

# Magnetism in Communications

## Spin waves for signal processing

Philipp Pirro

Fachbereich Physik and Landesforschungszentrum OPTIMAS,  
Rheinland Pfälzische Technische Universität Kaiserslautern-Landau,  
Kaiserslautern, Germany



CoSpIN



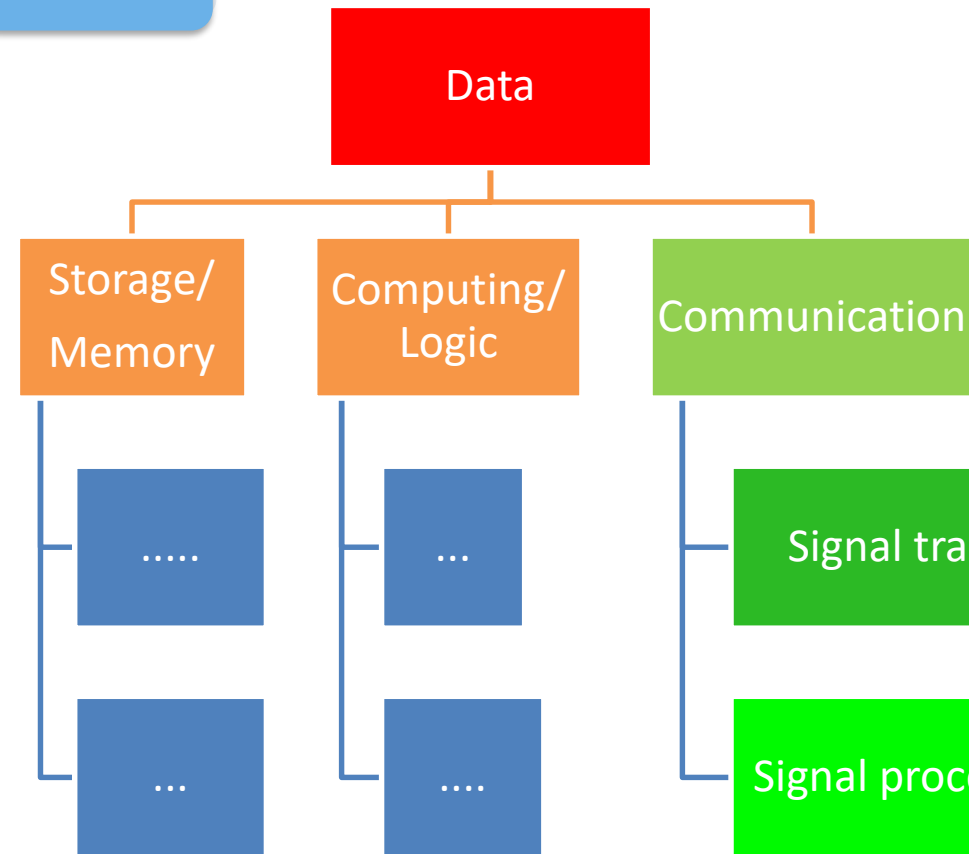
European Research Council  
Established by the European Commission



SPIN+X  
SFB/TRR 173  
Kaiserslautern • Mainz

DFG

- **Introduction: Data, Communication, Waves ...**
- **Spin waves for communication**
  - **Linear devices**
    - Filters, phase- and time delay units
  - **Nonlinear devices**
    - Limiter, signal-to-noise enhancement



## Physical/technical side

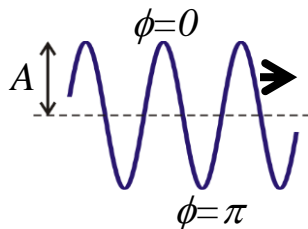
⇒ Communication: How to transfer data from A to B



⇒ Mainly *electromagnetic*: high speed + low attenuation of photons

⇒ Different *solid state* solutions

Communication requires **spatially moving dynamics** => **waves** (sound, light, spin ...)



$$\vec{A} = \vec{A}_0 \exp(\vec{k} \cdot \vec{x} - \omega t)$$

Wave equation

Interference phenomena

Linear dispersion relation => Group velocity

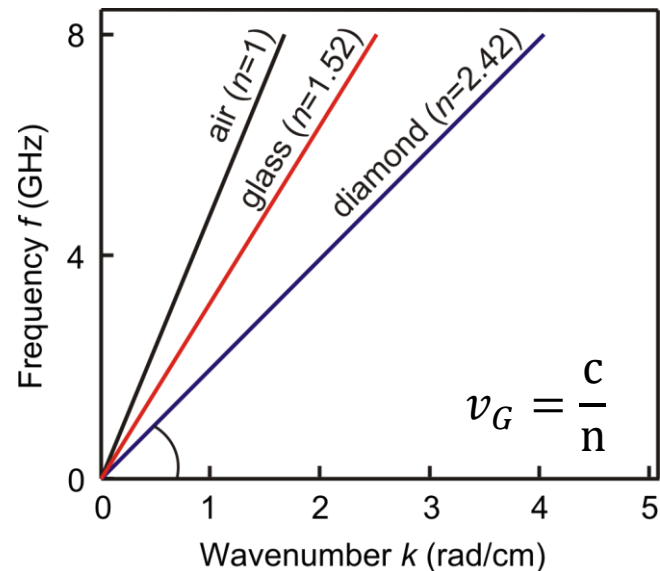
- is equal to phase velocity
- is independent of frequency/wave vector

**Speed** of the wave packet:

$$v_G = \frac{\partial \omega}{\partial t} = 2\pi \frac{\partial f}{\partial t}$$

**Group velocity defines speed of energy/information propagation.**

Example: Dispersion relation of light waves in different media



How “fast” is my communication?

1. **Delay time  $t$**  of the signal is given by the group velocity  $v_G$  of the wave.
2. **Data transfer rate  $R$**  of the communication is given by the **instantaneous bandwidth  $\Delta f$**  of the signal and the signal processing devices:

$$R \propto \Delta f$$

=> high bandwidth  $\Delta f$  needed for high speed data

# Electro-magnetic waves as signal carrier

## Visible light in optical fibers

$$\lambda_{em} = 1064 \text{ nm} \quad (f \approx 280 \text{ THz})$$



- High bandwidth
- Low attenuation

MAGNETISM?

## Radio-frequencies /Microwaves

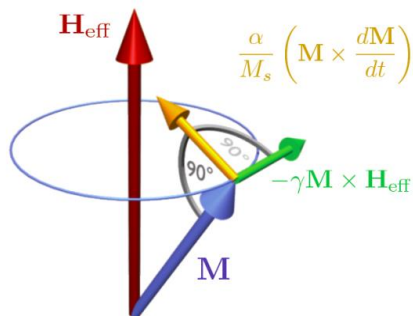
$$f=30 \text{ MHz}- 300 \text{ GHz} / \lambda_{em} =10 \text{ km}- 1\text{mm}$$



(also still used: MHz-GHz electromagnetic waves in cables)

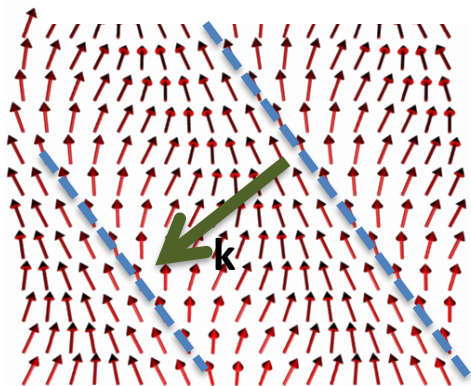
- Free-space propagation  
=> wireless & mobile communications

# Magnetization dynamics



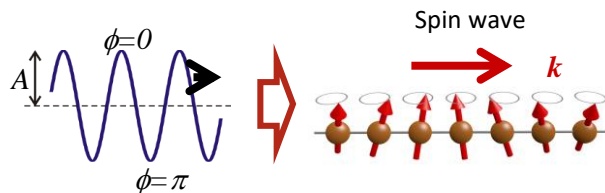
- Coherent coupled oscillations of magnetic moments
- Described by the Landau-Lifshitz-Gilbert equation

⇒ Fundamentals:  
See the two lectures  
of B. Hillebrands  
last Wednesday



P. Pirro, V. I. Vasyuchka, A. A. Serga, and B. Hillebrands,  
*Advances in Coherent Magnonics*, Nat Rev Mater **6**, 1114 (2021).

