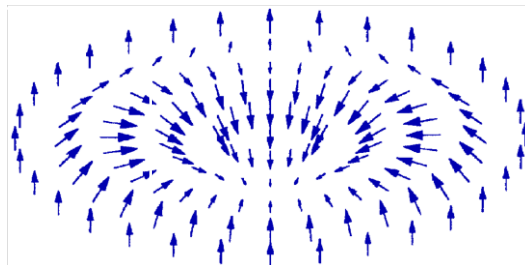


P. Milde *et al.*, *Science* **340**, 1076 (2013)

Skyrmions in bulk magnetic materials

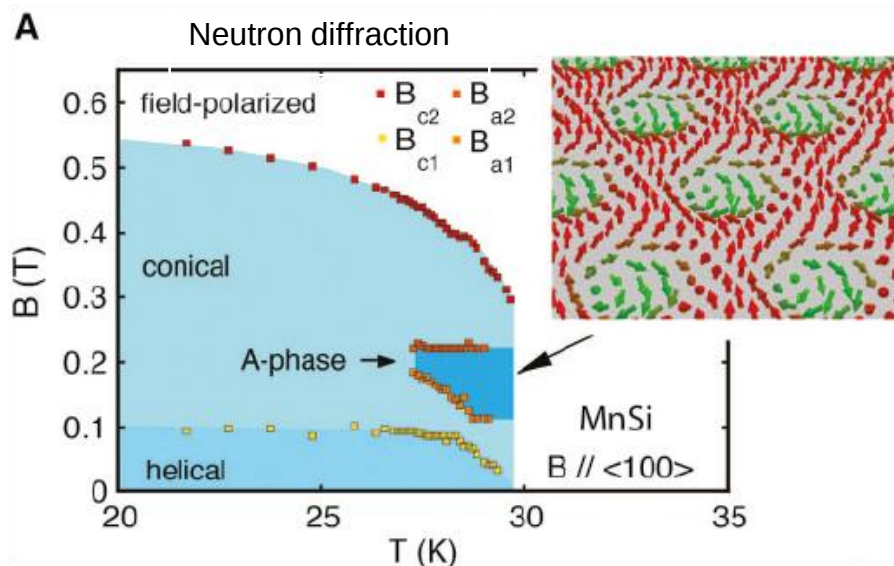
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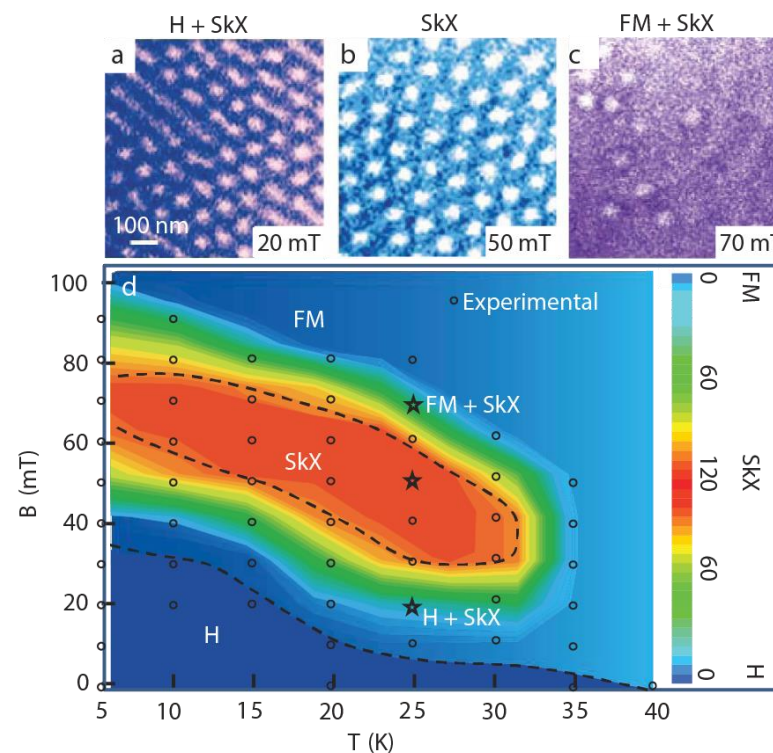


A. N. Bogdanov & D. A. Yablonskii, *Sov. Phys. JETP* **68**, 101 (1989)

Experimental discovery



Lorentz TEM



S. Mühlbauer *et al.*, *Science* **323**, 915 (2009)

X. Z. Yu *et al.*, *Nature* **465**, 901 (2010)

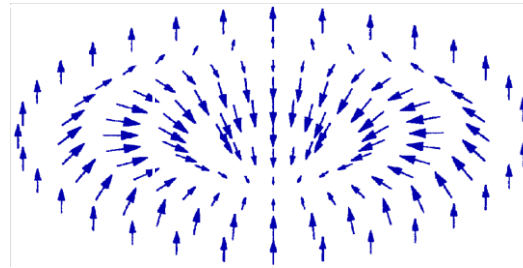
A. Neubauer *et al.*, *PRL* **102**, 186602 (2009)

M. Lee *et al.*, *PRL* **102**, 186601 (2009)

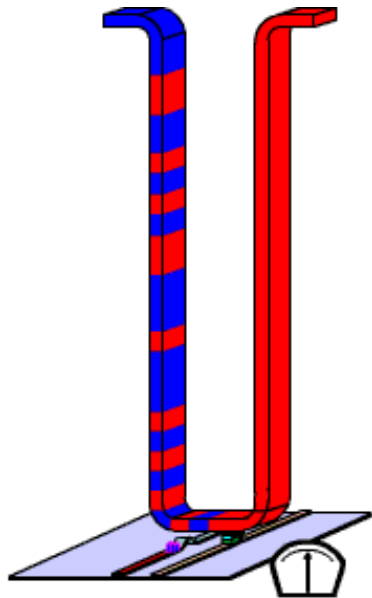
Skyrmions in bulk magnetic materials

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A. N. Bogdanov & D. A. Yablonskii, *Sov. Phys. JETP* **68**, 101 (1989)



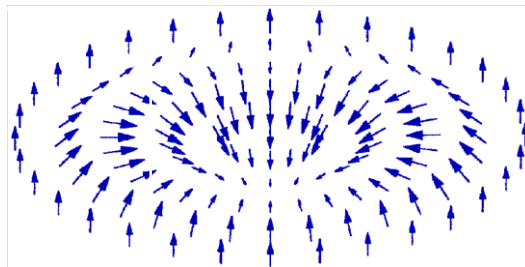
race-track

S. Parkin *et al.*, *Science* **320**, 190 (2008)

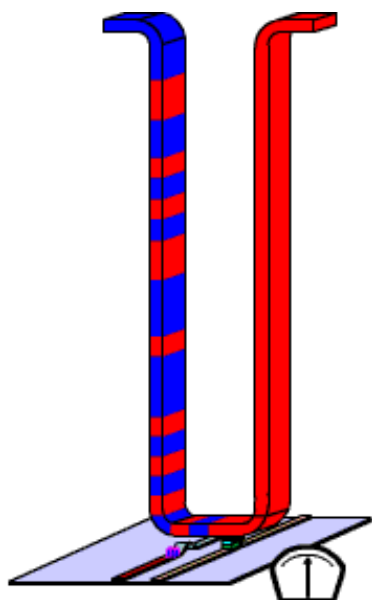
Skyrmions in bulk magnetic materials

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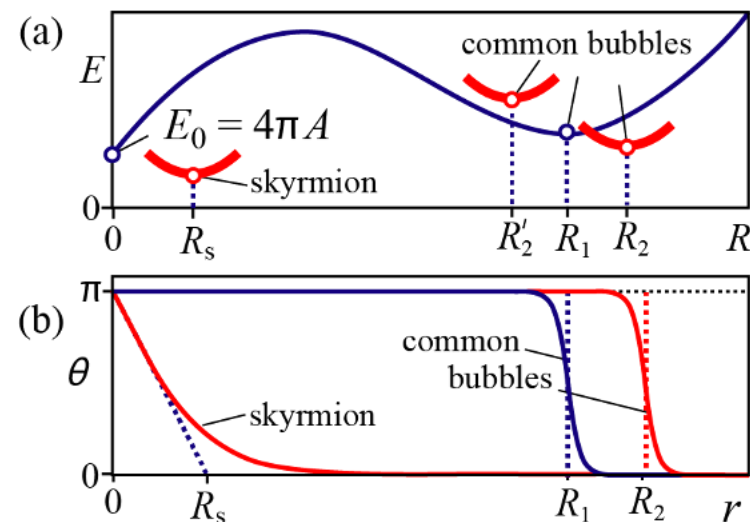
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A. N. Bogdanov & D. A. Yablonskii, *Sov. Phys. JETP* **68**, 101 (1989)



Skyrmion race-track



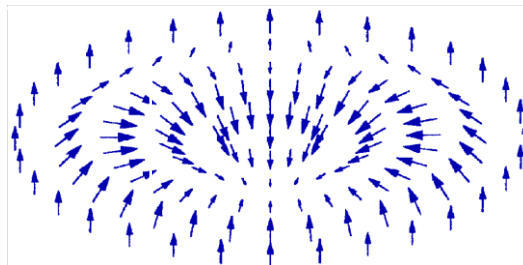
S. Parkin *et al.*, *Science* **320**, 190 (2008)
A. Fert *et al.*, *Nature Nanotech.* **8**, 152 (2013)
W. Jiang *et al.*, *Science* **349**, 283 (2015)
K. Litzius *et al.*, *Nature Phys.* **13**, 170 (2016)

N. S. Kiselev *et al.*, *J. Phys D* **44**, 39201 (2011)

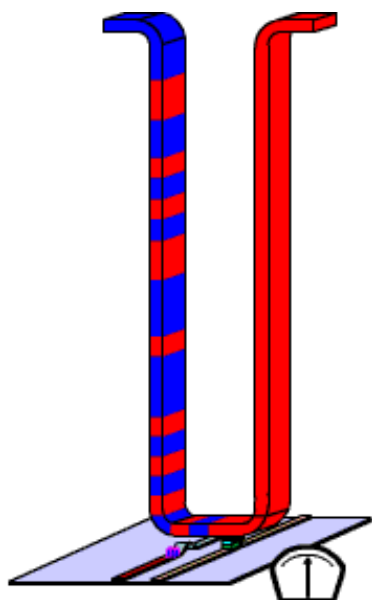
Skyrmions in bulk magnetic materials

Micromagnetic model prediction

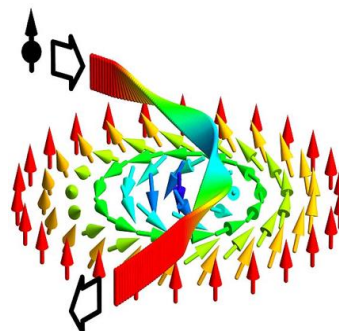
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A. N. Bogdanov & D. A. Yablonskii, *Sov. Phys. JETP* **68**, 101 (1989)



Skyrmion race-track

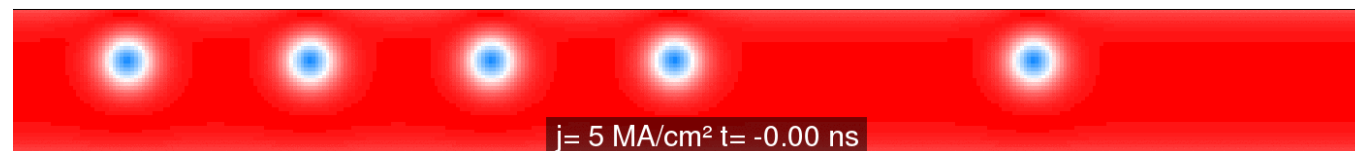


$$\mathbf{F}_{\sigma k} = e\mathbf{E} + \mathbf{F}_H + q_{\sigma}^e \cdot \mathbf{v}_k \times \mathbf{B}^e$$

$$\mathbf{B}^e = \frac{\hbar}{2} \mathbf{n} (\nabla_x \mathbf{n} \times \nabla_y \mathbf{n})$$

$$\mathbf{E} = \hbar \mathbf{n} (\nabla_r \mathbf{n} \times \nabla_t \mathbf{n})$$

T. Schulz *et al.*, *Nature Physics* **8**, 301 (2012)



S. Parkin *et al.*, *Science* **320**, 190 (2008)
A. Fert *et al.*, *Nature Nanotech.* **8**, 152 (2013)
W. Jiang *et al.*, *Science* **349**, 283 (2015)
K. Litzius *et al.*, *Nature Phys.* **13**, 170 (2016)

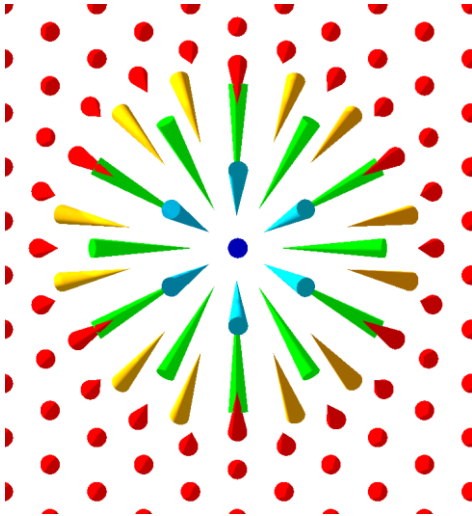
Skyrmions and topological charge

Skyrmion number (topological charge) of a vector field $\mathbf{n}(x,y)$:

$$S = \frac{1}{4\pi} \int \mathbf{n} \cdot (\nabla_x \mathbf{n} \times \nabla_y \mathbf{n}) dA$$

Ferromagnet ($S=0$): topologically trivial state

Skyrmion ($S=+1$)



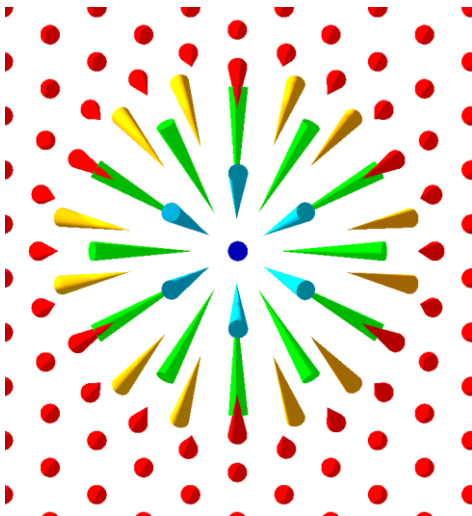
Skyrmions and topological charge

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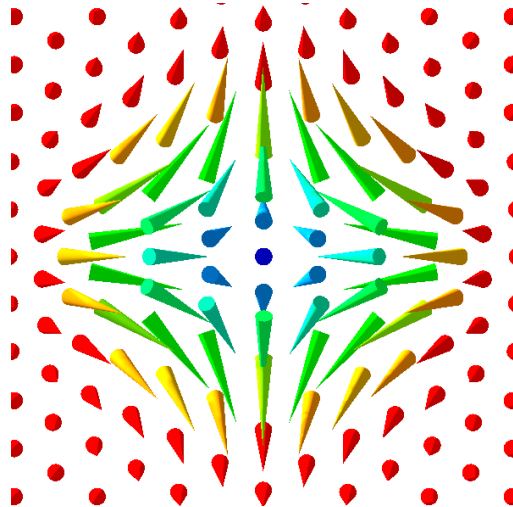
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Ferromagnet ($S=0$): topologically trivial state

Skyrmion ($S=+1$)

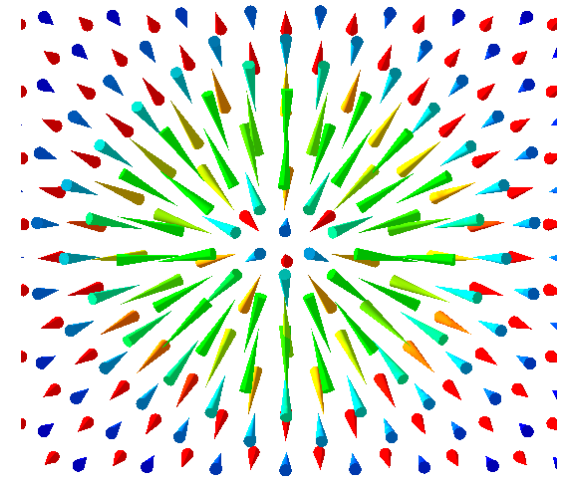


Antiskyrmion ($S=-1$)



L. Camosi *et al.*, *Phys. Rev. B* **95**, 214422 (2017)
M. Hoffmann *et al.*, *Nature Commun.* **8**, 308 (2017)
A. Nayak *et al.*, *Nature* **548**, 561 (2017)

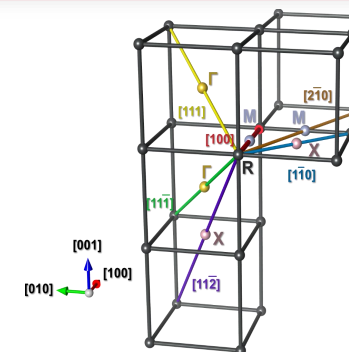
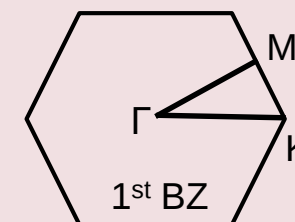
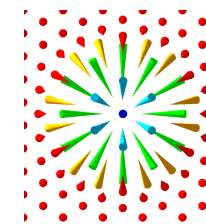
Skyrmion AFM ($S=+1$)



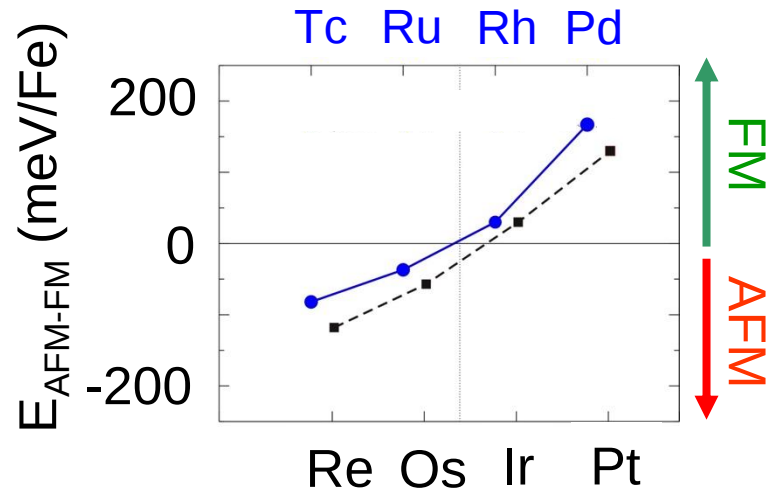
J. Barker *et al.*, *Phys. Rev. Lett.* **116**, 147203 (2016)
S. Meyer *et al.*, *Phys. Rev. B* **96**, 094408 (2017)
W. Legrand *et al.*, *Nat. Materials* (2019)
T. Dohi *et al.*, *Nat. Commun* (2019)
L. Desplat *et al.*, *Phys. Rev. B* **104**, L060409 (2021)

Outline

- Introduction on magnetic skyrmions
 - The discovery
 - The racetrack memory
 - Topology and stability
- The different magnetic interactions
 - Origin of the DM interaction
 - Obtaining magnetic interactions from density functional theory
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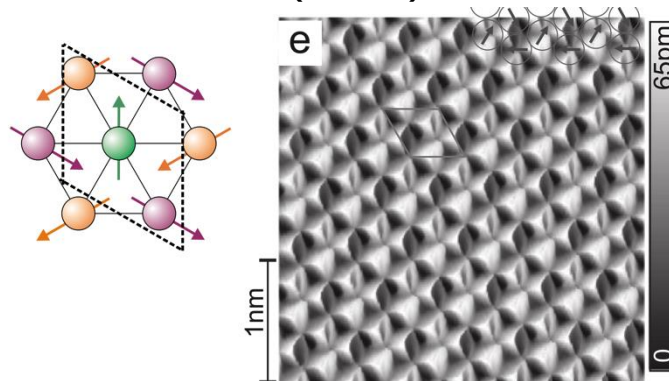


Tuning the magnetic exchange interaction in Fe monolayer



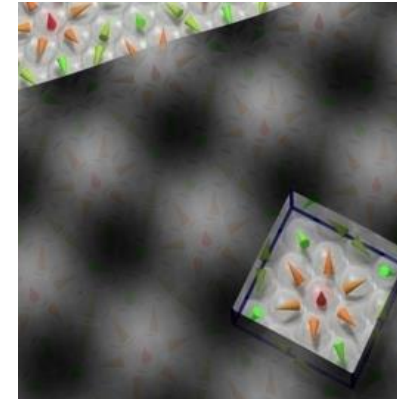
B. Hardrat *et al* PRB **79**, 094411 (2009)

Fe/Re(0001) Néel state



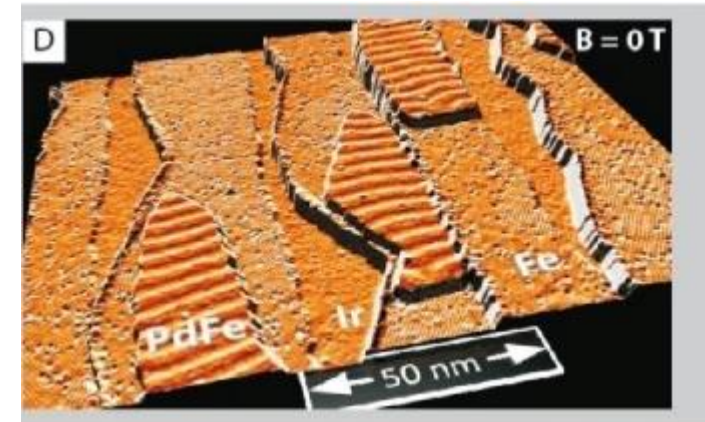
S. Ouazi *et al* Sur. sci. **630**, 280 (2014)
S. Ouazi *et al* PRL **112**, 076102 (2014)

Fe/Ir(111): nano-skyrmion lattice



S. Heinze *et al* Nat. Phys. **7**, 713 (2011)

Pd/Fe/Ir(111) spin spiral ground state



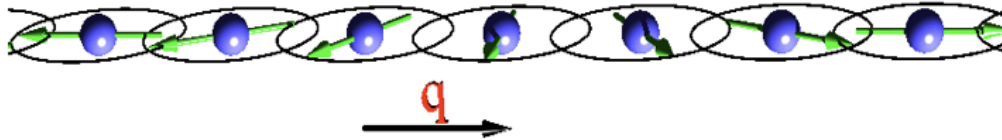
N. Romming *et al* Science **341**, 636 (2013)
B. Dupé *et al* Nature Comm. **5**, 4030 (2014)

Spin spiral energy calculation

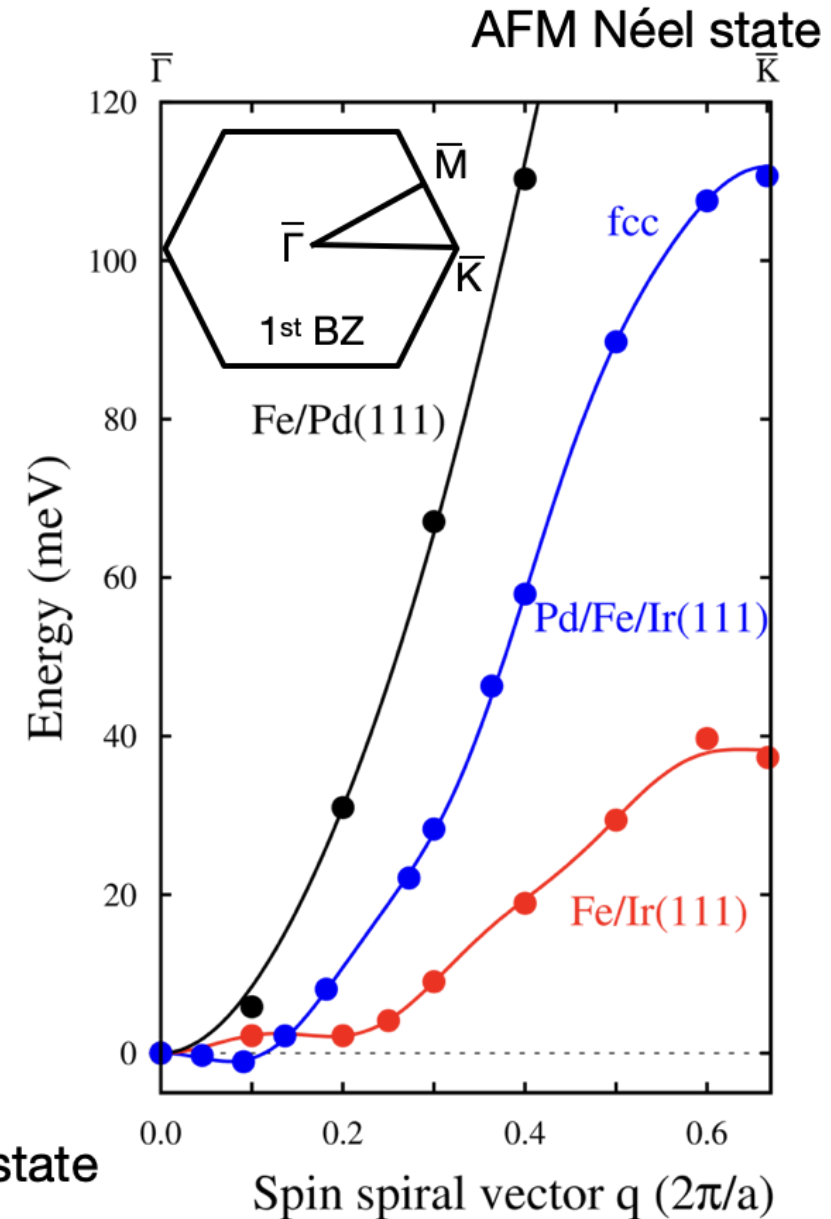
Density-functional theory (DFT) using the FLEUR code:

- energy of non-collinear magnetic structures
- energies of spiral spin-density waves

Spin spirals



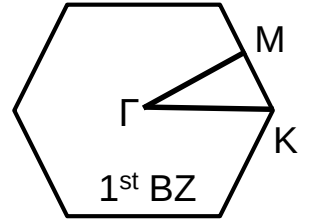
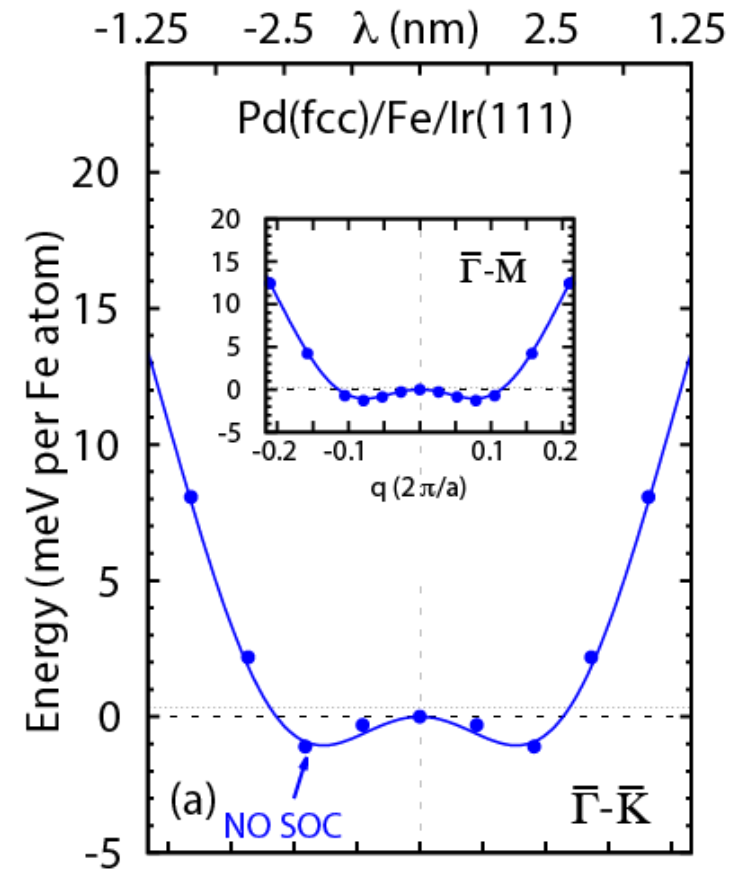
$$E(\mathbf{q}) = -N \sum_{\delta} J_{0,\delta} e^{-i\mathbf{R}_{\delta}\mathbf{q}}$$



Heisenberg Hamiltonian parametrization Pd/Fe/Ir(111)

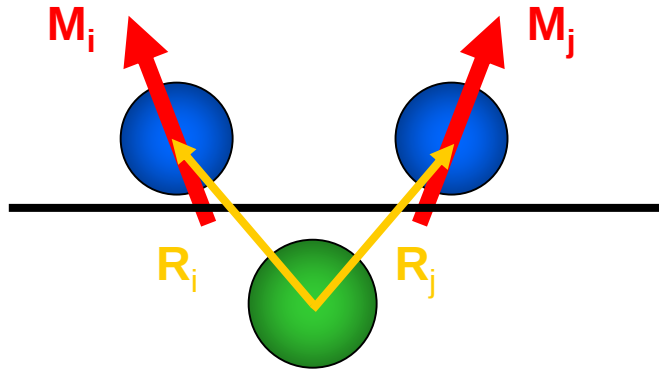
Spin Hamiltonian

$$H = -\sum_{i,j} \mathbf{M}_i J_{ij} \mathbf{M}_j$$



Origin of the DM interaction at metal surface

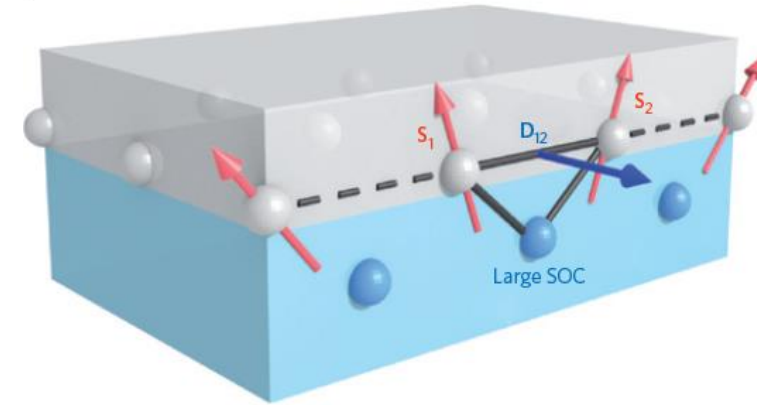
Interaction of two magnetic atoms via a surface atom with strong spin-orbit coupling



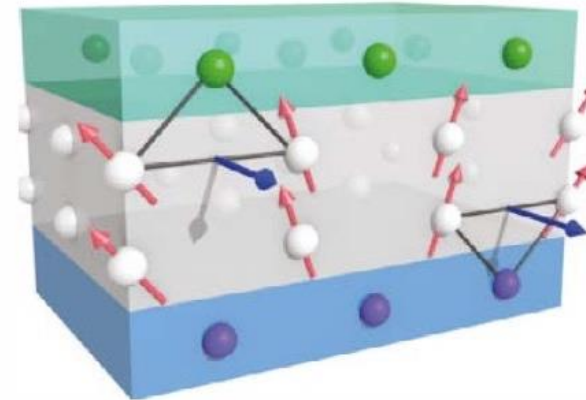
$$H_{DM} \propto -V(\xi)(\mathbf{R}_i \times \mathbf{R}_j)(\mathbf{M}_i \times \mathbf{M}_j)$$

$$H_{DM} \propto -\mathbf{D}_{ij}(\mathbf{M}_i \times \mathbf{M}_j)$$

Levy & Fert, Phys. Rev Lett. **44**, 1538 (1980)



J. Sampaio *et al.*, Nature Nanotech. **8**, 839 (2013)
O. Boulle *et al.*, Nature Nanotech. **5**, 449 (2016)



C. Moreau-Luchaire *et al.*, Nature Nanotech. **11**, 444 (2016)

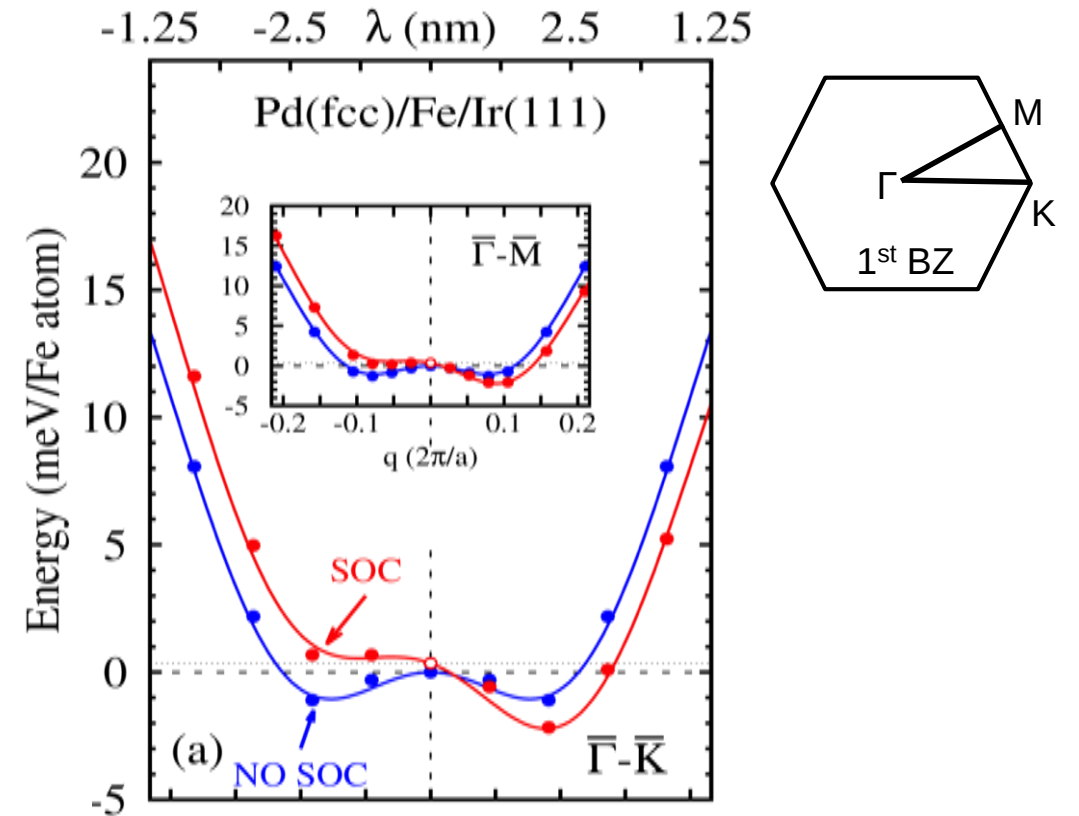
Heisenberg Hamiltonian parametrization: DMI

Spin Hamiltonian

$$H = -\sum_{i,j} \mathbf{M}_i J_{ij} \mathbf{M}_j$$

$$-\sum_i K M_z^2$$

$$-\sum_{i,j} \mathbf{D}_{ij} \cdot (\mathbf{M}_i \times \mathbf{M}_j)$$



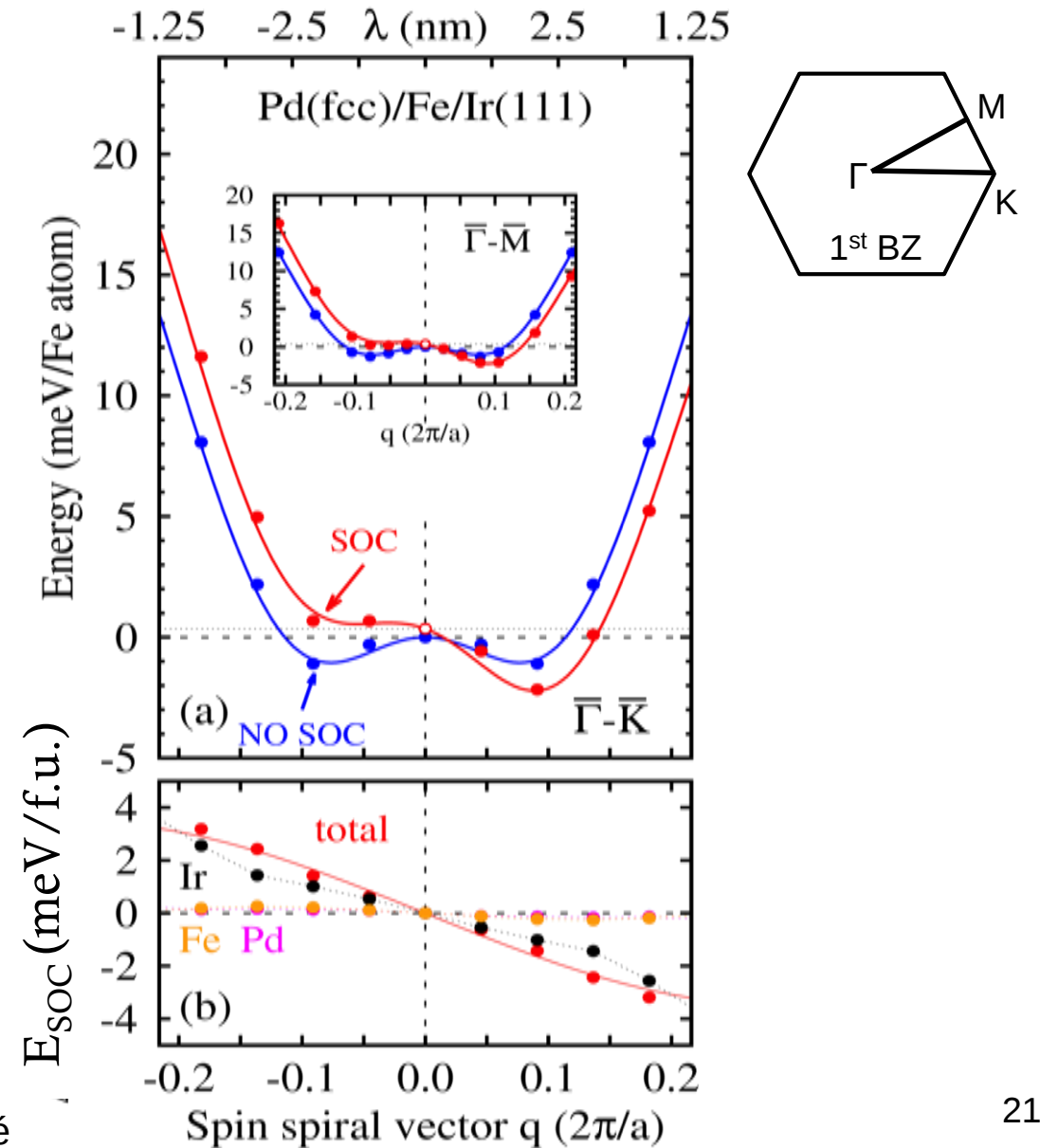
Heisenberg Hamiltonian parametrization: DMI

Spin Hamiltonian

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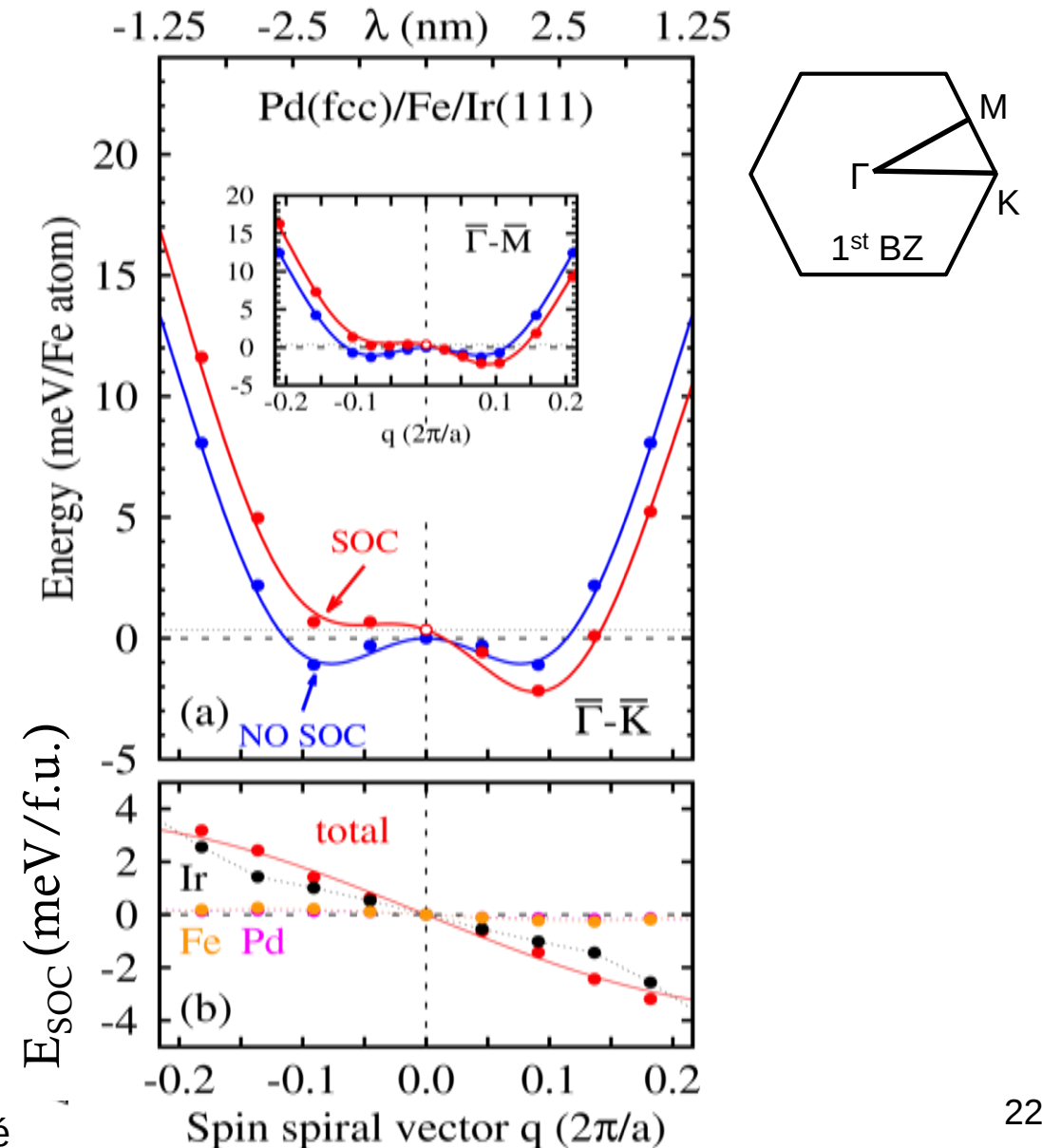
$$-\sum_{i,j} \mathbf{D}_{ij} \cdot (\mathbf{M}_i \times \mathbf{M}_j)$$