## Wave-based data processing



- Coherent propagation
- Interference effects and logic
- Intrinsic nonlinearity
- Scalability: down to nanometer wavelength
- Frequency multiplexing
- **Typical frequencies: GHz**

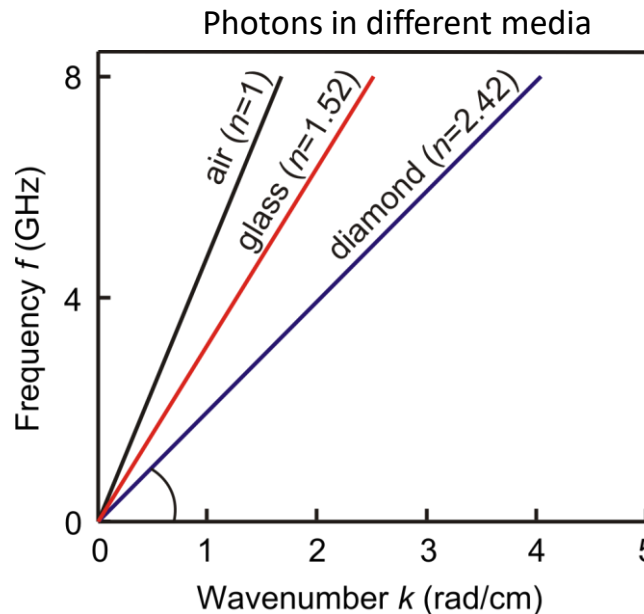
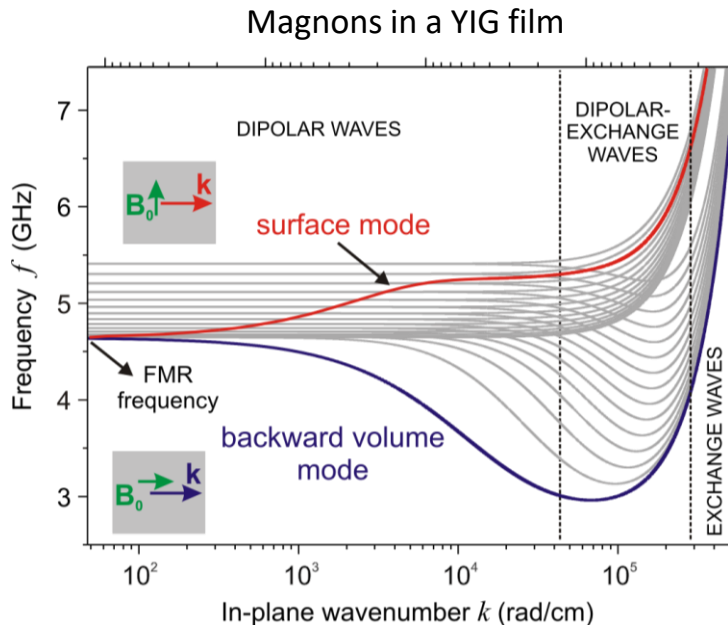
Ferromagnetic resonance of in-plane magnetized film:

$$f_{FMR} = \gamma \sqrt{\mu_0 H (\mu_0 H + \mu_0 M_s)}$$

(effective) magnetic field      saturation magnetisation

$$\gamma \approx 28 \frac{\text{GHz}}{\text{T}}, \mu_0 H = 0 - 1 \text{ T}, \mu_0 M_s = 0.175 \text{ T (YIG)} - 2.2 \text{ T (Fe)}$$

## Magnon versus photon dispersion relations

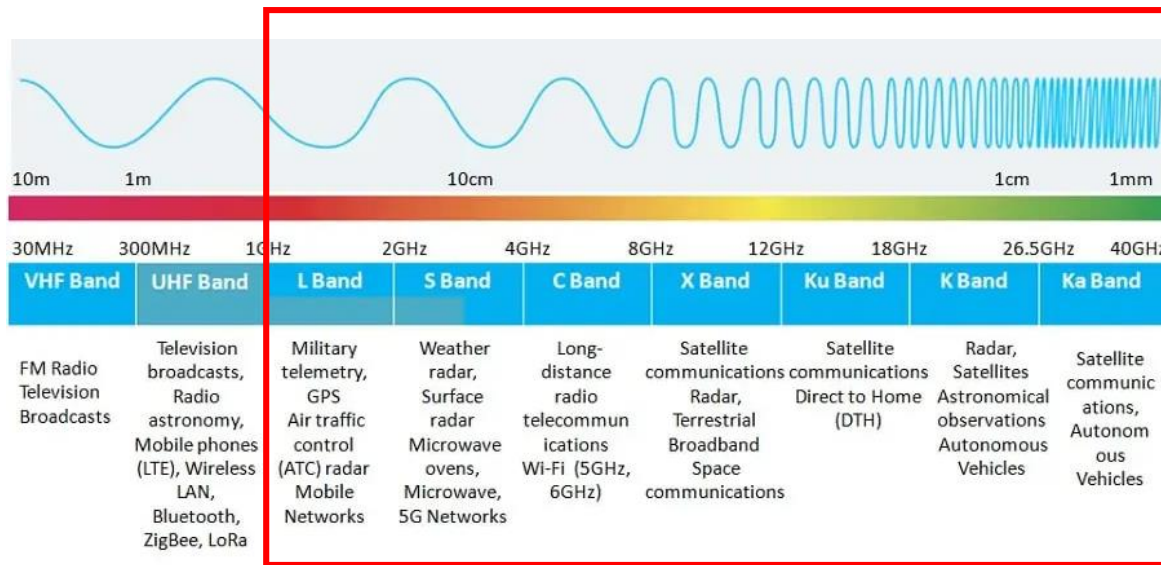


Magnons: about  $10^3$ - $10^5$  higher wavenumbers for the same frequency range:

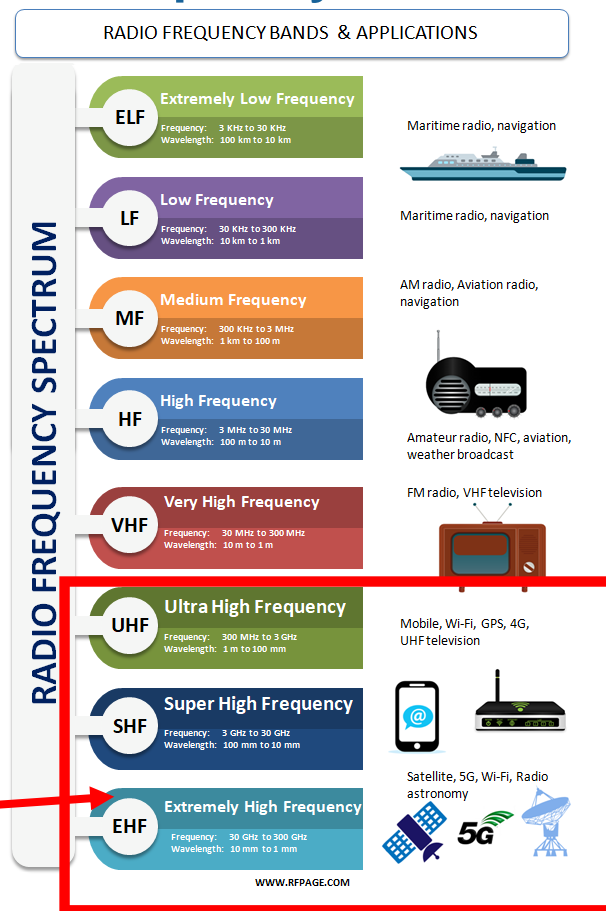
- **Miniaturization** is facilitated: RF devices based on magnons are much smaller
- Magnons travel much **slower** than photons: **RF time** and **phase delay** possible in microstructures