Heisenberg Hamiltonian parametrization

Spin Hamiltonian solved by Monte-Carlo & spin dynamics

$$\begin{aligned}
 H = & -\sum_{i,j} \mathbf{M}_i J_{ij} \mathbf{M}_j \\
 & -\sum_i K M_z^2 \\
 & -\sum_{i,j} \mathbf{D}_{ij} \cdot (\mathbf{M}_i \times \mathbf{M}_j) \\
 & -\sum_i \mathbf{M}_i \mathbf{B} \\
 & -\sum_{i,j} B_{ij} (\mathbf{M}_i \cdot \mathbf{M}_j)^2 + 4\text{spin}
 \end{aligned}$$



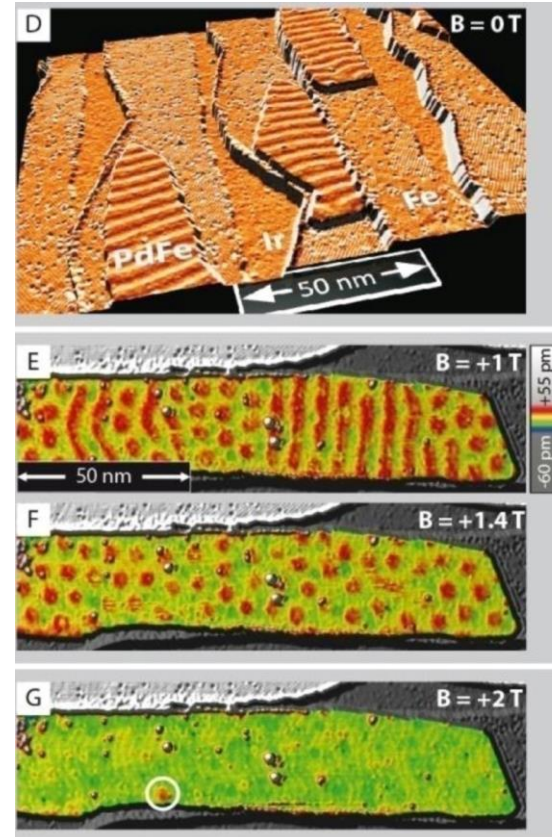
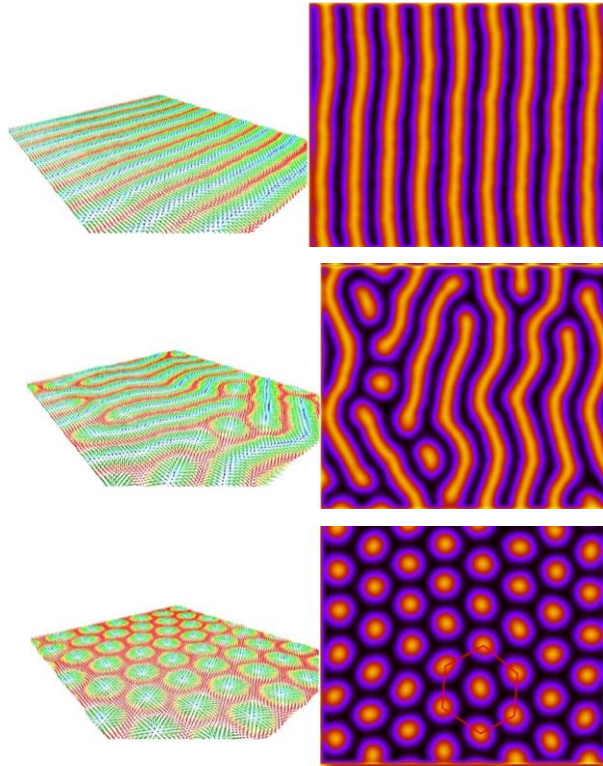
Phase diagram of Pd/Fe(fcc)/Ir(111) at 0K

Spin spiral pitch from Monte Carlo: 5.4nm

Monte-Carlo
Simulation

Experiments

B

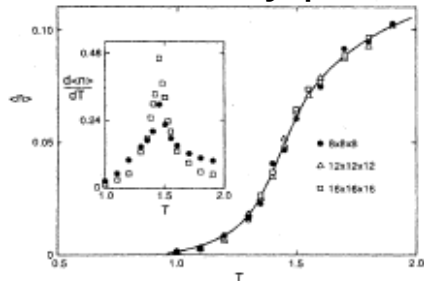


B. Dupé *et al* Nature Commun. **5**, 4030 (2014)
E. Simon *et al* Phys. Rev. B **90**, 094410 (2014)

N. Romming *et al* Science **341**, 636 (2013)
N. Romming *et al* Phys. Rev. Lett. **114**, 177203 (2015)

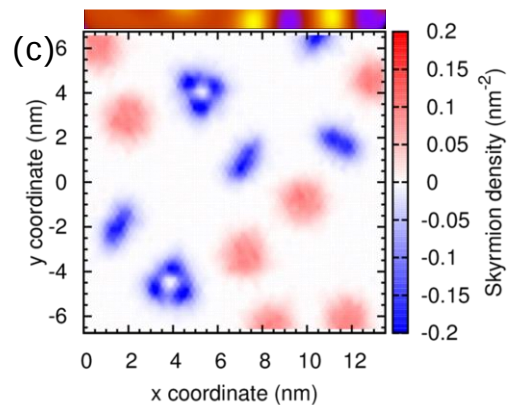
Phase diagram of Pd/Fe(fcc)/Ir(111)

Creation of skyrmion/antiskyrmion density pairs



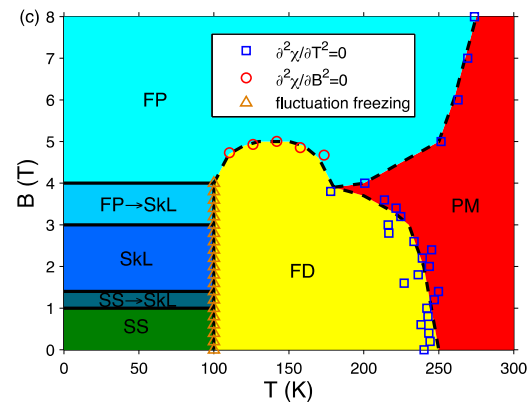
M. Lau *et al.*, Phys Rev B **39**, 7212 (1989)
Y. Nahas *et al.*, Phys. Rev. Lett. **119**, 117601 (2017)

Integrated skyrmion density is no order parameter



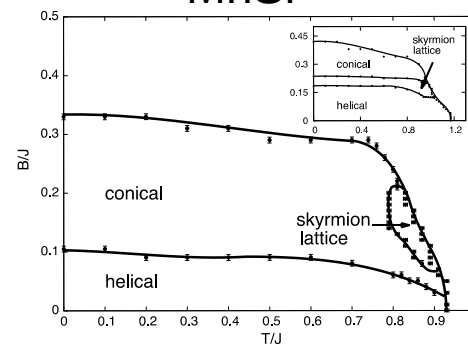
B. Dupé *et al.*, New J. Phys **18**, 055015 (2016)

Pd/Fe/Ir(111)



L. Rosza *et al.*, PRB **93**, 024417 (2016)

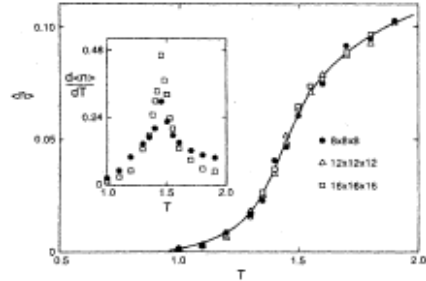
MnSi



S. Bührandt *et al.*, PRB **88**, 195137 (2013)

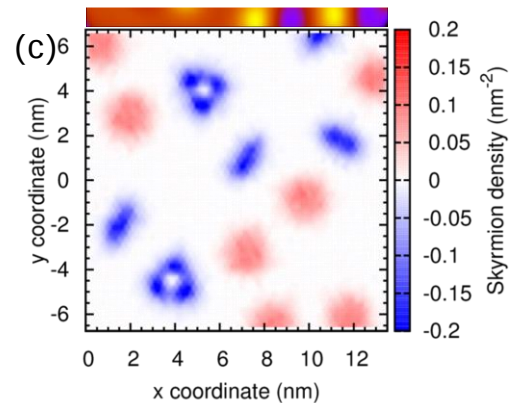
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Creation of skyrmion/antiskyrmion density pairs



M. Lau *et al.*, Phys Rev B **39**, 7212 (1989)
Y. Nahas *et al.*, Phys. Rev. Lett. **119**, 117601 (2017)

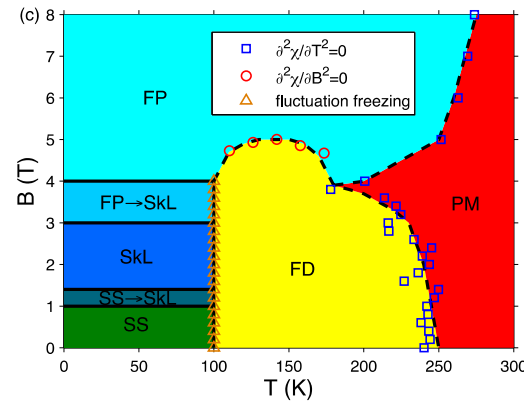
Integrated skyrmion density is no order parameter



B. Dupé *et al.*, New J. Phys **18**, 055015 (2016)

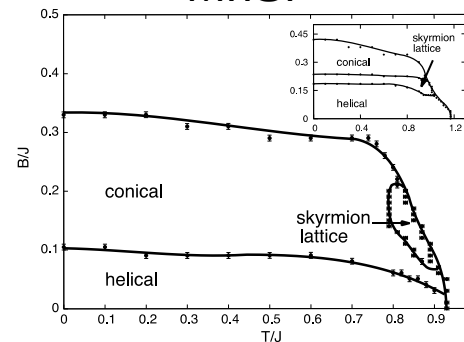


Pd/Fe/Ir(111)



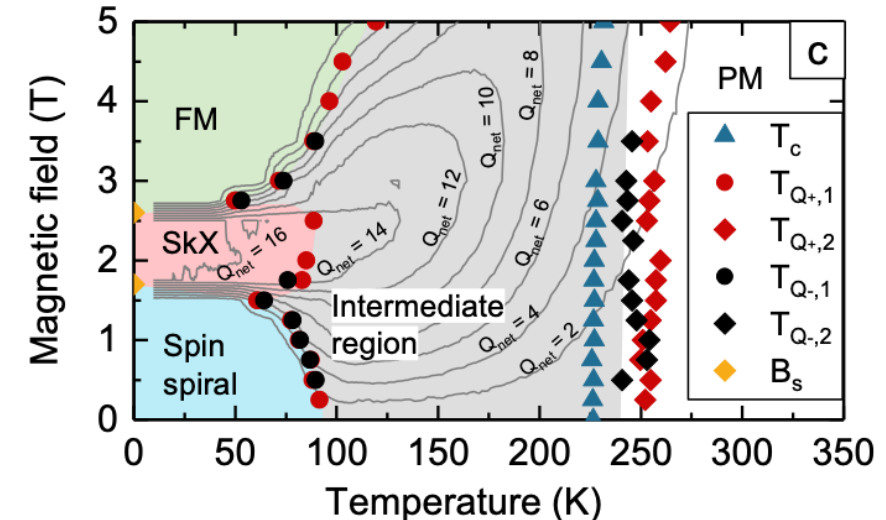
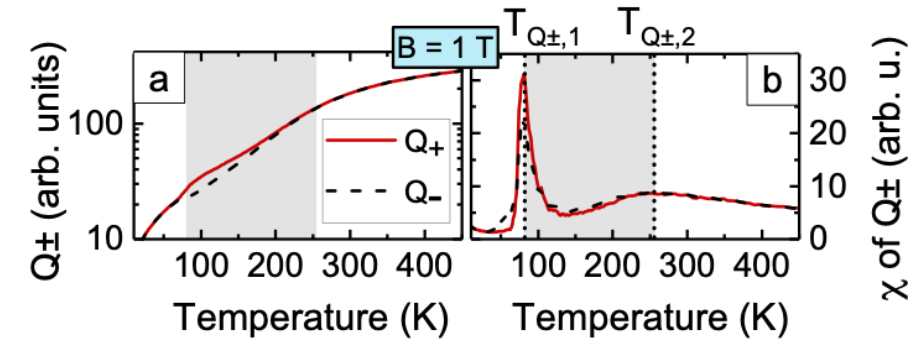
L. Rosza *et al.*, PRB **93**, 024417 (2016)

MnSi



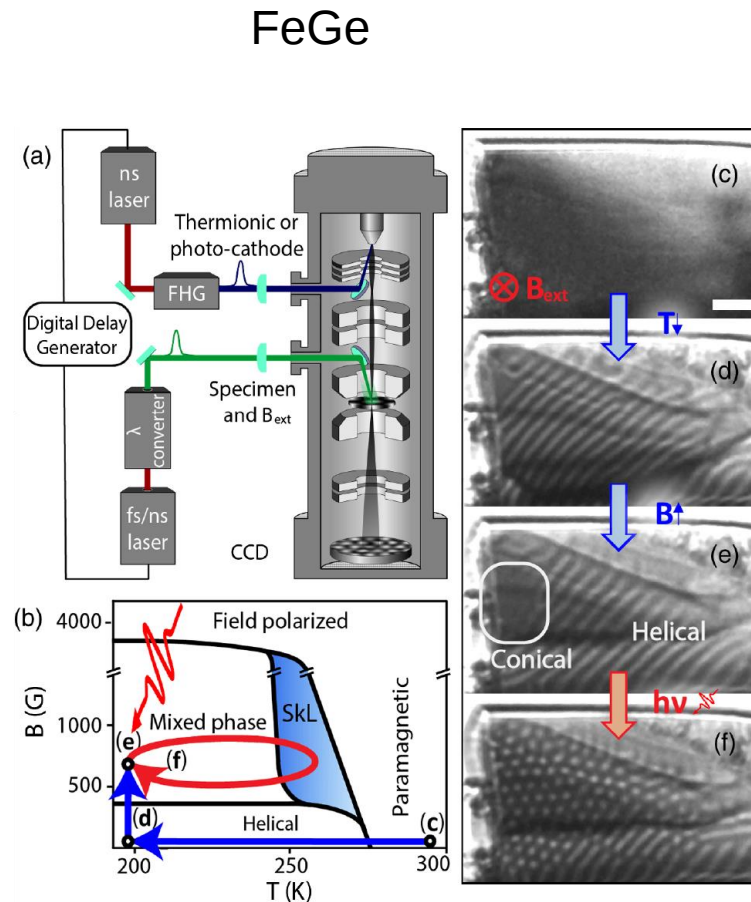
S. Bührandt *et al.*, PRB **88**, 195137 (2013)

Parallel tempering Monte Carlo



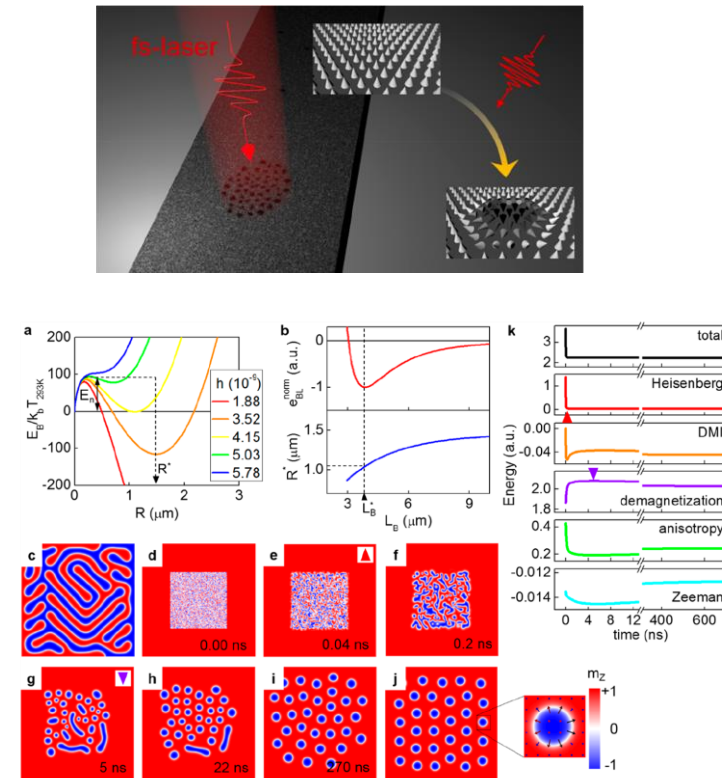
M. Böttcher *et al.*, New Journ. Phys. **20**, 103014 (2018)

Temperature skyrmion-antiskyrmion generation



G. Berruto *et al.* *Phys. Rev. Lett.* **120**, 117201 (2018)

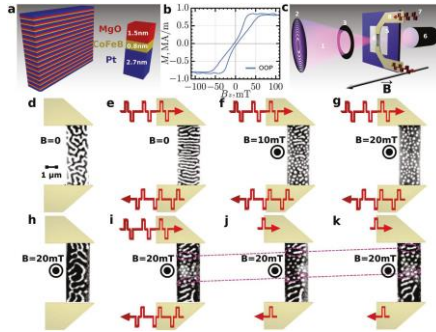
TaOx/CoFeB/Ta



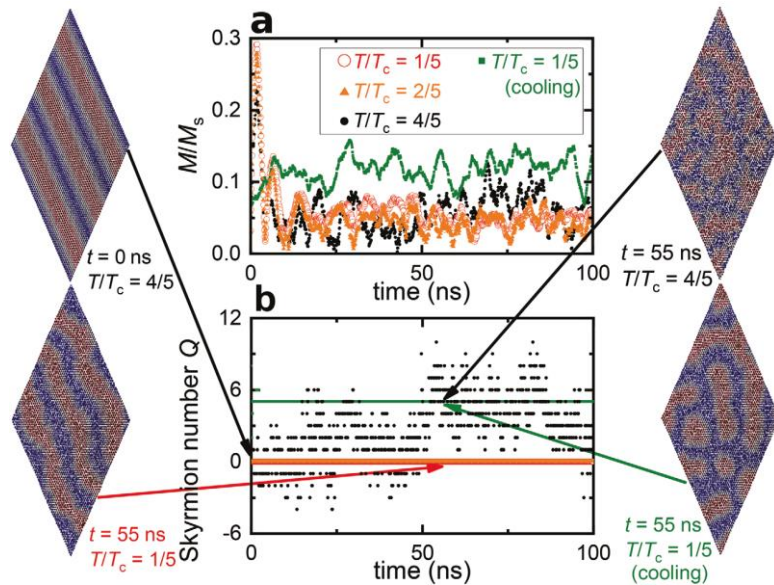
S.-G. Je *et al.* *NanoLetters* **18**, 7362 (2018)

Temperature skyrmion-antiskyrmion generation

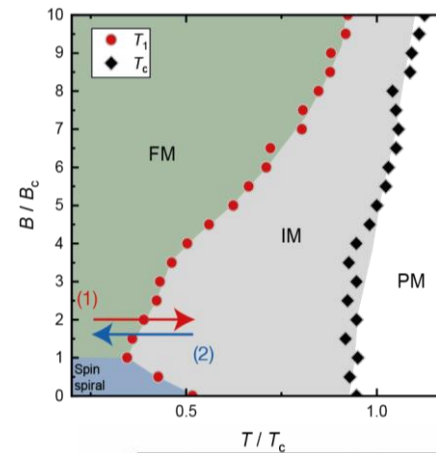
[MgO/Co/Pt]_n



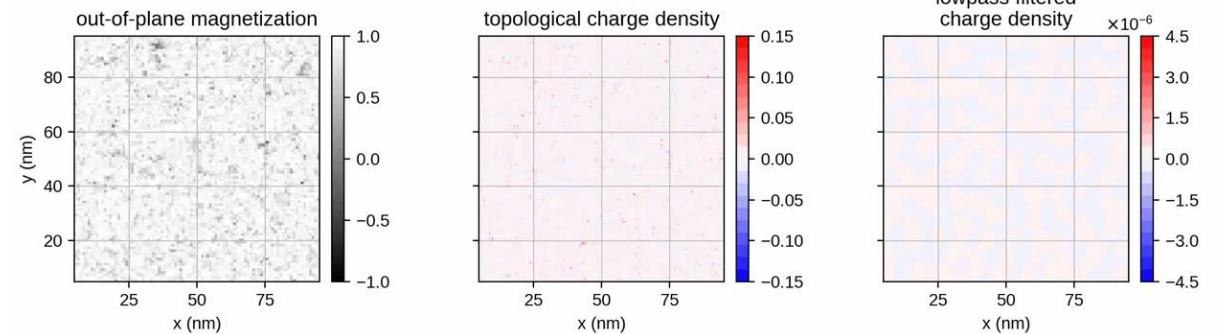
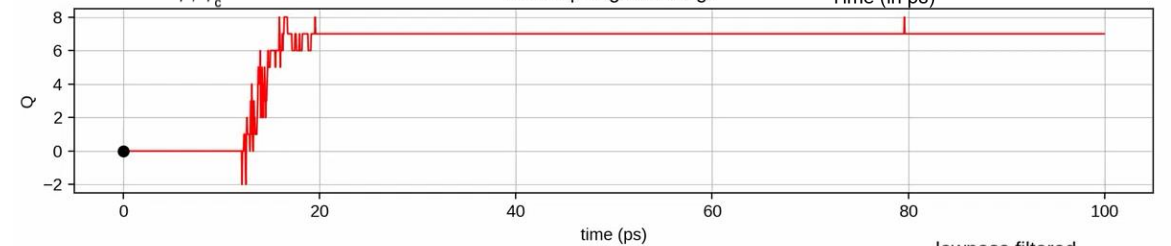
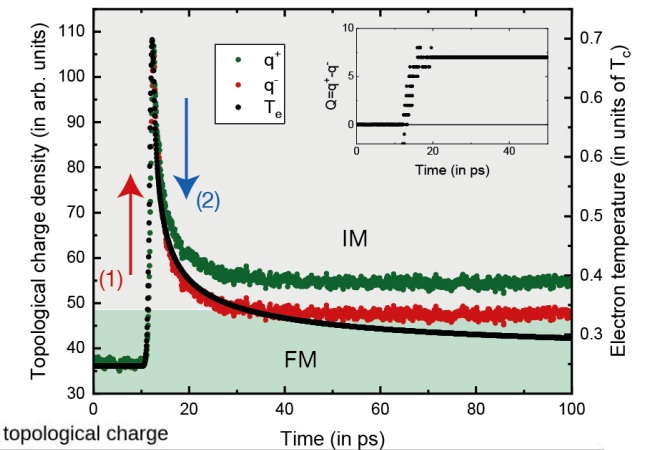
Magnetization dynamics simulations



Phase diagram of [MgO/Co/Pt]_n



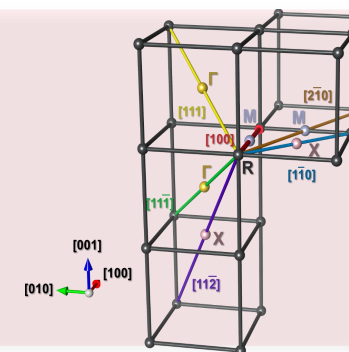
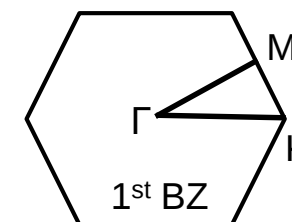
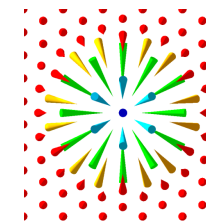
electron temperature



I. Lemesh, BD *et al.*, Advanced Materials **30**, 1805461 (2018)
F. Büttner *et al.*, Nature Mater. **20**, 30 (2021)

Outline

- Introduction on magnetic skyrmions
 - The discovery
 - The racetrack memory
 - Topology and stability
- The different magnetic interactions
 - Origin of the DM interaction
 - Obtaining magnetic interactions from density functional theory
 - Some examples of use
- Isolated skyrmions in Multiferroics
 - BiFeO_3 (BFO) as a good playground
 - Determining the magnetic interactions in BFO
 - Using strain to tune the interactions in BFO



Order parameters of BiFeO₃

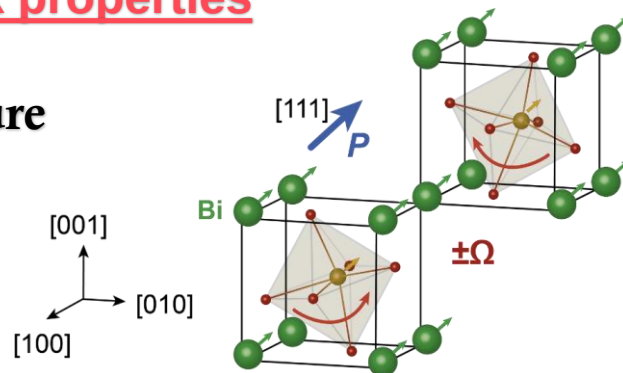
Bulk properties

1. Perovskite structure

Bi Cations

Fe Cations

O Anions



2. Fe-Bi displacement

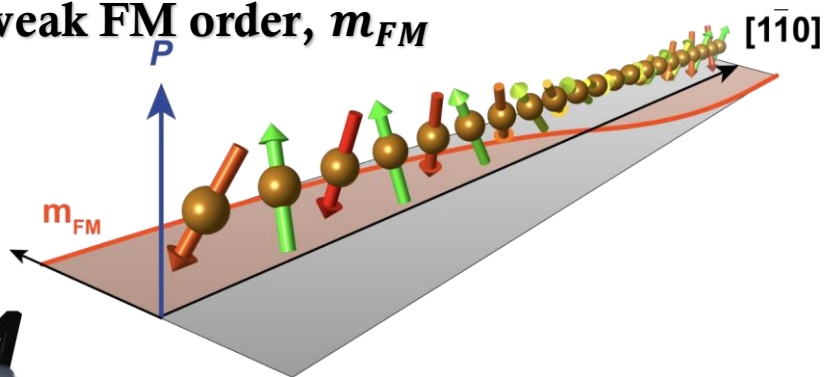
Polarization along [111] direction

$$P \approx \frac{90\mu\text{C}}{\text{cm}^2}, T_c = 1123 \text{ K}$$

$R3c$

3. Oxygen tilt in antiphase $\pm\Omega$

4. Local antiferromagnetic order with long-ranged spin cycloid, $\lambda \approx 62 \text{ nm}$ + weak FM order, m_{FM}



Order parameters of BiFeO₃

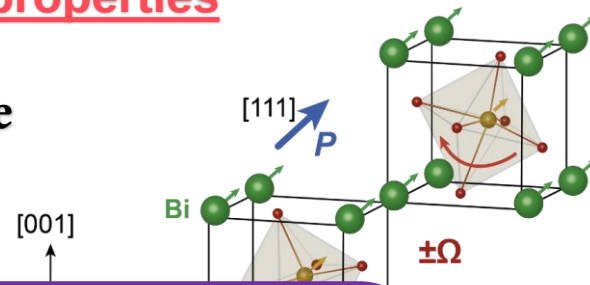
Bulk properties

1. Perovskite structure

Bi Cations

Fe Cations

O Anions



Objective:

Investigation of different spin
cycloids

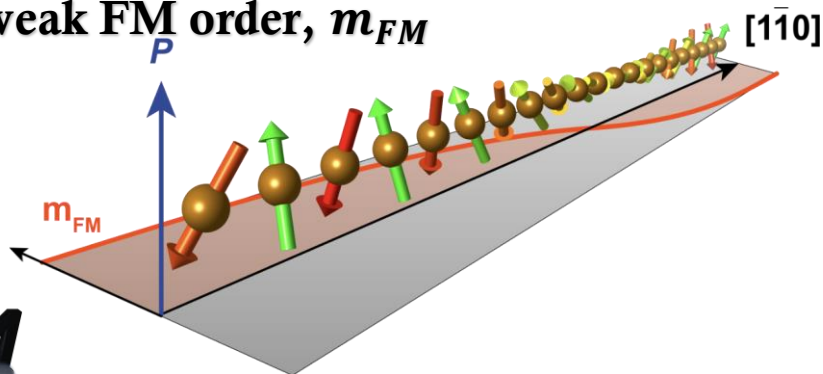
βc

2. Fe-

Pol

3. Oxygen tilt in antiphase $\pm\Omega$

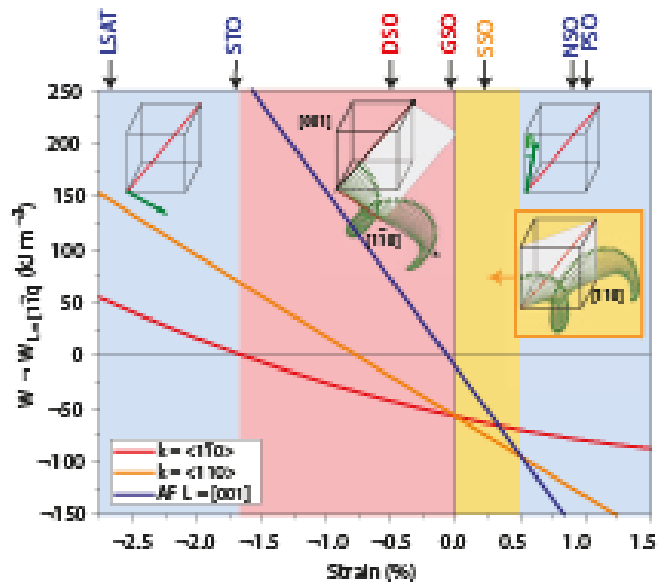
4. Local antiferromagnetic order with long-ranged spin cycloid, $\lambda \approx 62$ nm + weak FM order, m_{FM}



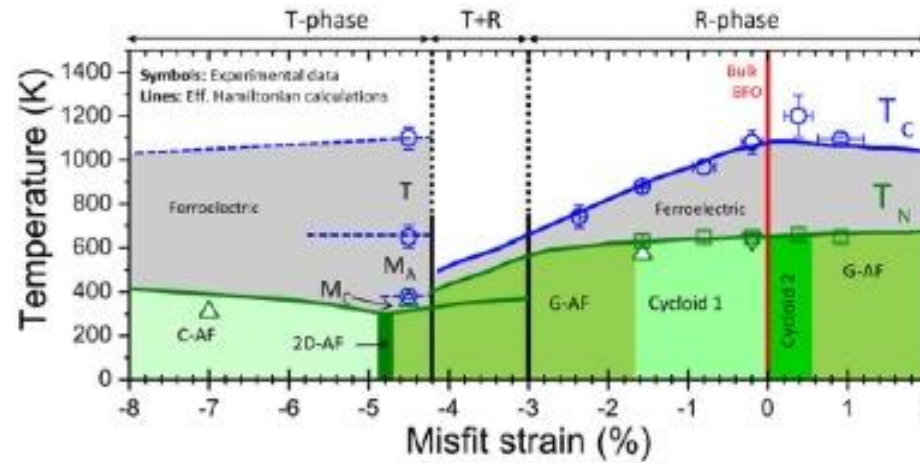
Magnetic properties and orders in BFO

Computed phase diagram from effective Hamiltonian

Spin spirals from experiments



D. Sando, *et al.*, Nature Mater. **12**, 641-646 (2013)



D.Sando. *et al.*, J. Phys.: Condens. Matter **26**,
473201 (2014)

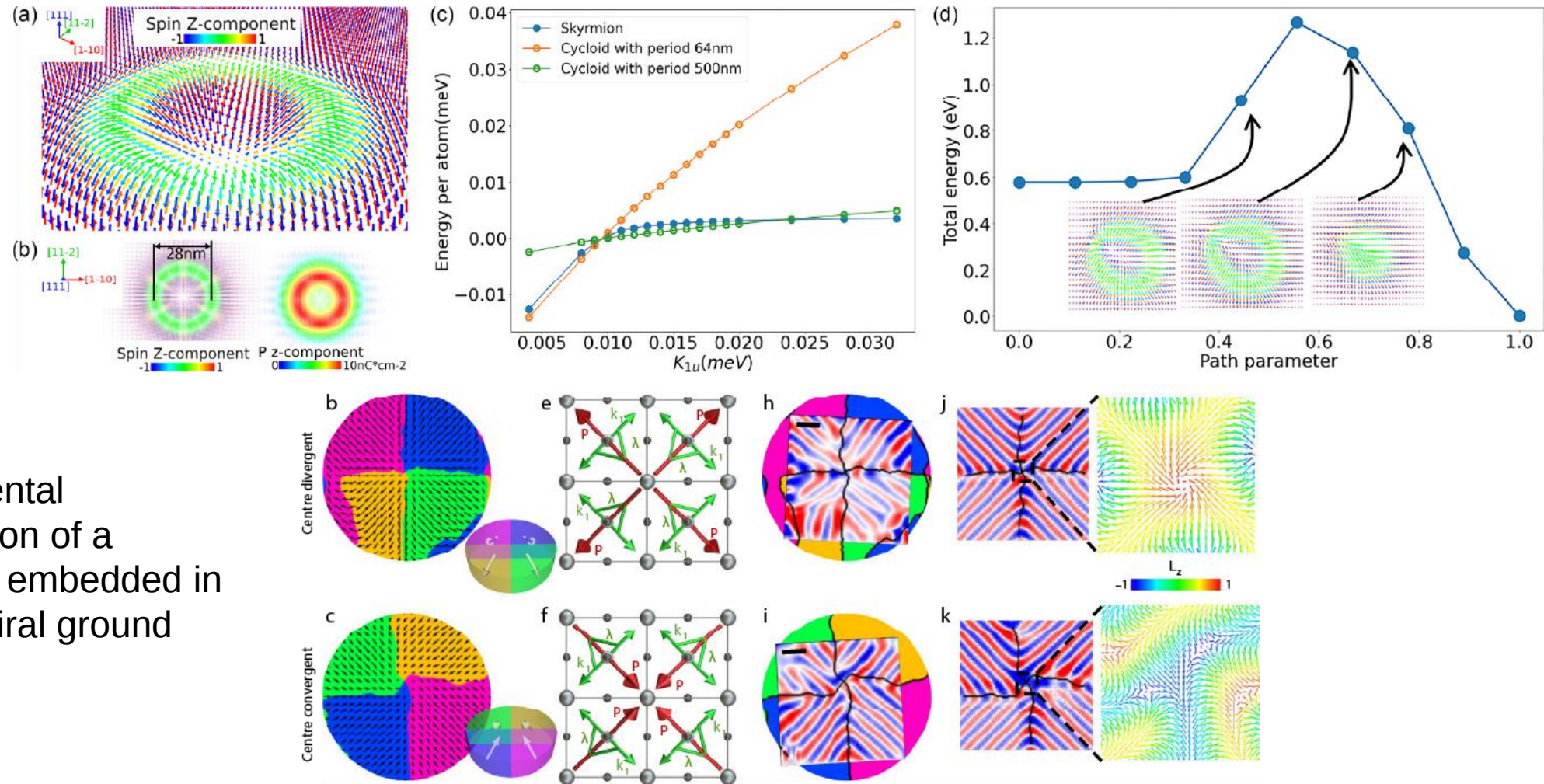
Experimental based Hamiltonian

Heisenberg Hamiltonian
obtained from Neutron
diffraction
Easy axis anisotropy

C. Weingart, *et al.*, Phys. Rev. B **86**, 094413 (2012)
R. Fishman, Phys. Rev. B **87**, 134416 (2013)

Topological textures in BFO

Stable Skyrmion in BFO with imposed out-of-plane anisotropy

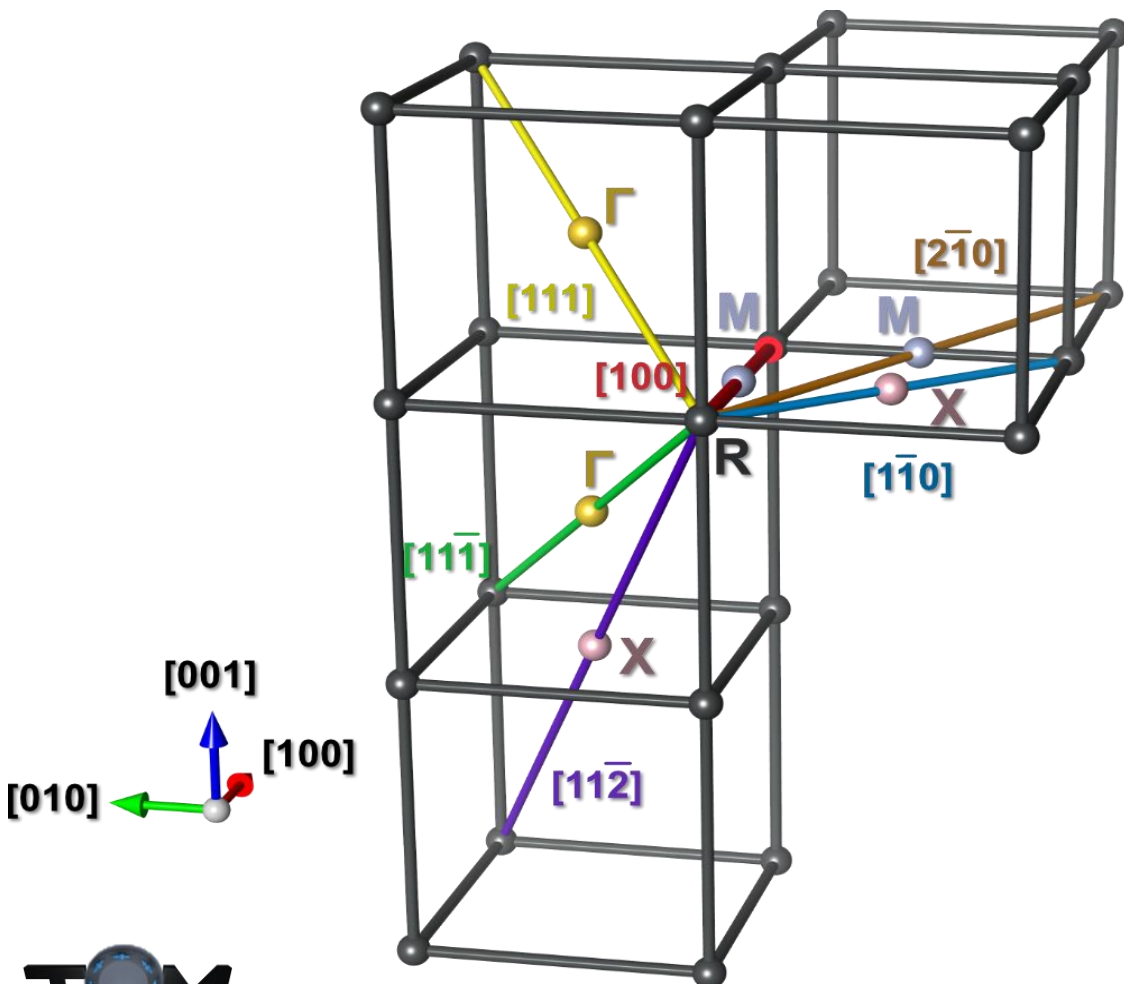


Experimental observation of a skyrmion embedded in a spin spiral ground state

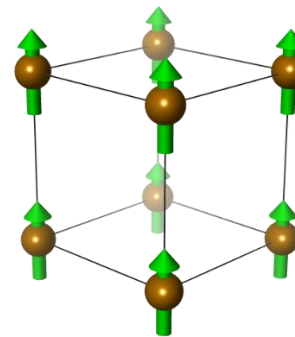
Z. Li, *et al.*, Phys. Rev. Research **5**, 043109 (2023)
A. Chaudron *et al.*, Nat. Matter **23**, 905-911 (2024)