# RF applications in different frequency bands



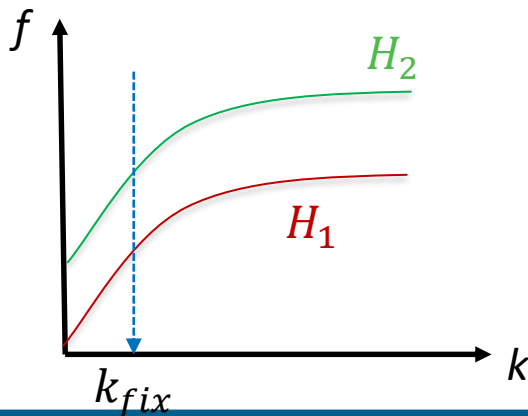
Frequency range mainly targeted by magnonic applications:  
1-60 GHz



### Linear systems

Linear SW propagation

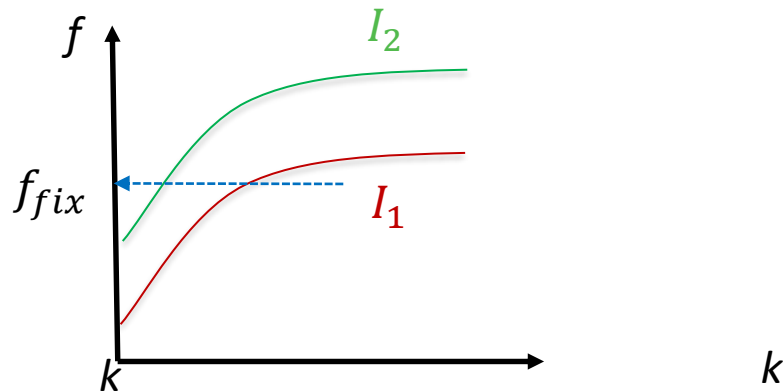
- RF phase shifter
  - RF delay lines
  - RF filters, isolators
- tunable via the applied magnetic fields ( $H_1, H_2$ )



### Nonlinear magnon systems

Nonlinear SW propagation

- Half-adder (AND gate)
  - Ring resonators
  - Magnon Transistor
- tunable via the magnon intensity ( $I_1, I_2$ )



## Linear systems

(Input-output relation independent of RF power/magnon intensity)

- **Phase delay unit**

shifts (only) the phase  $\phi$  of the signal by a shift  $\Delta\phi$  (which should be constant over the bandwidth  $\Delta f$ )

- **Time delay unit**

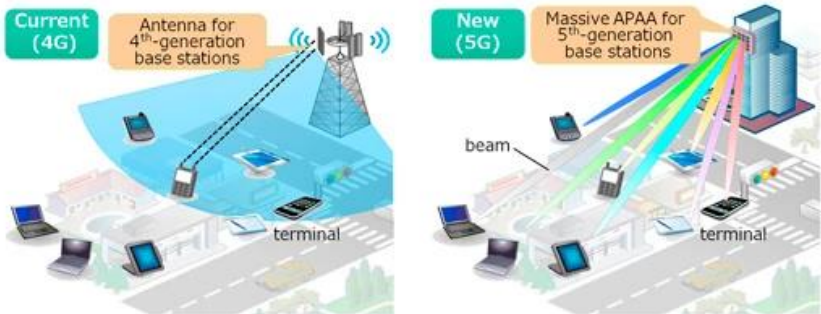
delays a signal by a time delay  $\Delta t$  (all frequency components in the bandwidth should have the same delay to avoid signal distortion)

- **Filter**

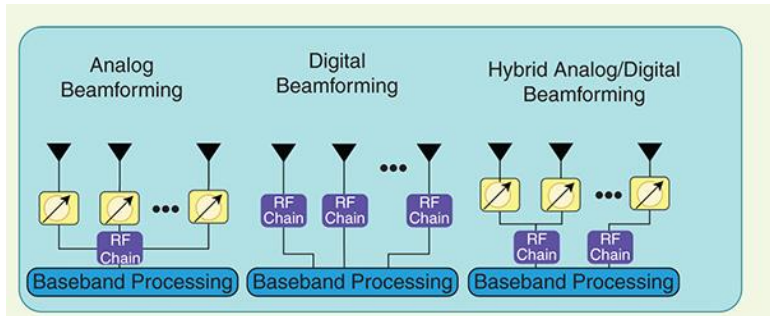
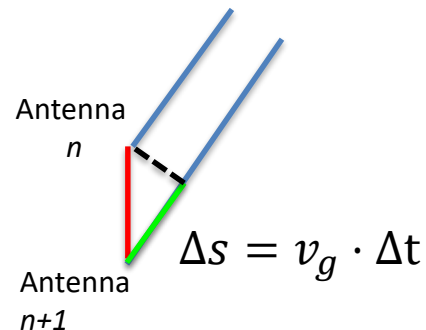
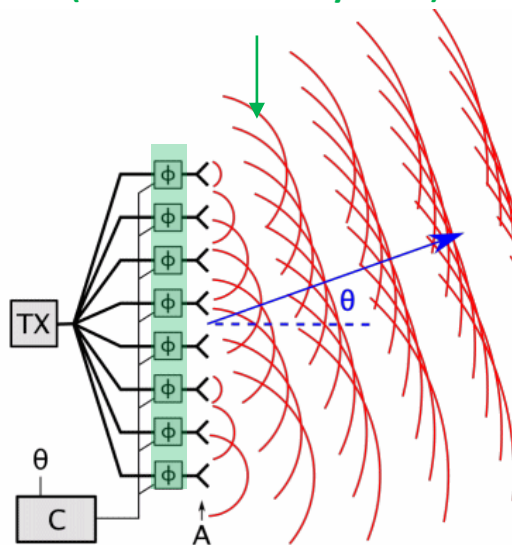
removes a selected frequency range from the spectrum

# Applications of RF delay units

Tunable Directional antennas, e.g. for massive MIMO (“Multiple Input Multiple Output”)



Tunable miniaturized phase delay units (better: time delay units)