Brillouin zone and high symmetry points of BFO

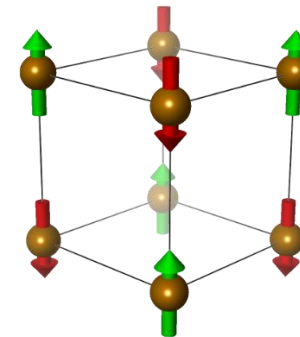
Detailed path of the q-vector



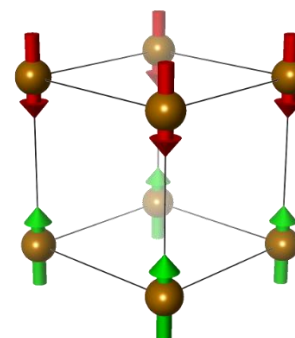
Γ
FM



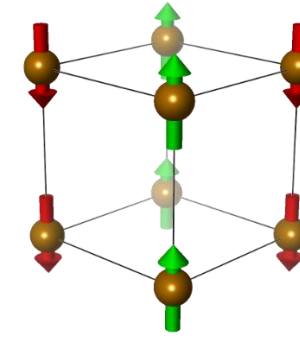
R
G-type AFM



X
A-type AFM

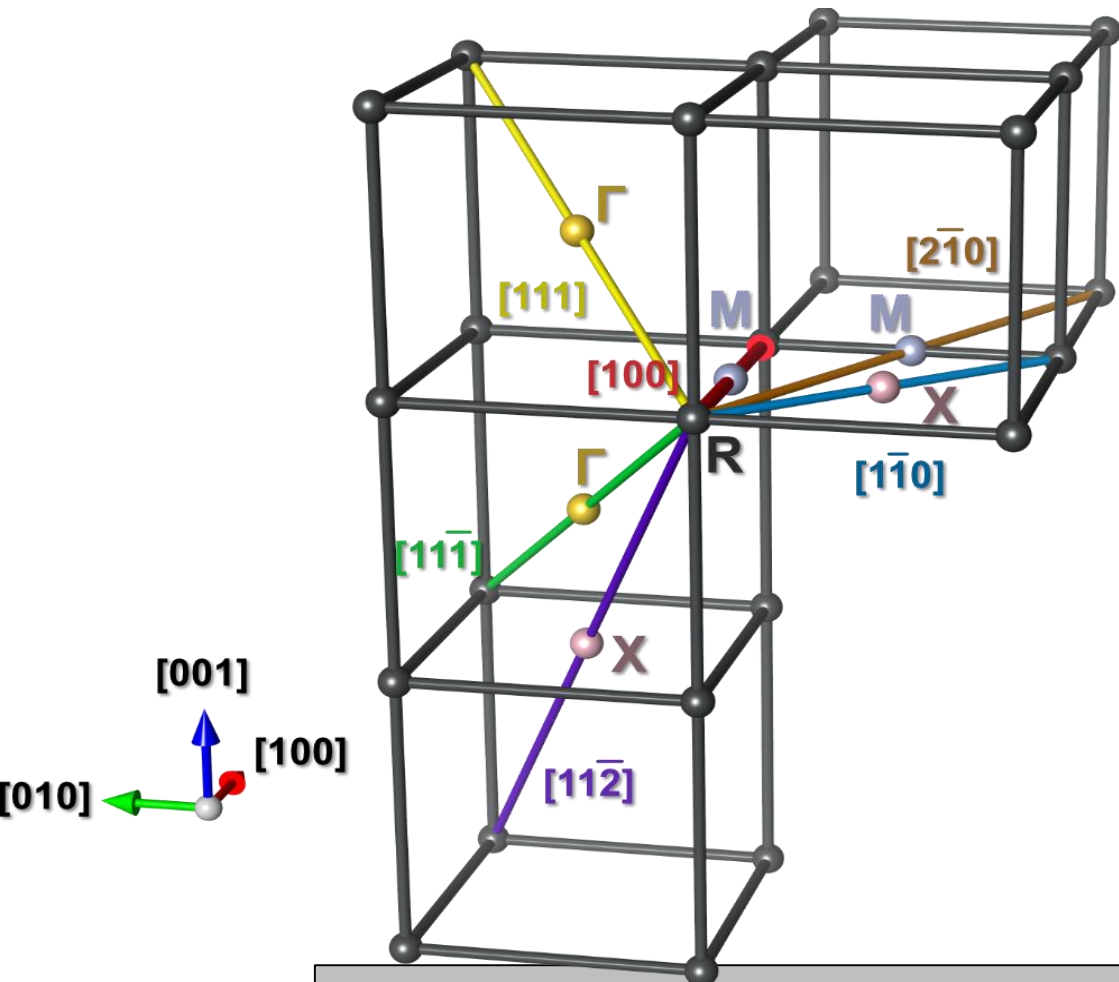


M
C-type AFM



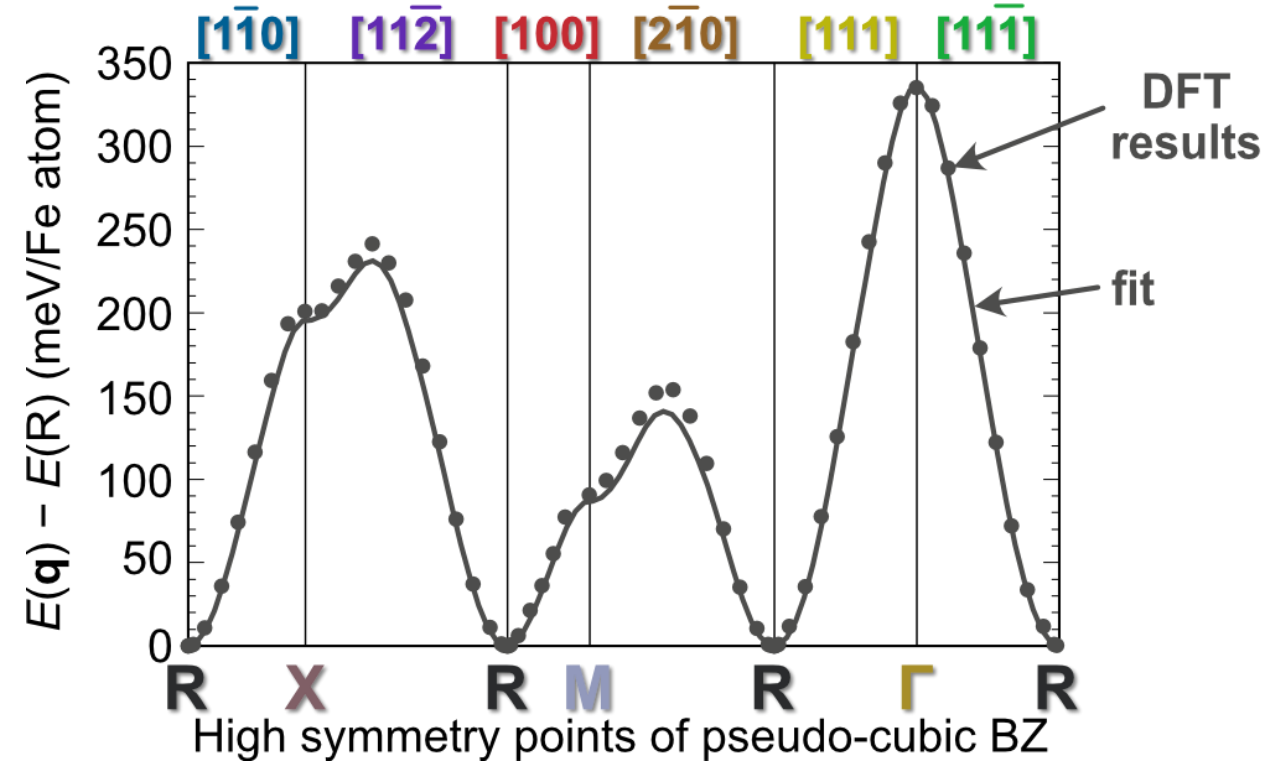
Spin spiral energy calculation in BFO

Detailed path of the q-vector



B. Xu, *et al.*, Phys. Rev. B **103**, 214423 (2023)
S. Meyer, *et al.*, Phys. Rev. B **108**, 024403 (2023)

Directions within the pseudo-cubic BZ

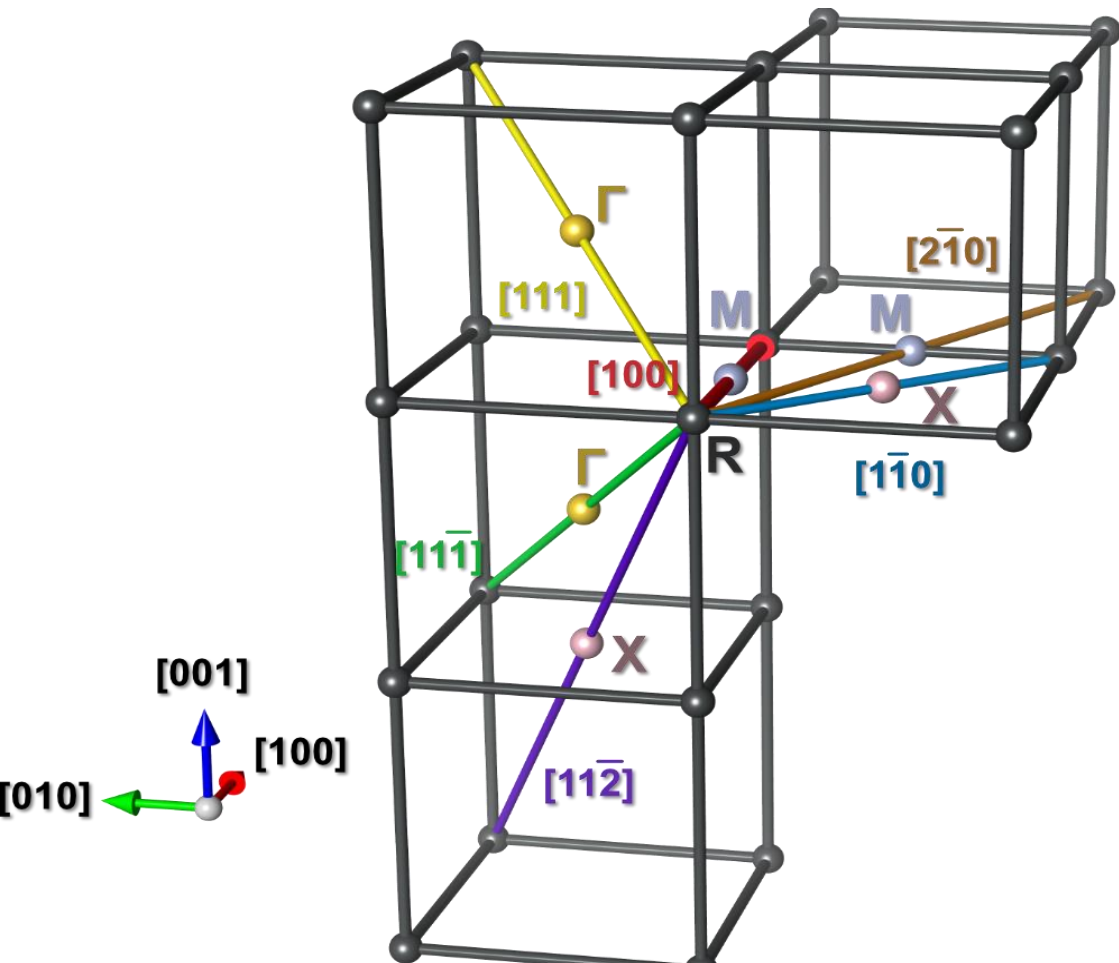


Conclusions

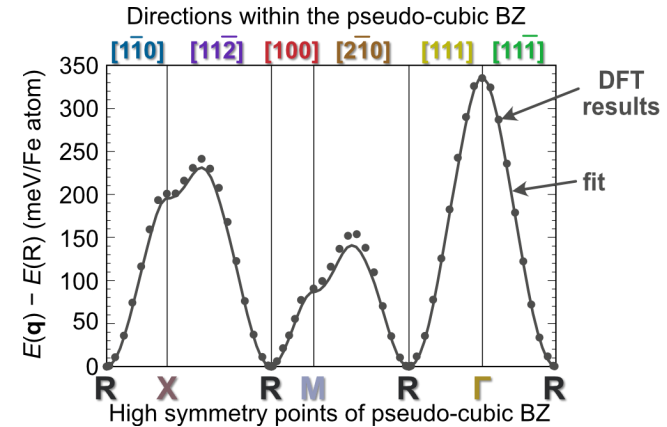
Dominant NN exchange J_1
Exchange fully isotropic in 3D BZ

Spin spiral energy calculation in BFO

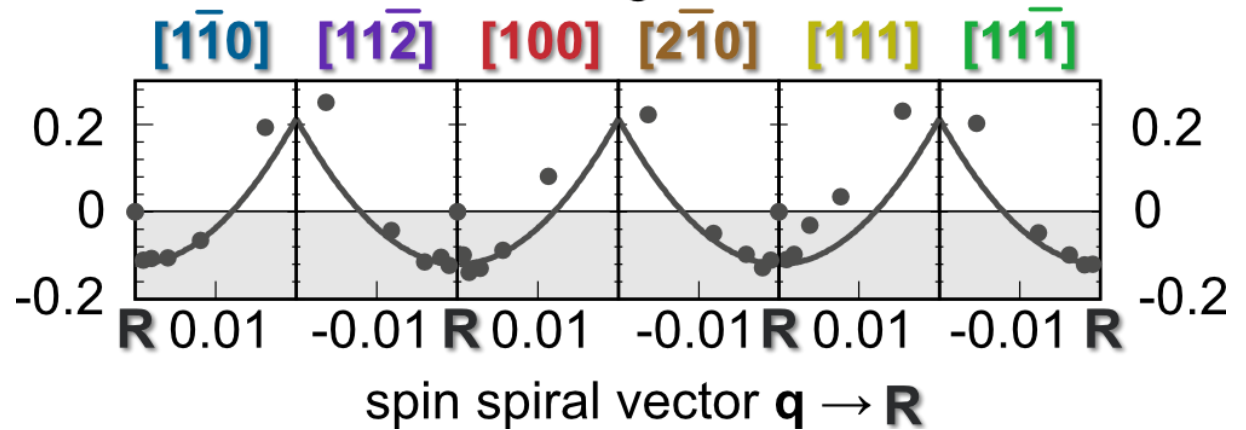
Detailed path of the q-vector



B. Xu, *et al.*, Phys. Rev. B **103**, 214423 (2023)
S. Meyer, *et al.*, Phys. Rev. B **108**, 024403 (2023)



Zoom around ground state

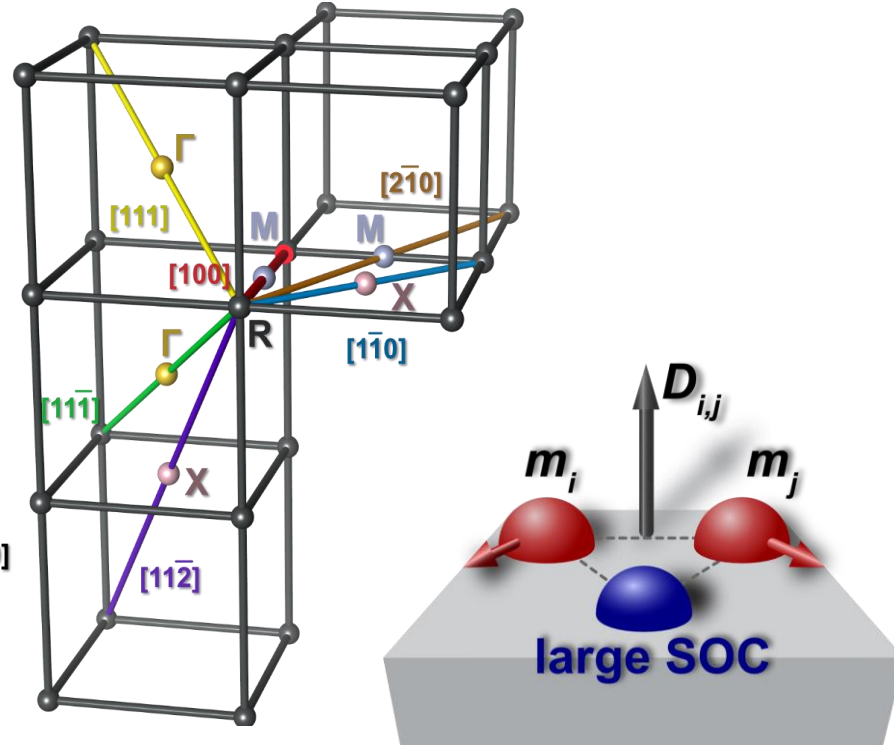


Conclusions

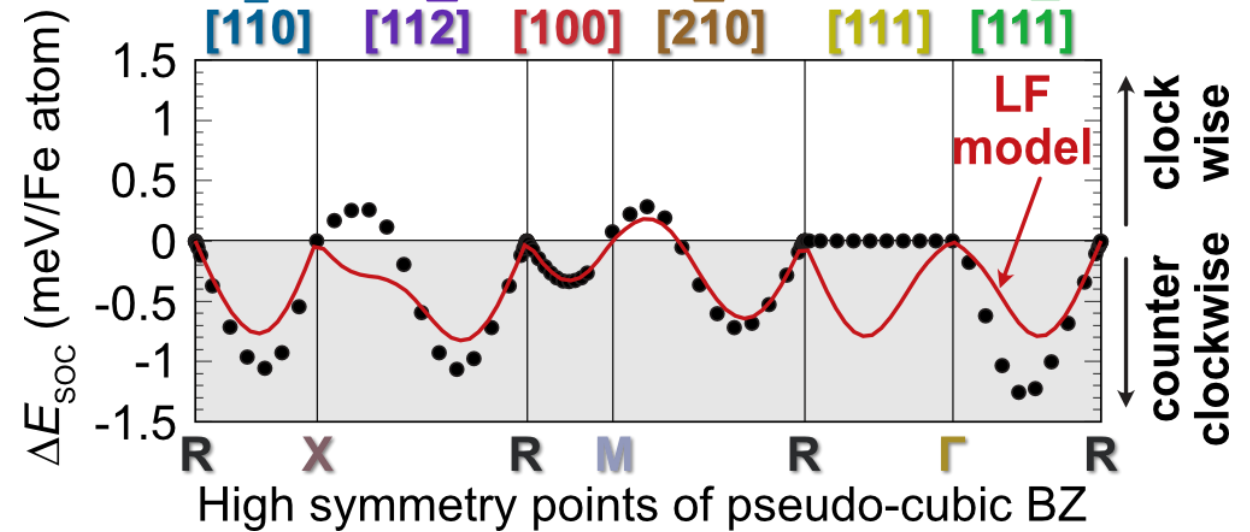
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DM interaction energy calculation in BFO

Detailed path of the q-vector



Directions within the pseudo-cubic BZ



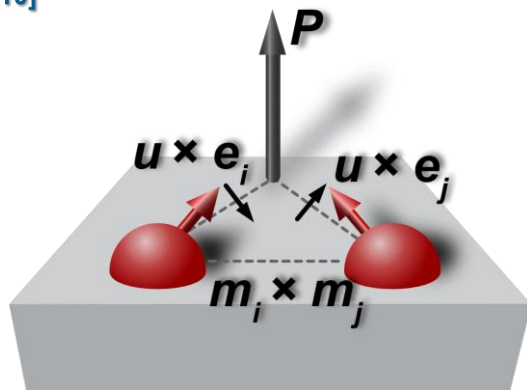
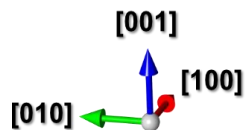
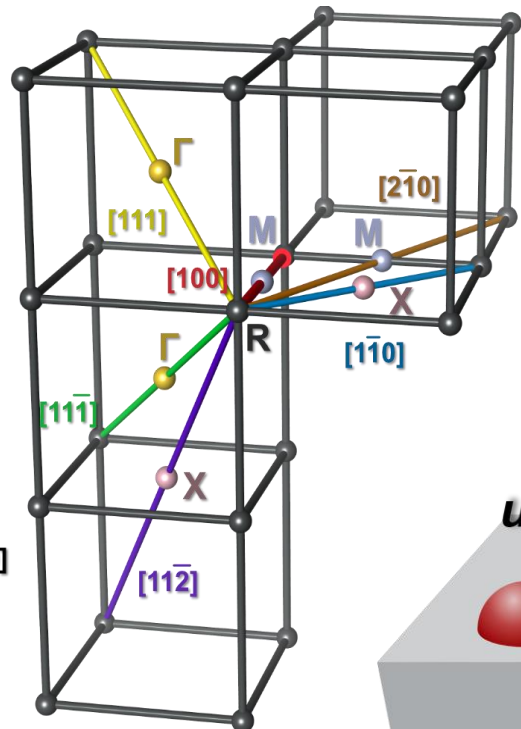
D_{ij} from Fe-O-Fe bonds

$$\mathcal{H}_{DMI}^{LF} = - \sum_{i,j} D_{ij} (\mathbf{m}_i \times \mathbf{m}_j)$$

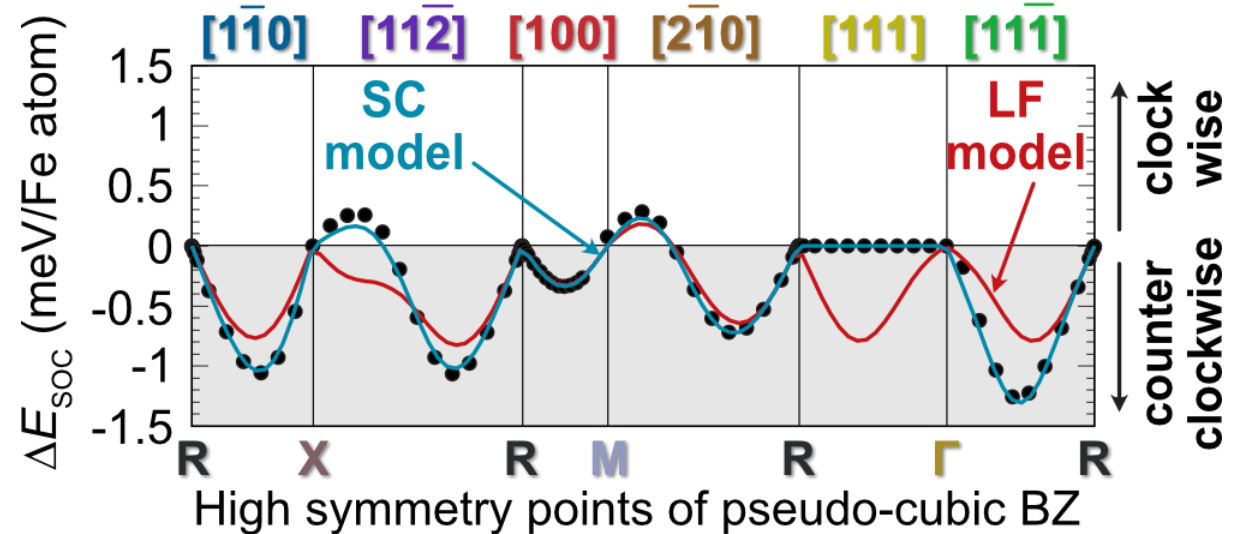
B. Xu, *et al.*, Phys. Rev. B **103**, 214423 (2023)
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DM interaction energy calculation in BFO

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Directions within the pseudo-cubic BZ

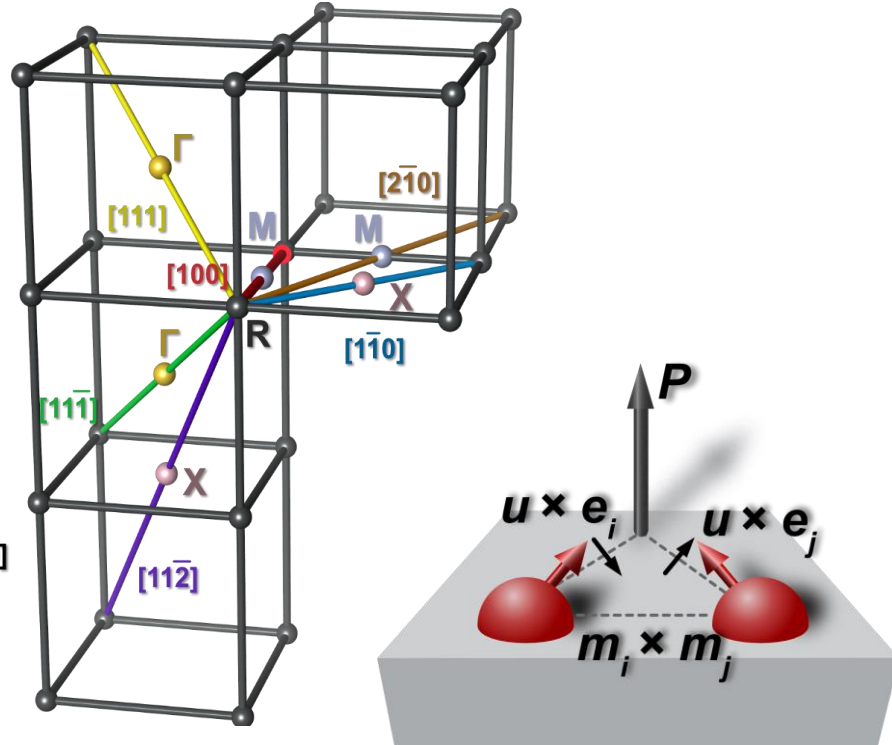


$$\mathcal{H}_{DMI}^{SC} = - \sum_{i,j} C_{ij} (\mathbf{u}_i \times \mathbf{e}_{ij}) (\mathbf{m}_i \times \mathbf{m}_j)$$