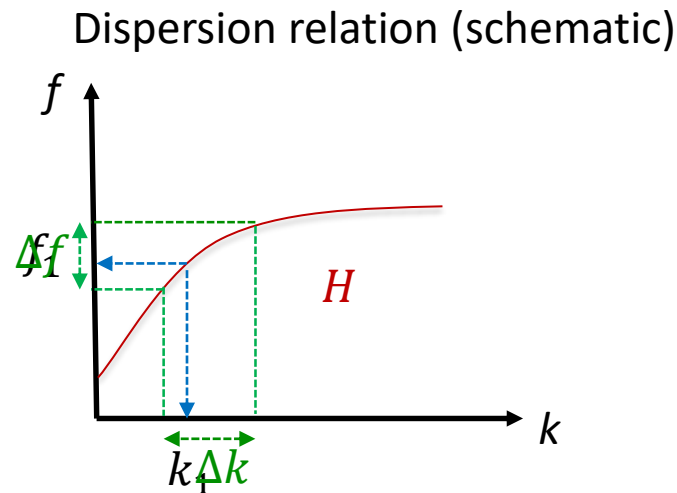
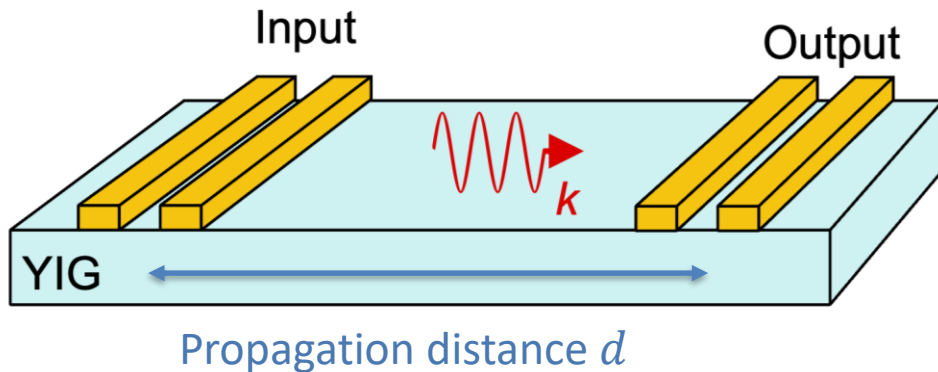


IEEE signal processing

Many RF antennas with different time delays

# 1. “Stationary” magnonic devices

One big advantage of magnonic devices is their reconfigurability using different external fields. For simplicity, let us first understand how a magnonic devices works at a fixed external field.



Phase delay:  $\Delta\phi(f_1) = k_1 \cdot d$  ( $\Delta\phi$  can be  $\gg 2\pi$ )

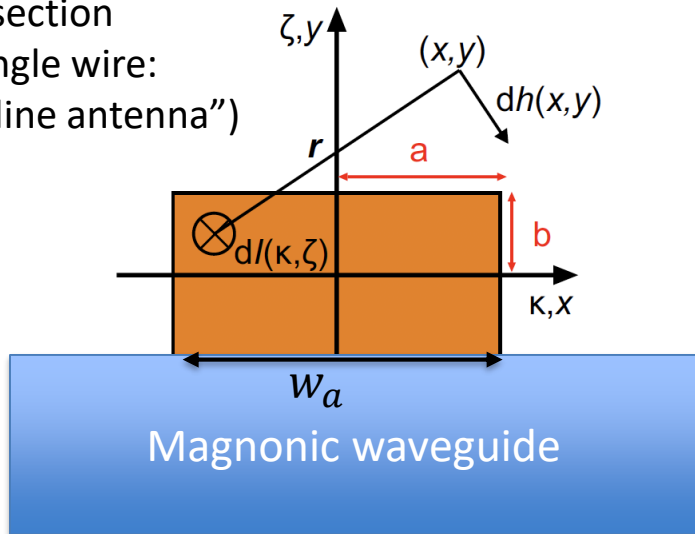
Time delay:  $\Delta\tau(f_1) = \frac{d}{v_G(k_1)} = \frac{d}{\partial f / \partial k(k_1)}$  ( $\Delta\phi$  can be  $\gg 2\pi$ )

Filter function? Bandwidth  $\Delta f$  ? Efficiency of energy transduction?

**General idea:**

- A microwave of frequency  $f$  propagating in a wire-like conductor is creating **AC electron currents** of the same frequency
- According to **Biot-Savart's law** the **AC currents create AC magnetic fields** which can excite the spin waves if their direction has a perpendicular to the static magnetization

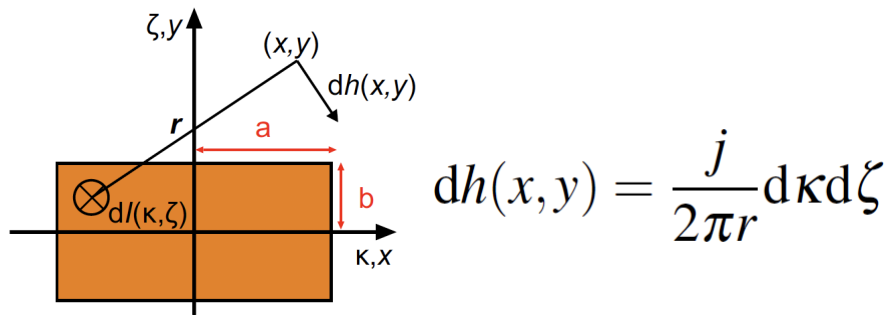
Cross section  
of a single wire:  
("stripline antenna")



In  $z$  direction, the wire is usually considered as infinitely extended

Conductor (orange): Au, Cu, Al, superconductor...

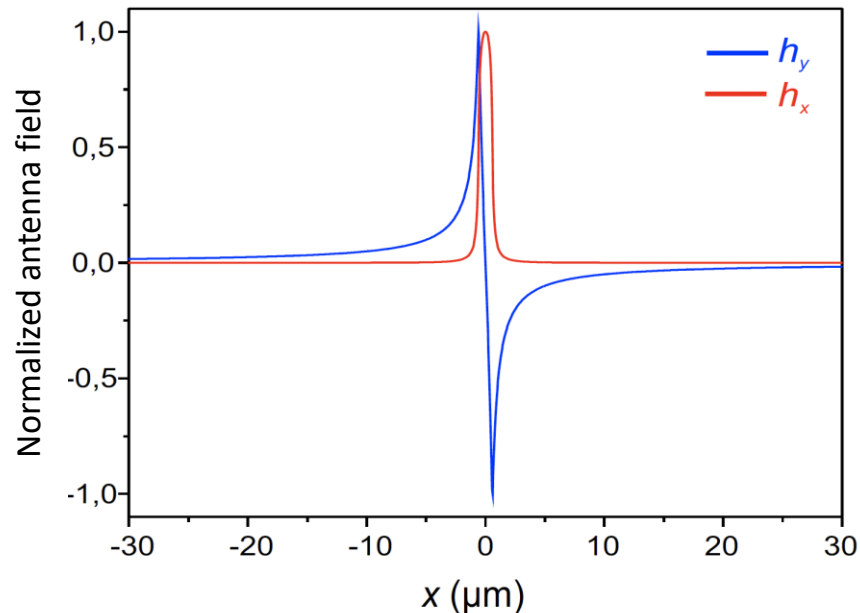
In the magneto-static approximation, Biot-Savart's law can be used to calculate the dynamic fields of the microwave currents:



Integration for a rectangular antenna cross section with width  $2a$  and thickness  $2b$ :

$$h_x(x, y) = -\frac{I}{8\pi ab} \left[ (a-x) \left[ \frac{1}{2} \ln \left( \frac{(b-y)^2 + (a-x)^2}{(-b-y)^2 + (a-x)^2} \right) + \frac{b-y}{a-x} \operatorname{atan} \left( \frac{a-x}{b-y} \right) - \frac{-b-y}{a-x} \operatorname{atan} \left( \frac{a-x}{-b-y} \right) \right] - (-a-x) \left[ \frac{1}{2} \ln \left( \frac{(b-y)^2 + (-a-x)^2}{(-a-x)^2 + (-b-y)^2} \right) + \frac{b-y}{-a-x} \operatorname{atan} \left( \frac{-a-x}{b-y} \right) - \frac{-b-y}{-a-x} \operatorname{atan} \left( \frac{-a-x}{-b-y} \right) \right] \right]$$

- Here,  $I$  is the total current running in the antenna
- Similar expression for the field in  $y$  direction



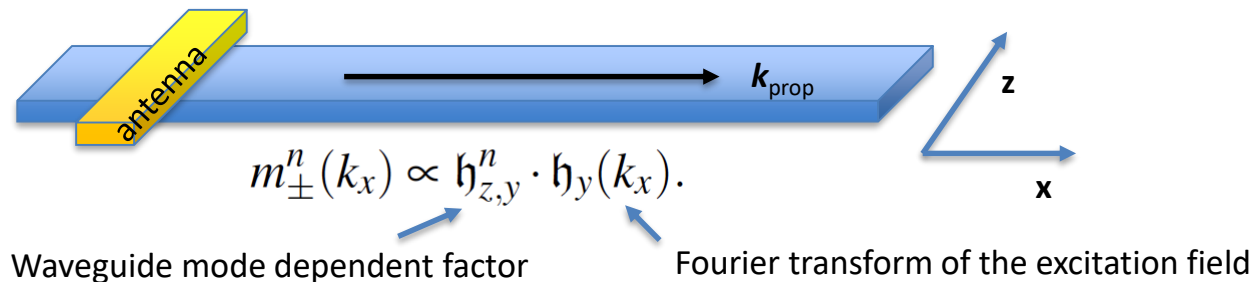
width:  $2a = 1.15 \mu\text{m}$ ,  
thickness:  $2b = 0.22 \mu\text{m}$

## Excitation efficiency of microwave antennas

The excitation efficiency of spin waves is calculated based on the general connection of the dynamic magnetization with the dynamic microwave field (external field) and the dynamic dipolar field of the spin wave via the susceptibility tensor:

$$\mathbf{m}^{\text{dyn}} = \hat{\chi}(\mathbf{h}^{\text{MW}} + \mathbf{h}^{\text{d}})$$

- Results depend on the geometry (Surface waves, Backward Volume, Forward Volume...)
- **Efficiency of excitation is always proportional to the Fourier transform of the microwave field along the propagation direction** (and also on the mode profile across the waveguide)



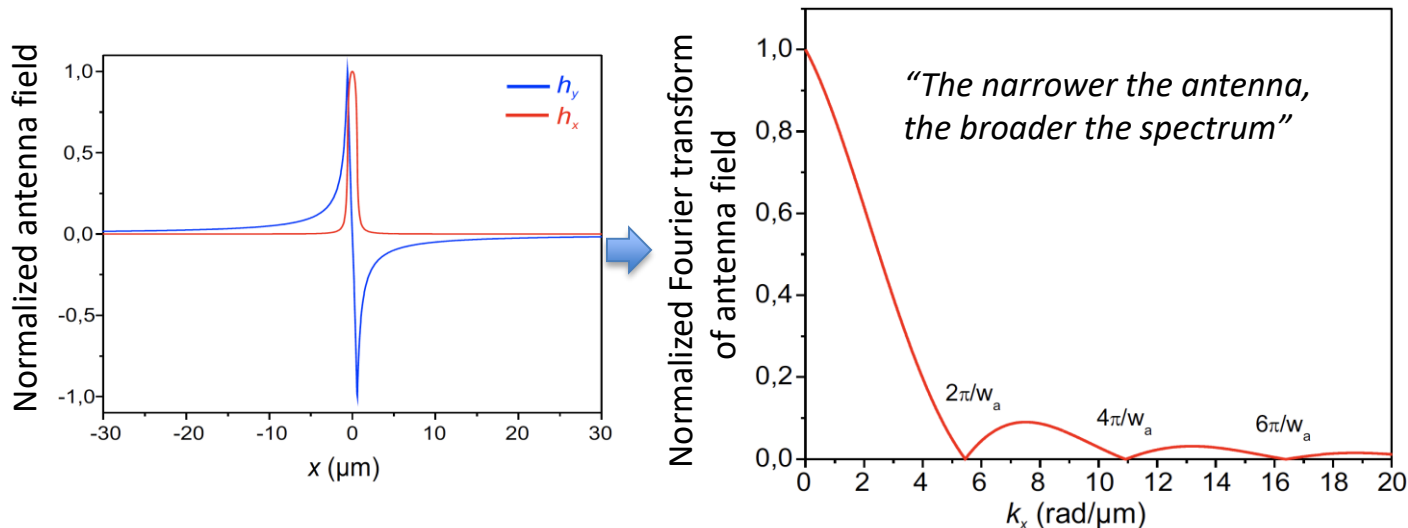
T. Schneider, A. A. Serga, T. Neumann, B. Hillebrands, M. P. Kostylev, Phase reciprocity of spin-wave excitation by a microstrip antenna, *PRB* **77**, 214411 (2008).

V. F. Dmitriev, B. A. Kalinikos, Excitation of propagating magnetization waves by microstrip antennas, *Sov. Phys. J.* **31**, 875 (1988).



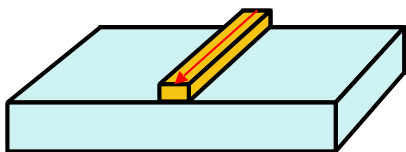
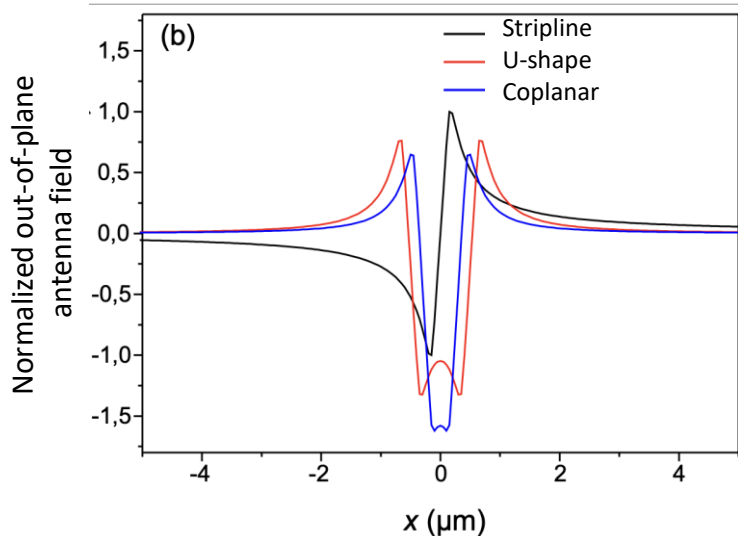
## Excitation efficiency: wave vector dependence

- Spatial extent of the antenna field defines the “effective linear momentum” which can be provided by the antenna to create linear momentum ( $p = \hbar k$ ) of spin waves
- The **wave-vector-dependent excitation efficiency** is mainly given by the **Fourier transform of the antenna fields**

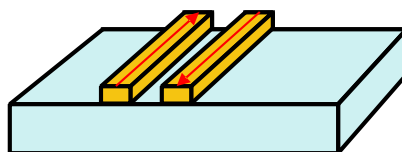
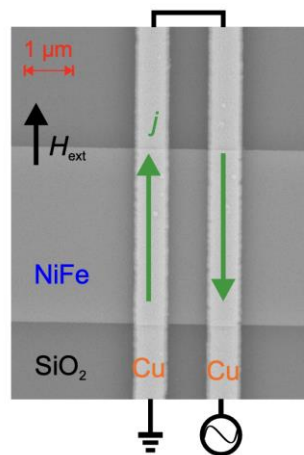