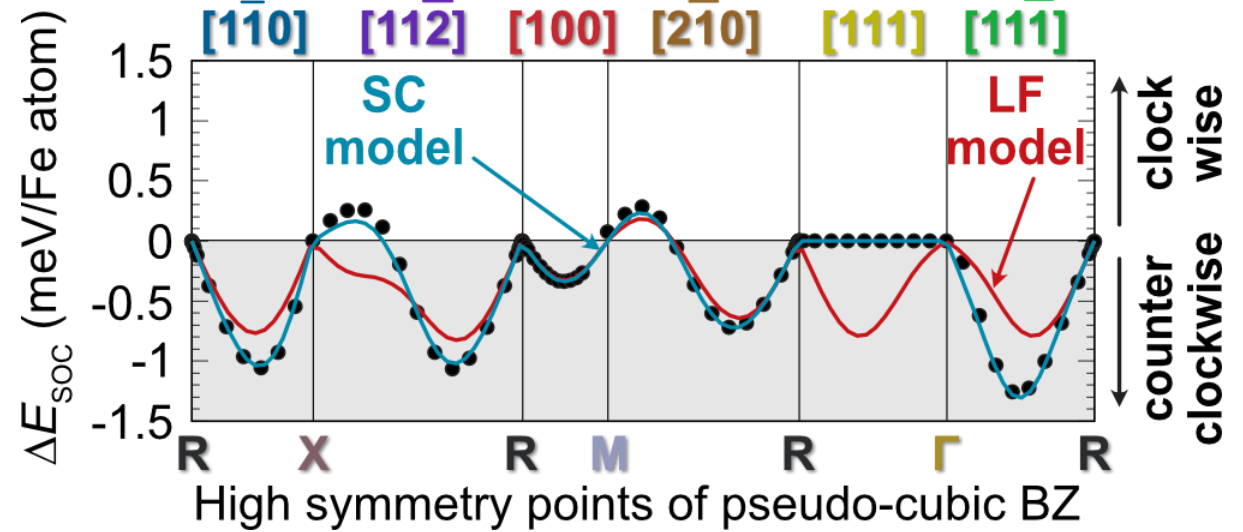
B. Xu, *et al.*, Phys. Rev. B **103**, 214423 (2023)
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Directions within the pseudo-cubic BZ



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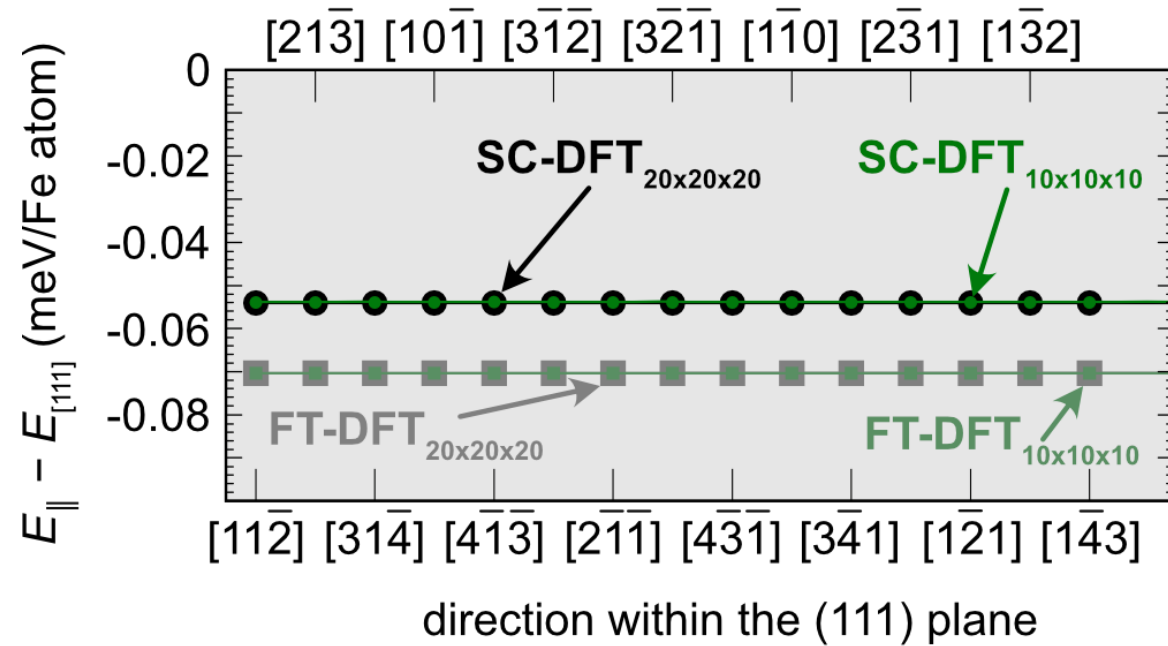
Conclusions

No SOC contribution in [111] direction
Disagreement with LF model
Spin-current model in agreement with data

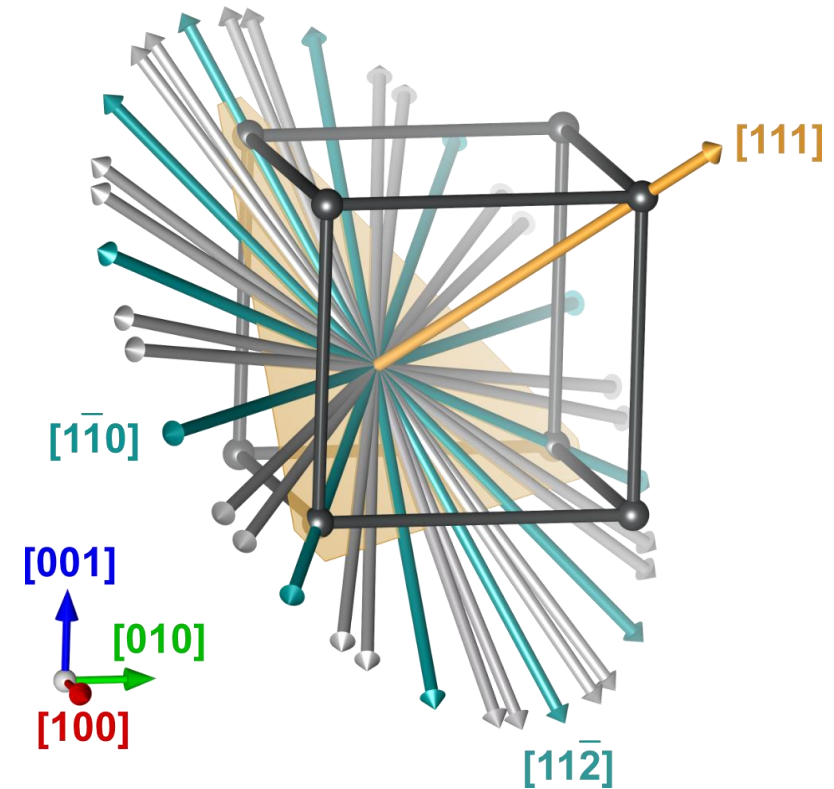
B. Xu, *et al.*, Phys. Rev. B **103**, 214423 (2023)
S. Meyer, *et al.*, Phys. Rev. B **108**, 024403 (2023)

Bertrand Dupé

Magnetocrystalline anisotropy in BFO

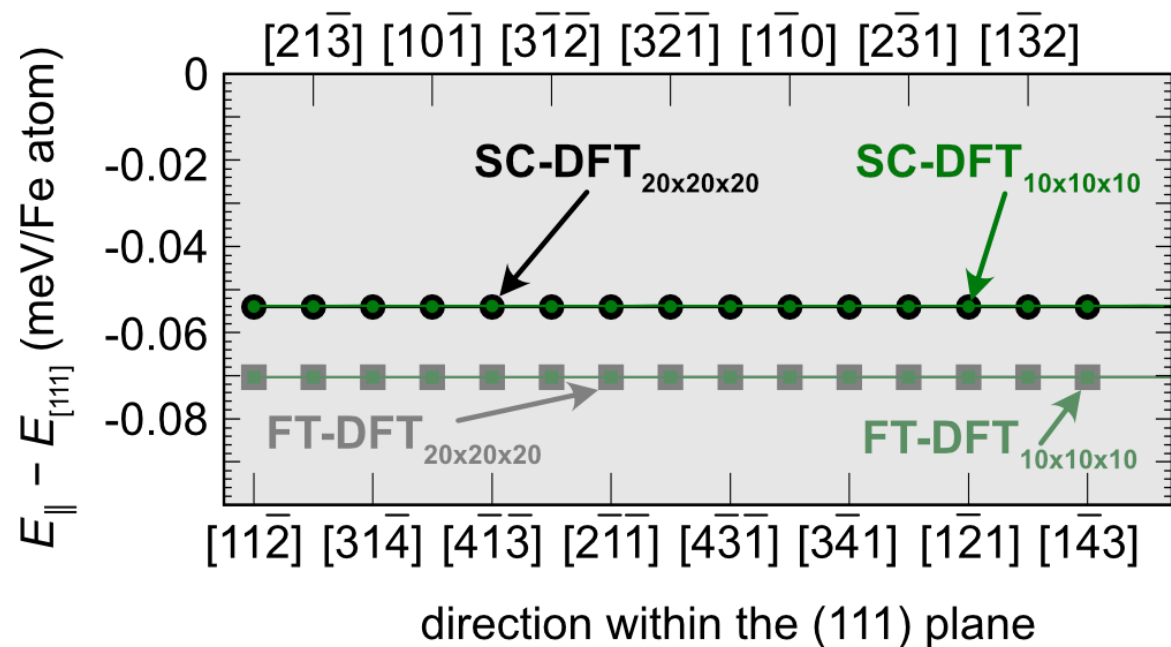


Calculations of preferred magnetization direction



B. Xu, *et al.*, Phys. Rev. B **103**, 214423 (2023)
S. Meyer, *et al.*, Phys. Rev. B **108**, 024403 (2023)

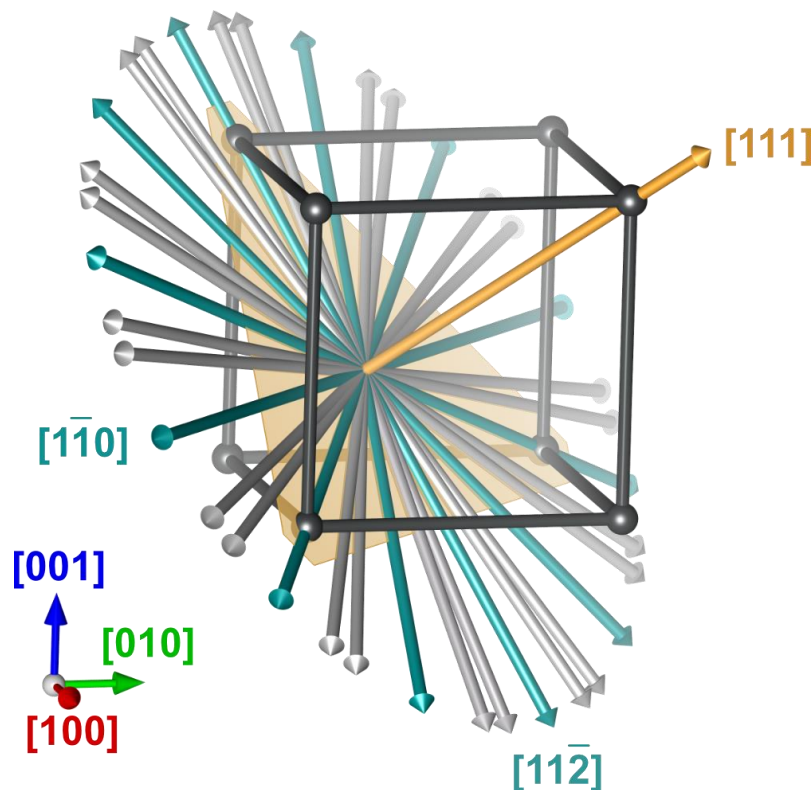
Magnetocrystalline anisotropy in BFO



Conclusions

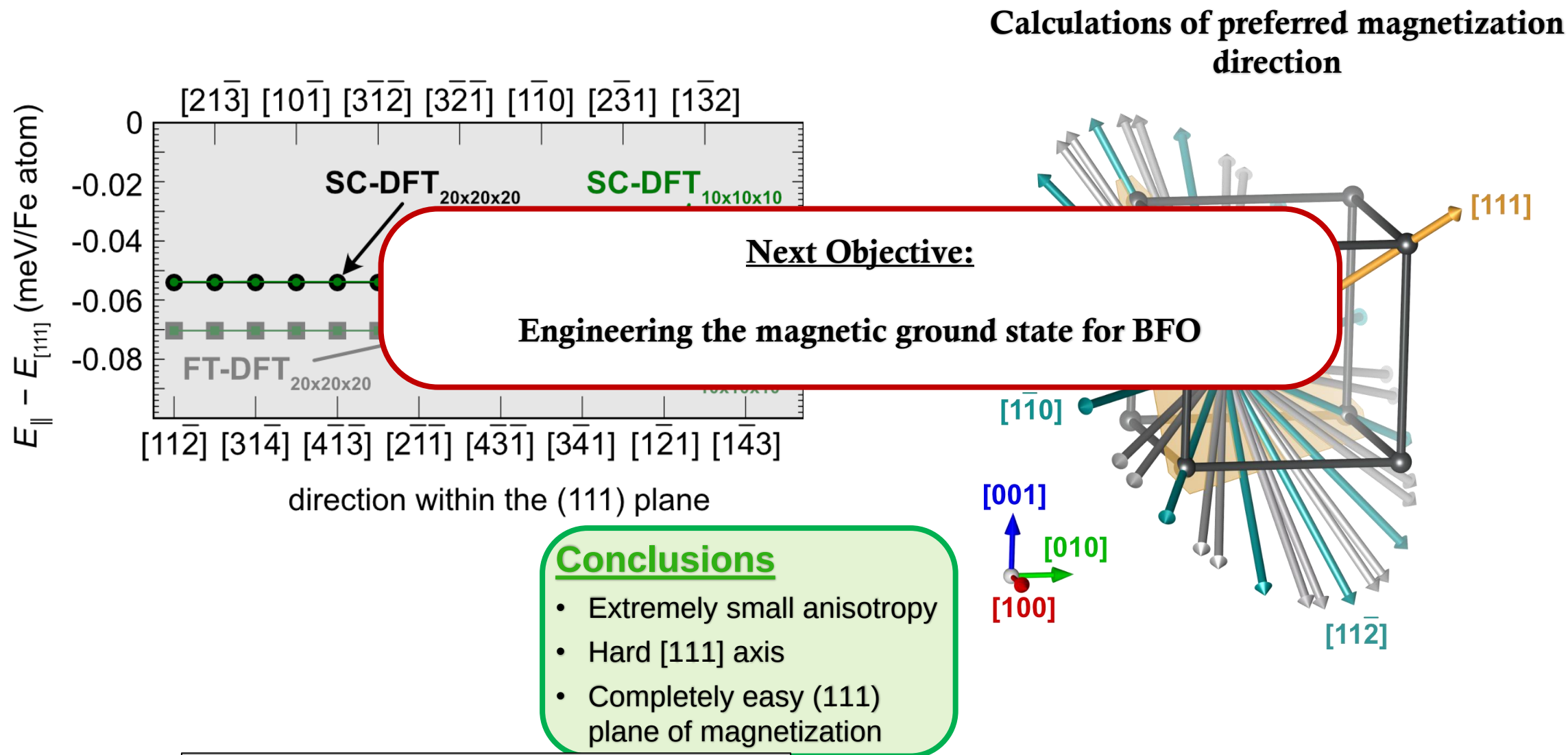
- Extremely small anisotropy
- Hard $[111]$ axis
- Completely easy (111) plane of magnetization

Calculations of preferred magnetization direction



B. Xu, *et al.*, Phys. Rev. B **103**, 214423 (2023)
S. Meyer, *et al.*, Phys. Rev. B **108**, 024403 (2023)

Magnetocrystalline anisotropy in BFO



B. Xu, *et al.*, Phys. Rev. B **103**, 214423 (2023)
S. Meyer, *et al.*, Phys. Rev. B **108**, 024403 (2023)

Order parameters of BiFeO₃

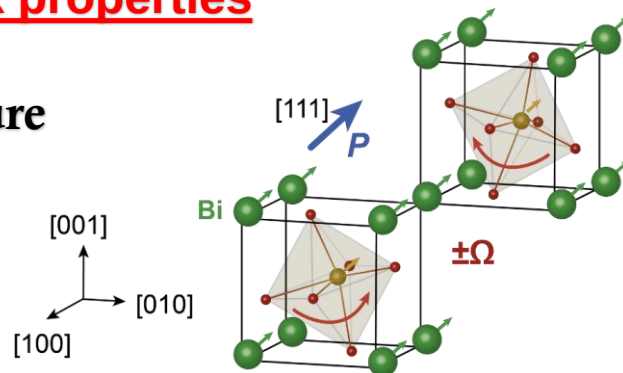
Bulk properties

1. Perovskite structure

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Fe Cations

O Anions



2. Fe-Bi displacement

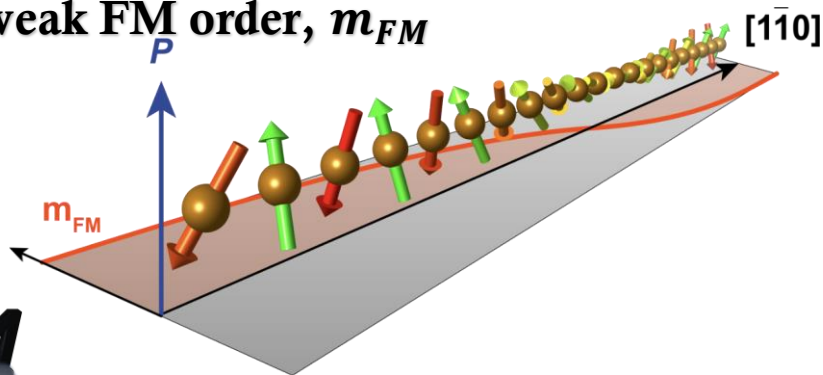
Polarization along [111] direction

$$P \approx \frac{90\mu\text{C}}{\text{cm}^2}, T_c = 1123 \text{ K}$$

$R3c$

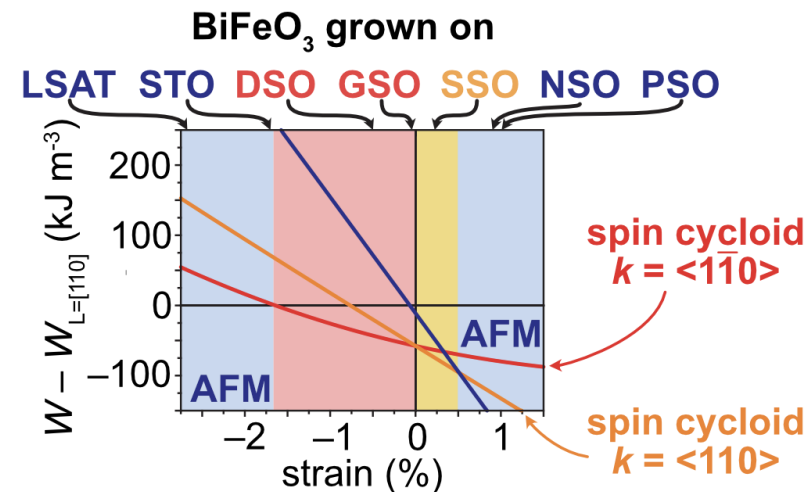
3. Oxygen tilt in antiphase $\pm\Omega$

4. Local antiferromagnetic order with long-ranged spin cycloid, $\lambda \approx 62 \text{ nm}$ + weak FM order, m_{FM}



Strain properties

1. Precise modification of material's lattice parameters
2. Modification of multiferroic properties:



D. Sando, et al., Nature Mater. 12, 641 (2013)

3. Collinear AFM state for strain $\varepsilon < -0.5\%$ and $\varepsilon > +0.5\%$
 4. Large sensitivity of magnetic ground states
- Different results for equivalent substrates