**Fourier transform of both field components is the same!**

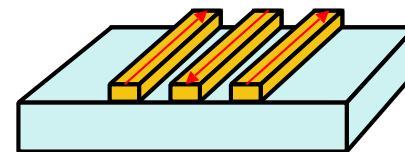
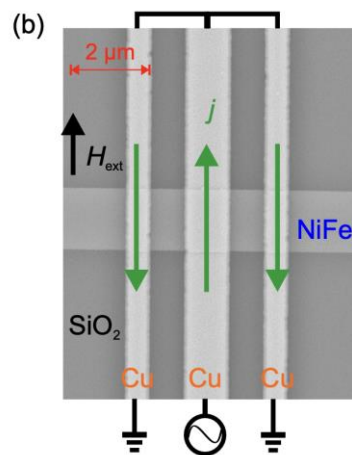
## Antenna type examples



Stripline

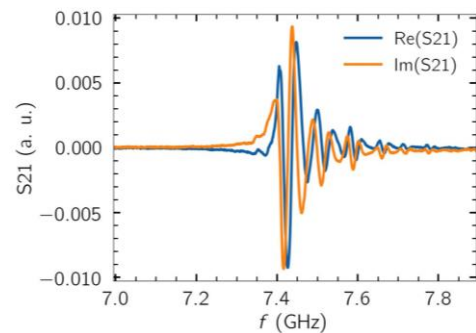
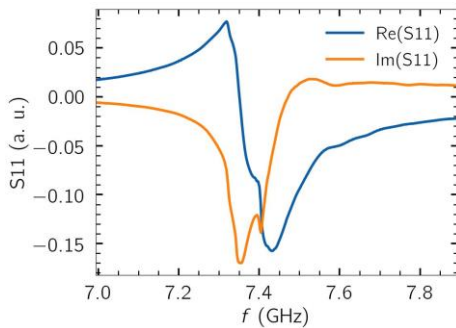
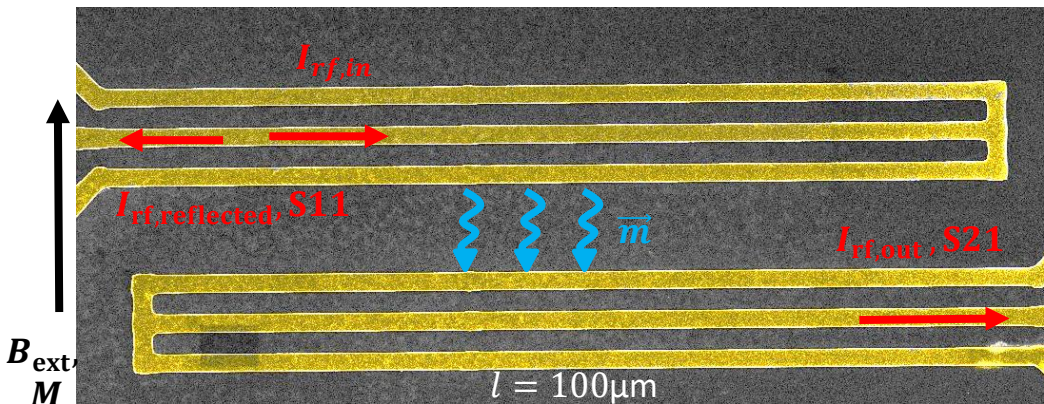
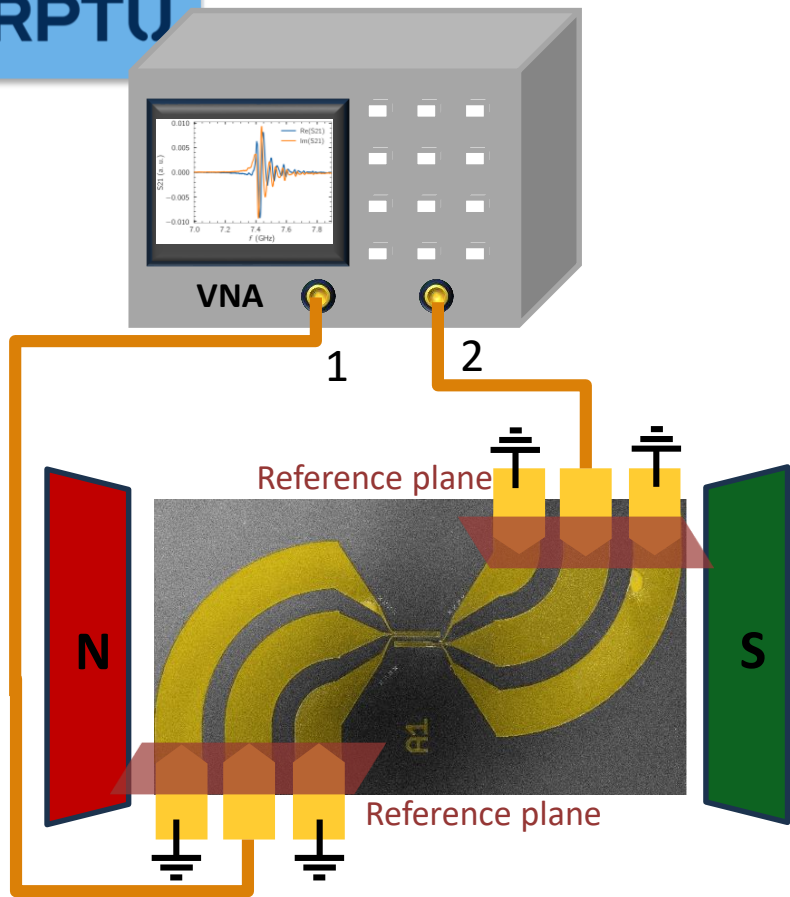


U-shape

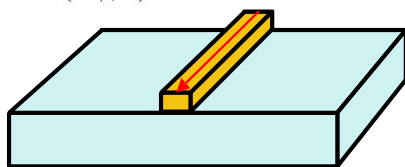
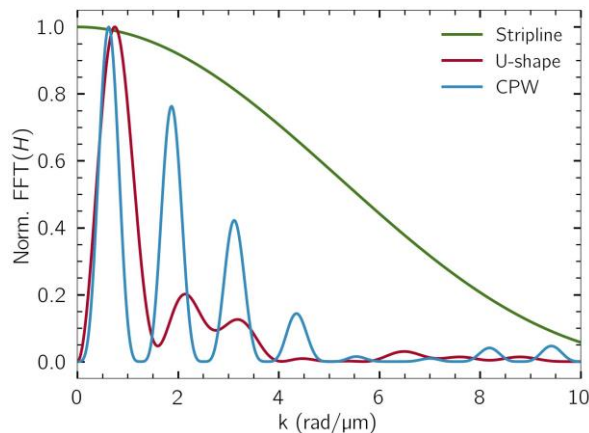
Coplanar waveguide  
(CPW)

Due to the alternating current directions, the out-of-plane field of U-shape and CPW is much better confined

# Coplanar waveguide (CPW)

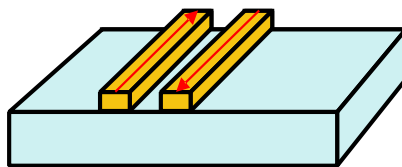
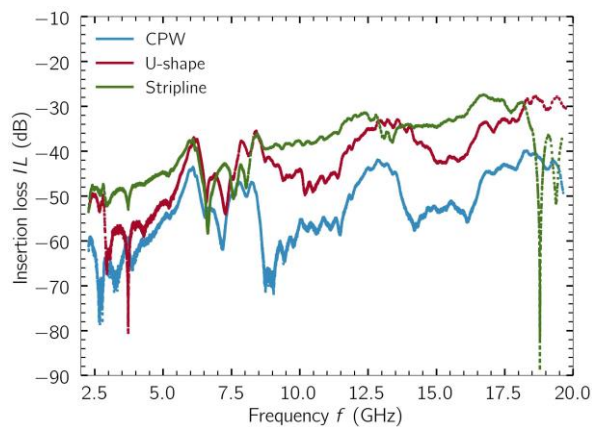


Wave-vector spectrum



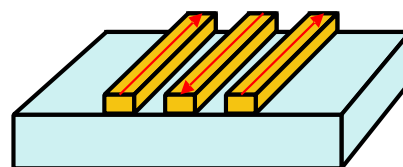
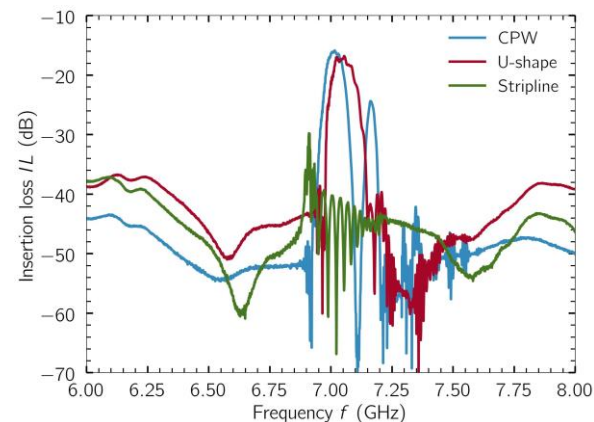
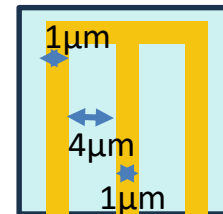
Stripline

Unwanted electro-magnetic crosstalk



U-shape

Signal Transmission

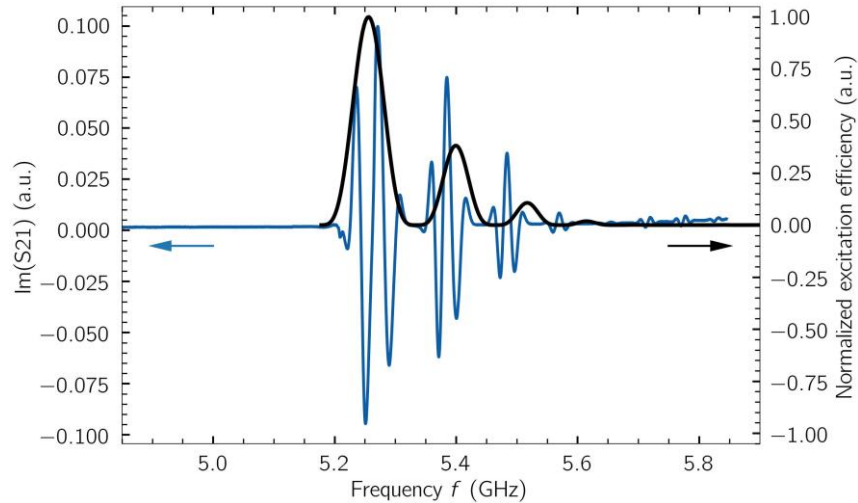
Coplanar waveguide  
(CPW)

➤ **CPW offers lowest e.m.crosstalk (in this example also the best transmission)**

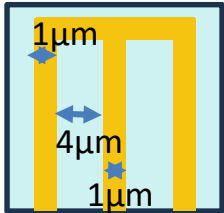
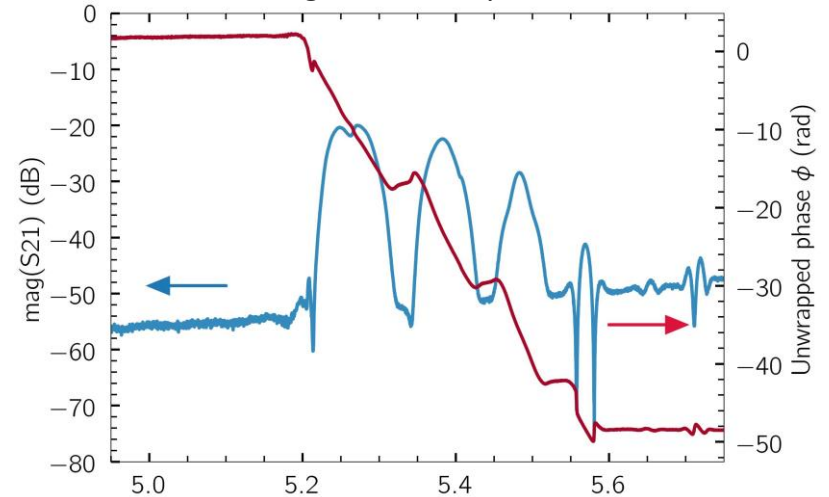
100nm YIG