Order parameters of BiFeO₃

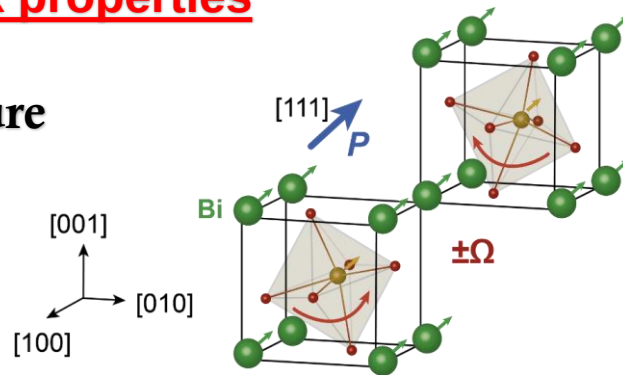
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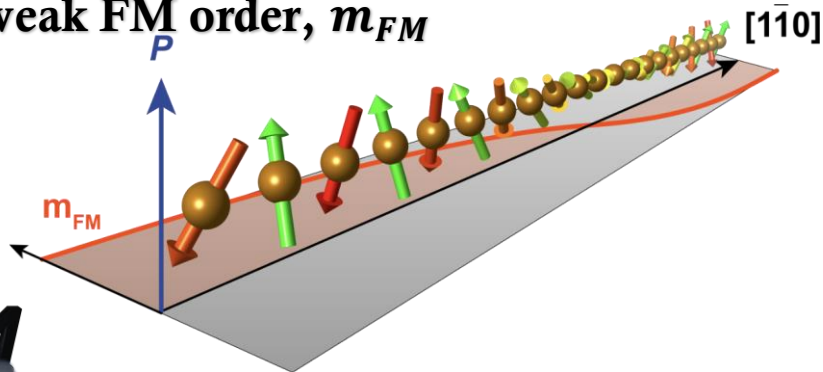
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$R\bar{3}c$

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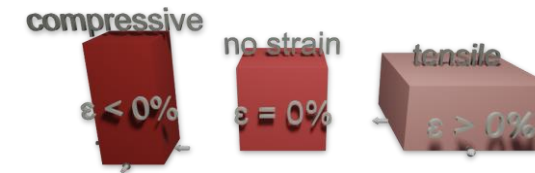
Strain properties

1. Precise modification of material's lattice parameters
2. Modification of multiferroic properties:

BiFeO₃ grown on

Objective:

Investigation of magnetic energies with epitaxial strain



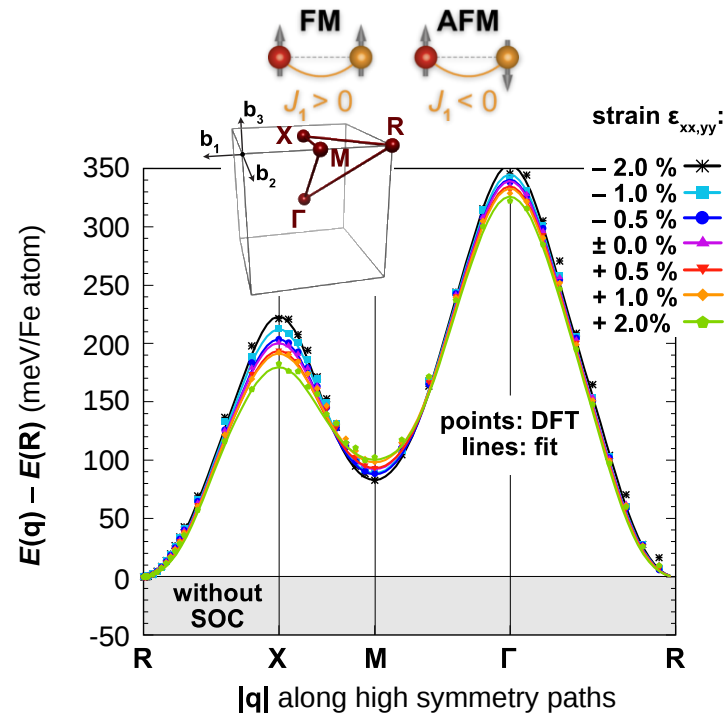
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4. Large sensitivity of magnetic ground states
Different results for equivalent substrates

Evolution of magnetic interactions with strain

S. Meyer, *et al.*, Phys. Rev. B **109**, 184431

(2024) **Exchange interaction**

$$\mathcal{H}_{ex} = - \sum_{i,j} J_{ij} (\mathbf{m}_i \cdot \mathbf{m}_j)$$



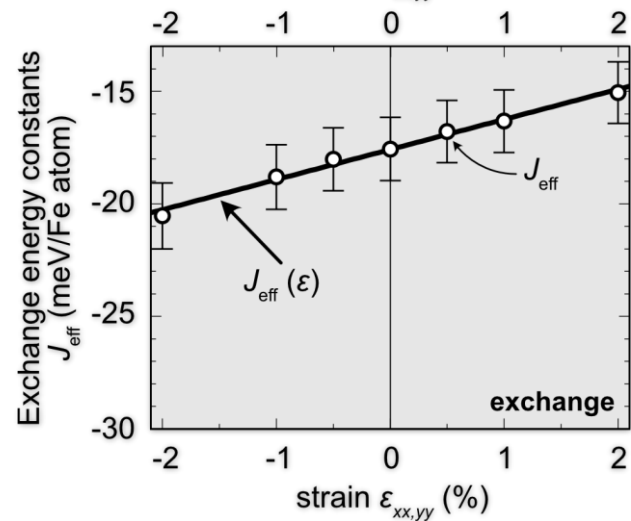
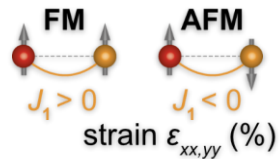
**Linearly decreasing
strength**

Evolution of magnetic interactions with strain

S. Meyer, *et al.*, Phys. Rev. B **109**, 184431

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**Linearly decreasing
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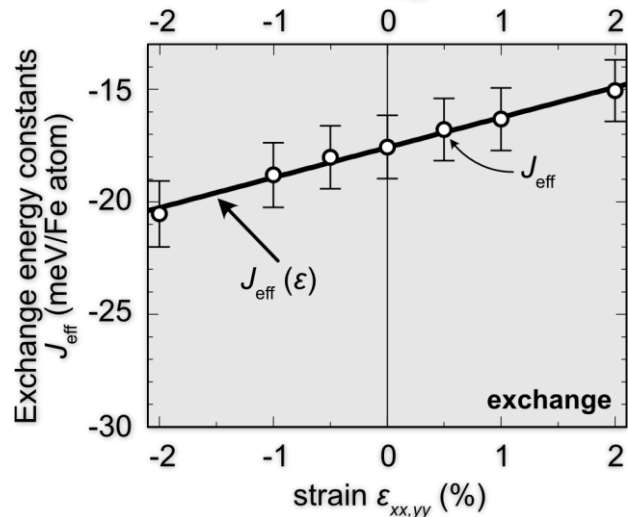
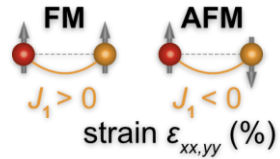
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S. Meyer, *et al.*, Phys. Rev. B **109**, 184431

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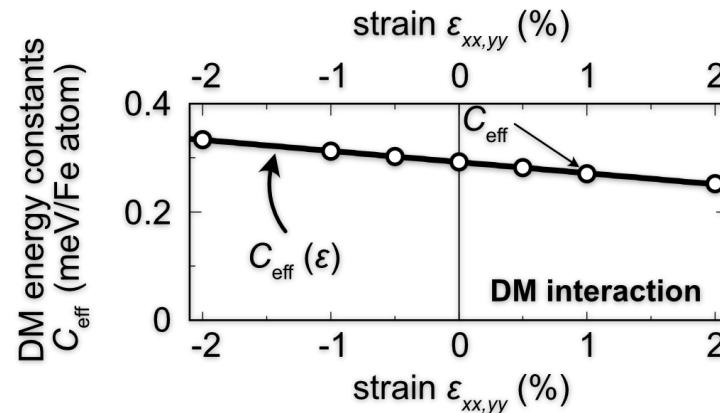
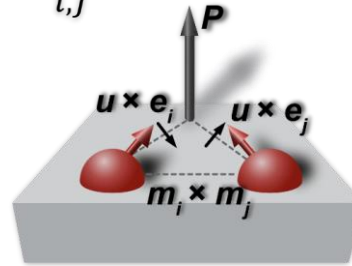


Linearly decreasing
strength



Dzyaloshinskii-Moriya interaction

$$\mathcal{H}_{DMI}^{SC} = - \sum_{i,j} C_{ij} (\mathbf{u} \times \mathbf{e}_{ij}) (\mathbf{m}_i \times \mathbf{m}_j)$$



Linearly decreasing
strength

Bertrand Dupé

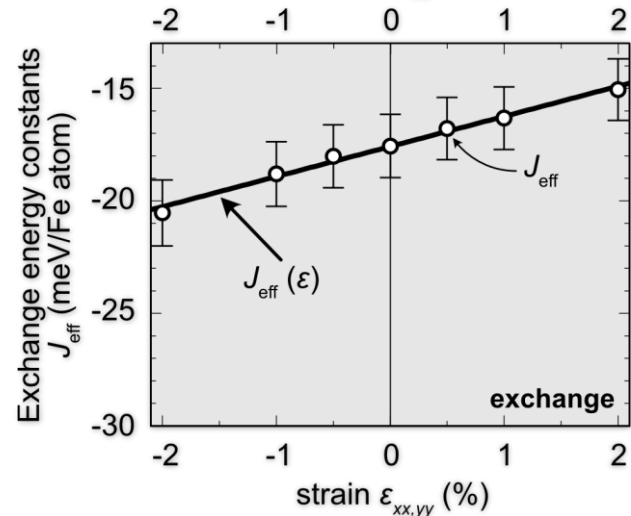
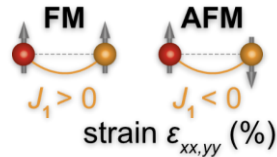
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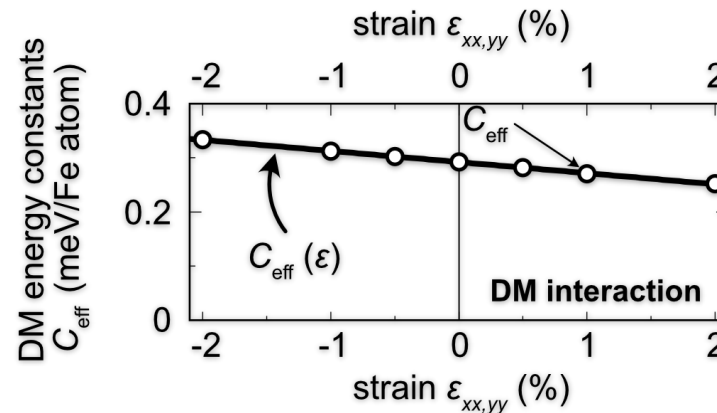
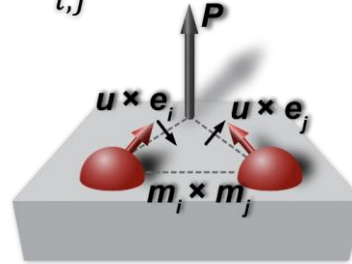


Linearly decreasing strength



Dzyaloshinskii-Moriya interaction

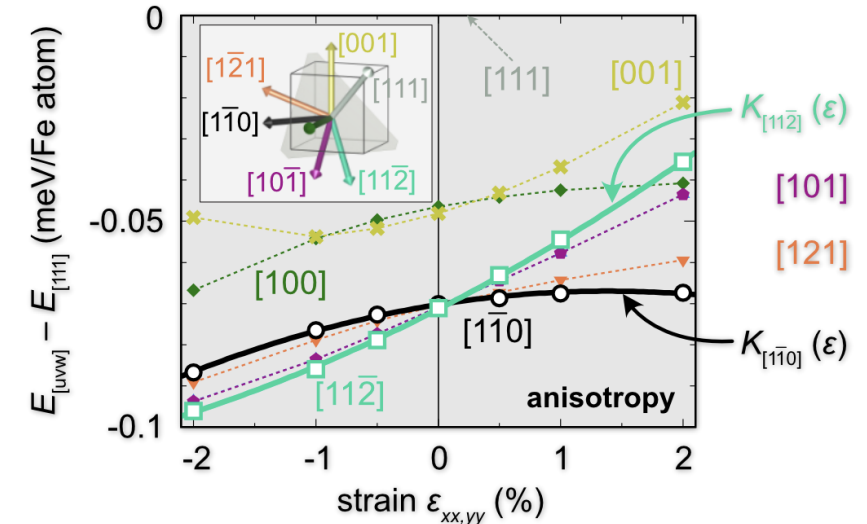
$$\mathcal{H}_{DMI}^{SC} = - \sum_{i,j} C_{ij} (\mathbf{u} \times \mathbf{e}_{ij}) (\mathbf{m}_i \times \mathbf{m}_j)$$



Linearly decreasing strength

Magnetocrystalline anisotropy

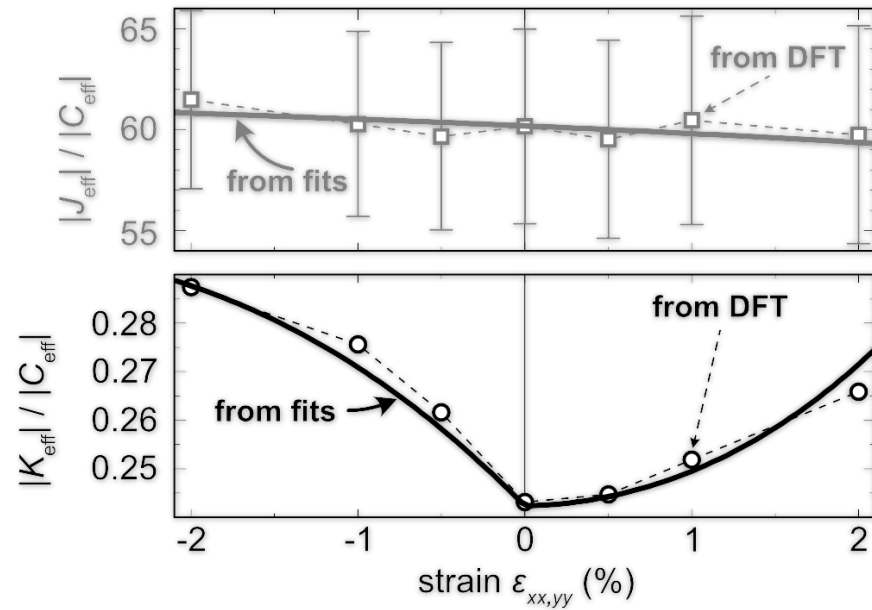
$$\mathcal{H}_{ani} = - \sum_i K m_{[111]}^2$$



$[11\bar{2}]$ easy axis for $\epsilon < 0$
 $[1\bar{1}0]$ easy axis for $\epsilon > 0$

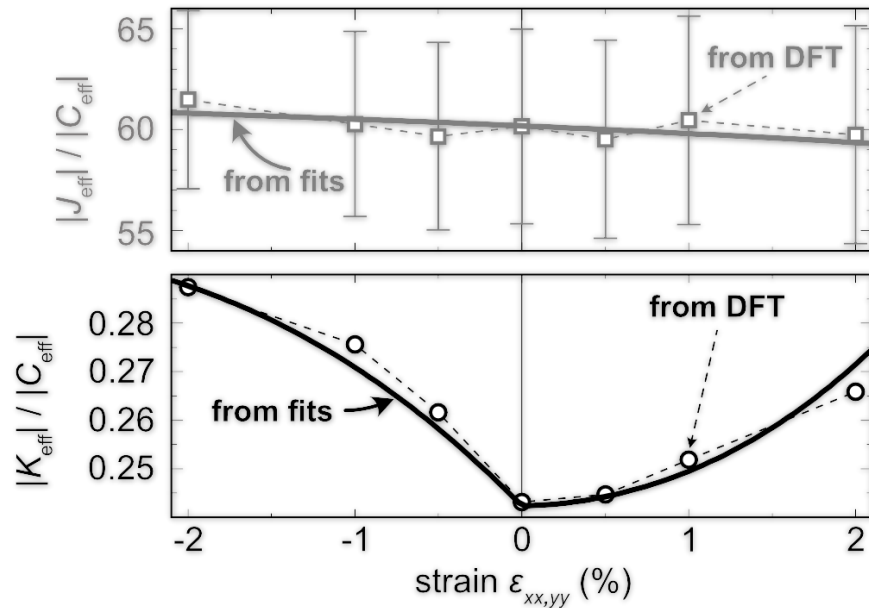
Evolution of the ratio between Exchange, anisotropy and DM

Ratios of magnetic interactions



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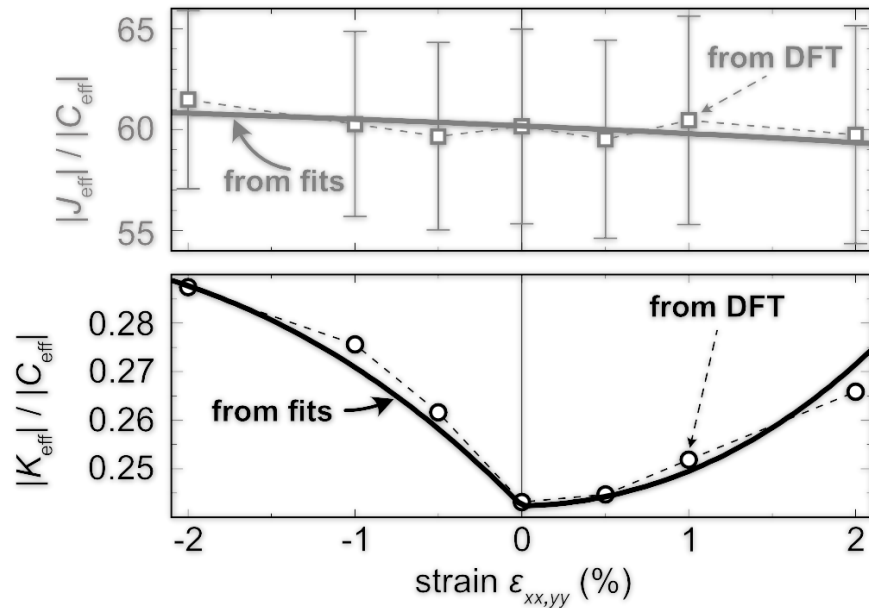


Conclusions

- $\varepsilon < 0$: Preferred collinear order from exchange and anisotropy
- $\varepsilon > 0$: Weaker collinear order from exchange vs. DMI, but Stronger due to MAE

Domain wall energy

Ratios of magnetic interactions



Conclusions

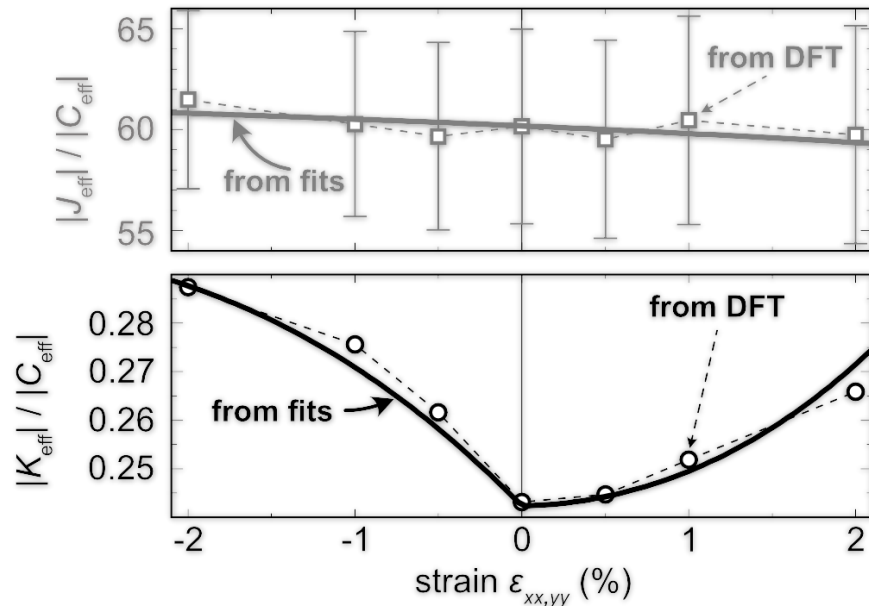
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Magnetic domain wall energies

$$E_{DW}(\varepsilon) = \frac{4}{a^2(\varepsilon)} \sqrt{J_{\text{eff}}(\varepsilon)K_{\text{eff}}(\varepsilon)} - \frac{2\pi}{a^2(\varepsilon)} C_{\text{eff}}(\varepsilon)$$

Domain wall energy

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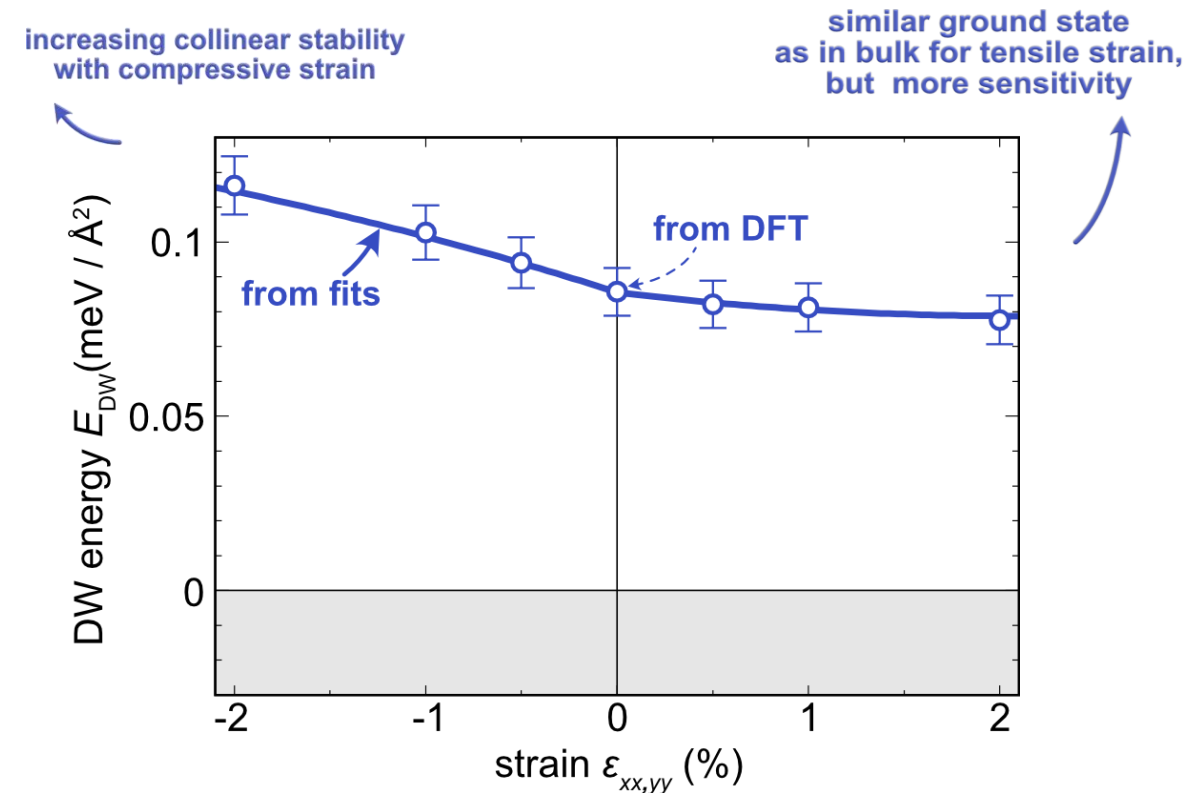


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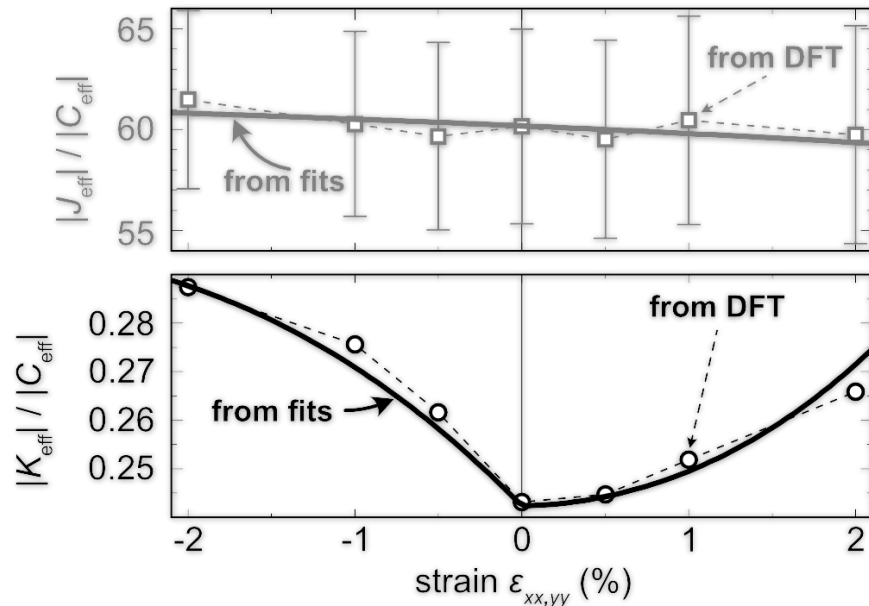
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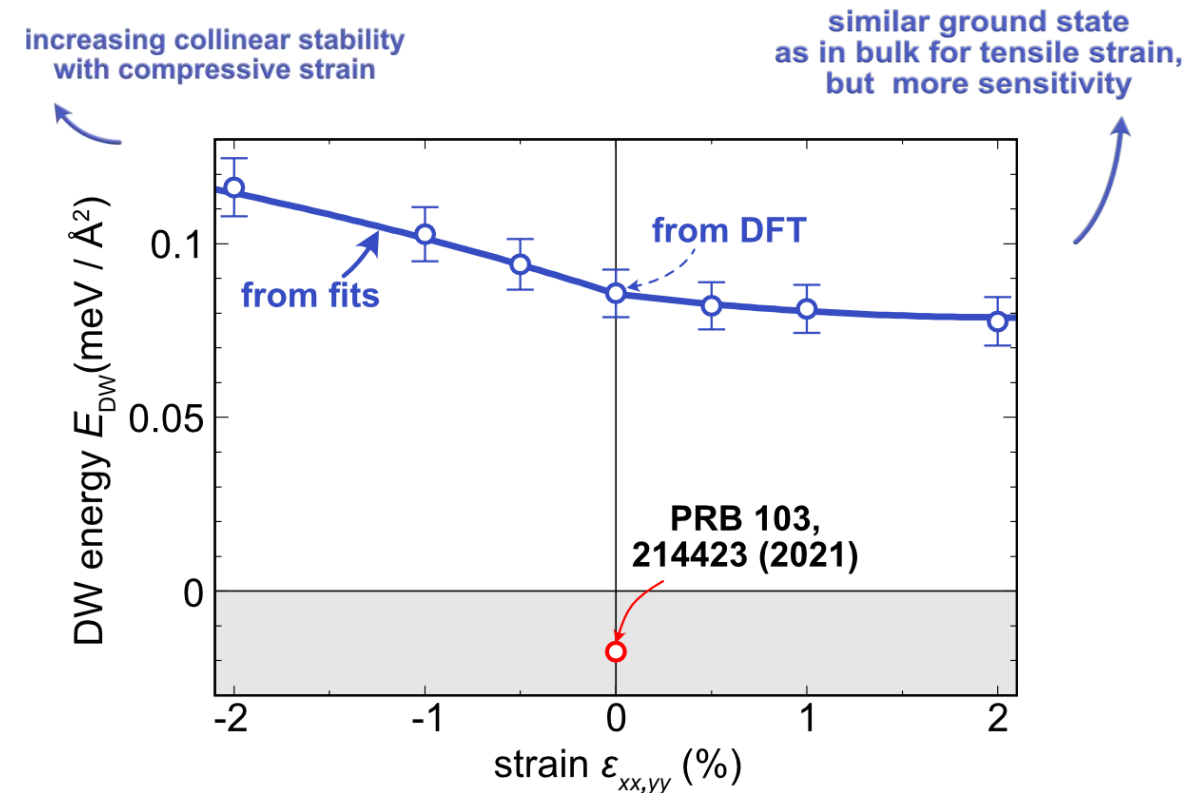


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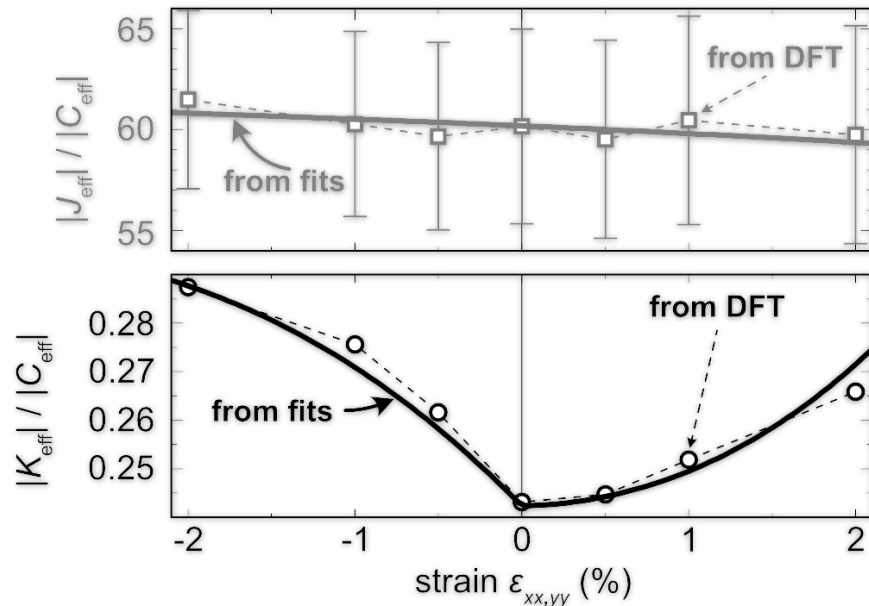
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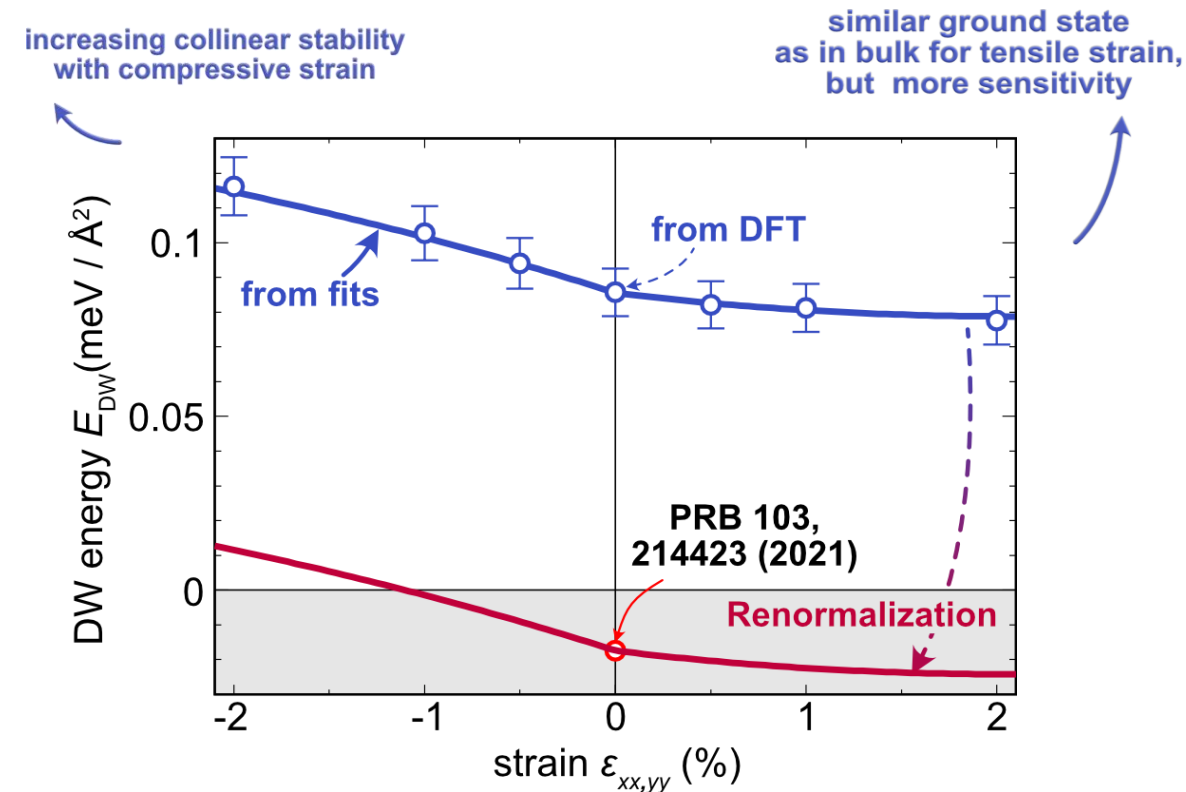


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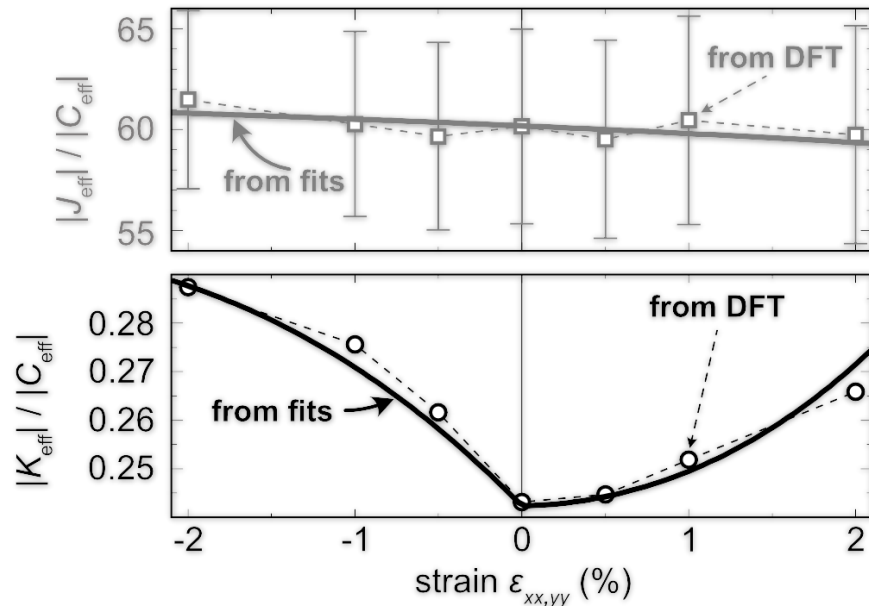
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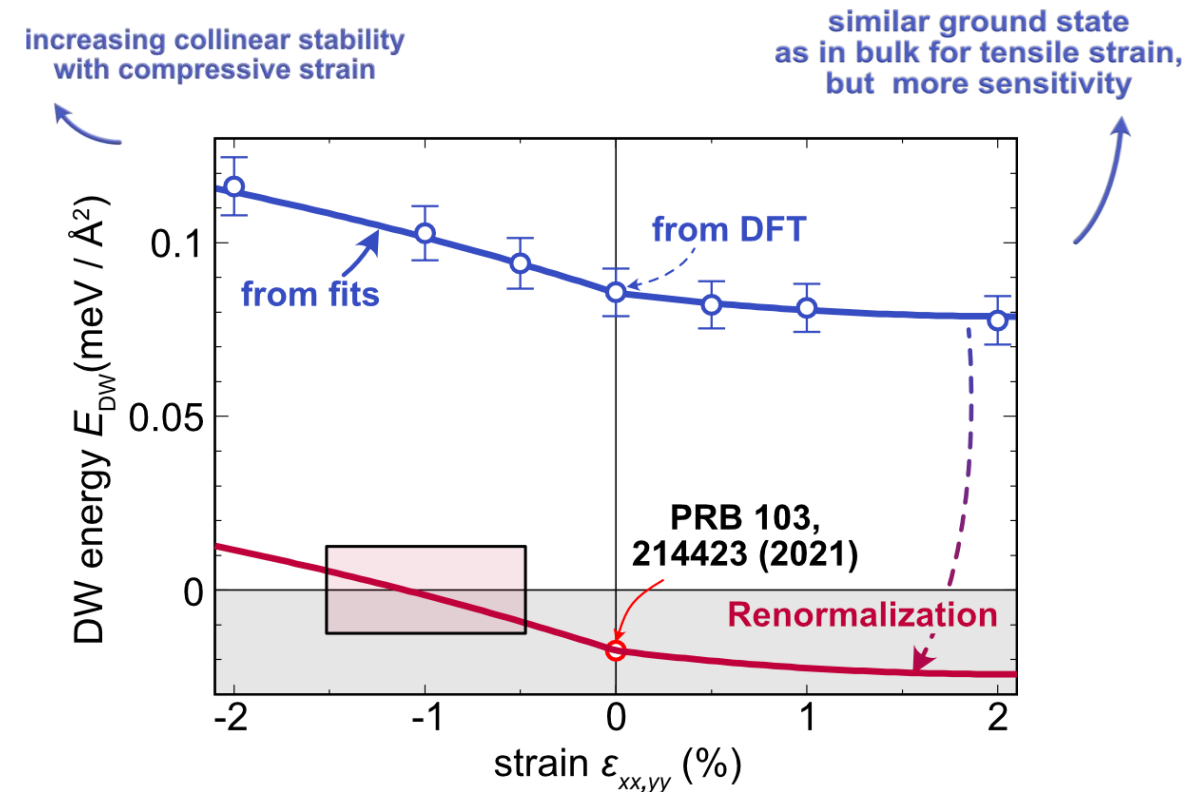


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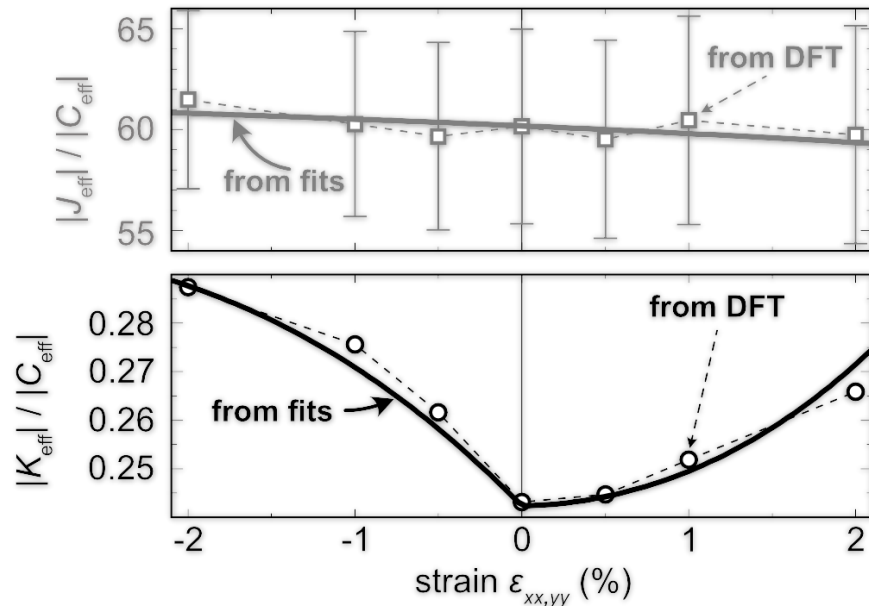
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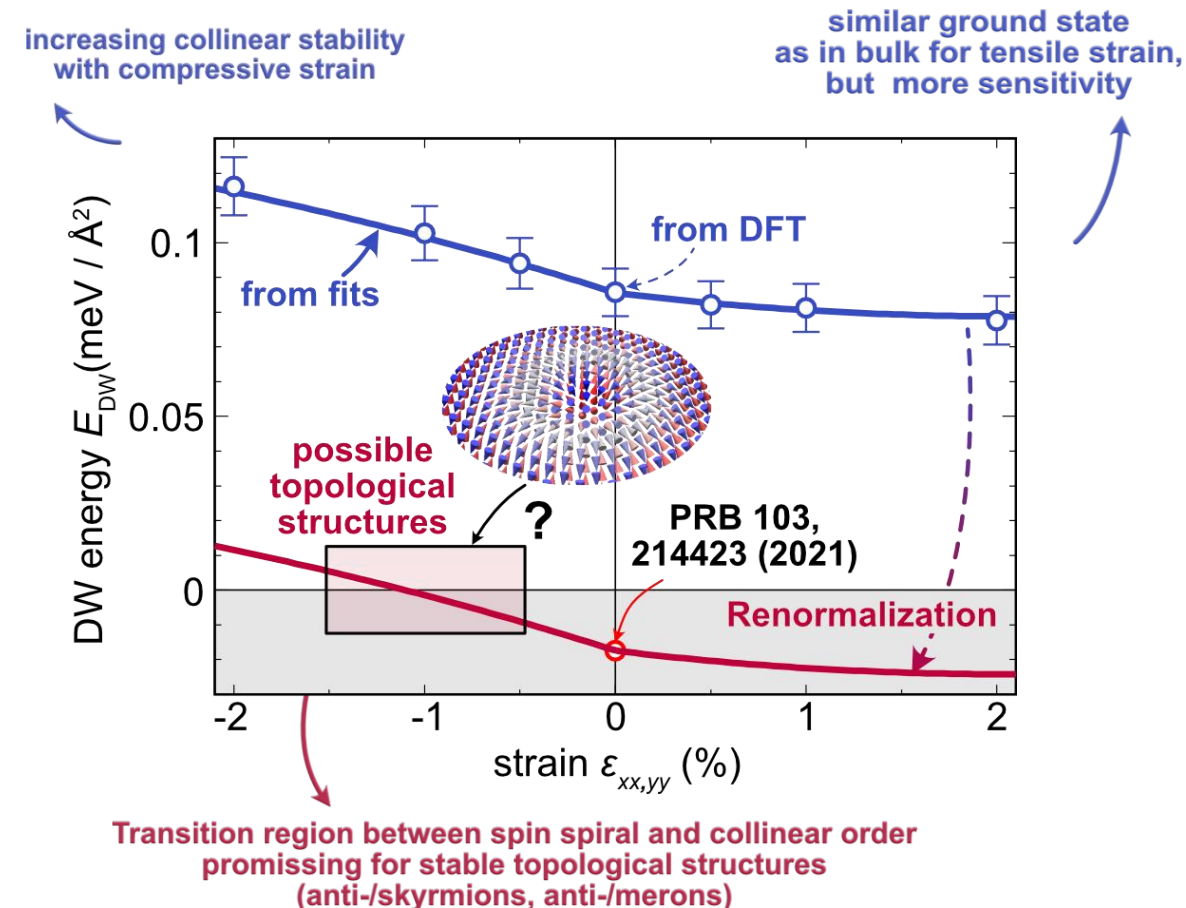


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Conclusion

- Topology is a very rich playground to do physics
- Multiferroics offer an excellent platform to stabilize ferroelectric and magnetic skyrmions
- Skyrmions can be functionalize for
 - neuromorphic computing,

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G. Yang, *et al.* Phys. Rev. B **93**, 224505 (2016)
- Check our code github.com/bertdupé/Matjes



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