 $P = -20\text{dBm}$ 

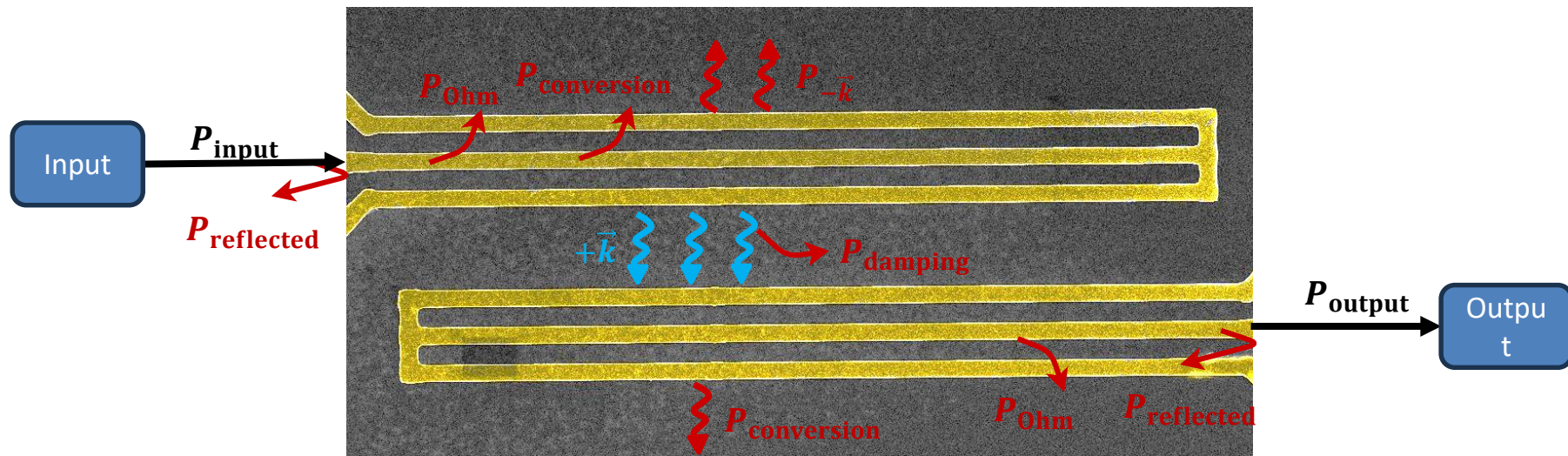
Transmission spectrum



Magnitude and phase

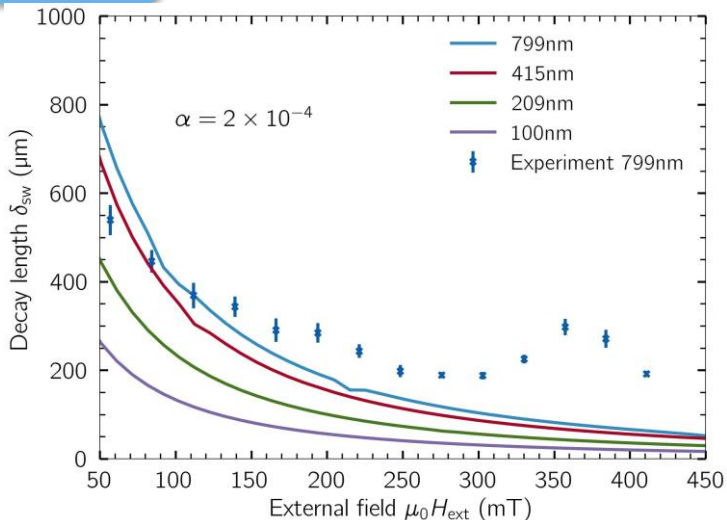


- Transmission spectrum defined by antenna=> use antenna shape to build RF filters
- Strong phase accumulation by spin-waves:



➤ How much can these losses be reduced?

## Propagation losses



- (Gilbert) damping leads to propagation loss of SW

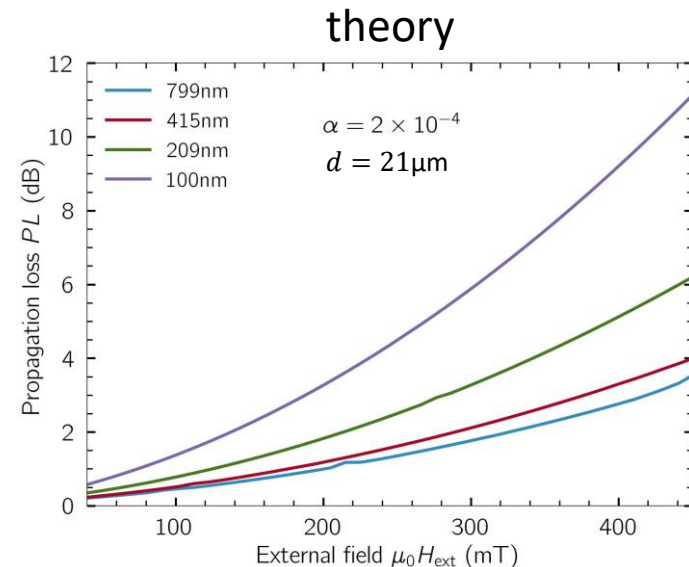
- $PL = -20 \log_{10}(\exp(-\frac{d}{\delta_{SW}}))$

- $\delta_{SW} = \tau_{SW} \cdot v_g$

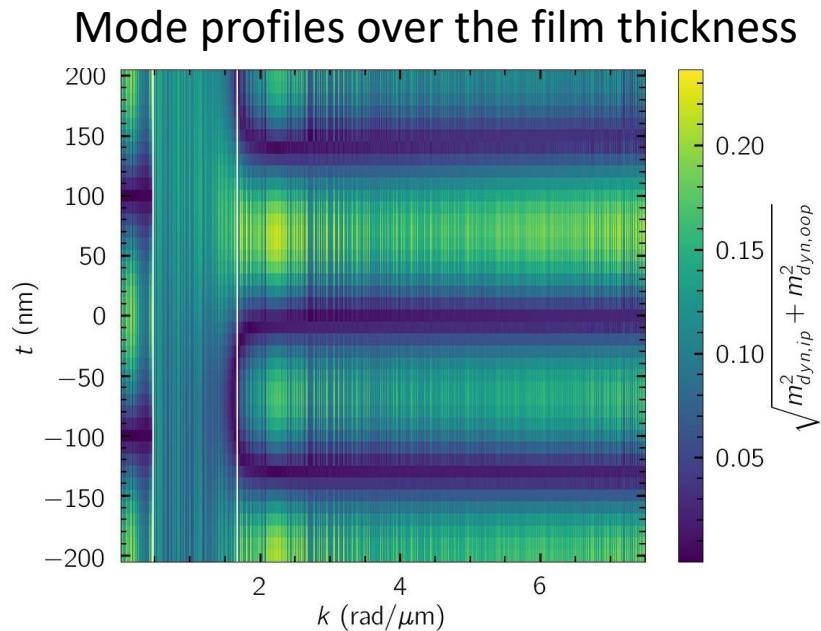
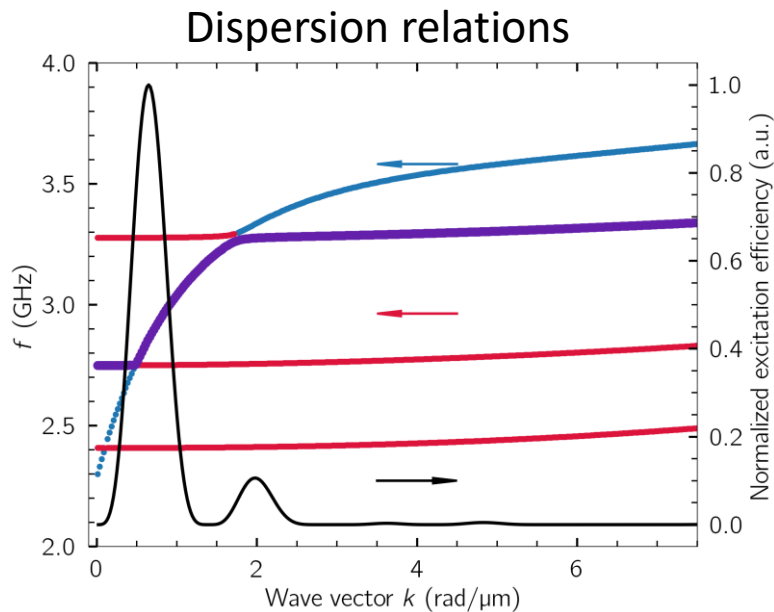
SW lifetime

$$\tau_{SW} \propto \frac{1}{(\alpha \cdot \omega)}$$

➤ Reduction of propagation losses for thicker films



YIG film, 400 nm thickness, 30 mT external field, surface wave geometry

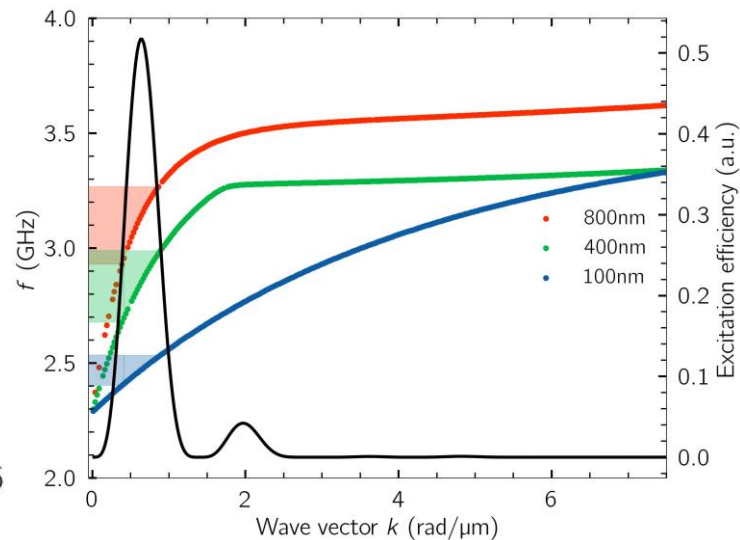
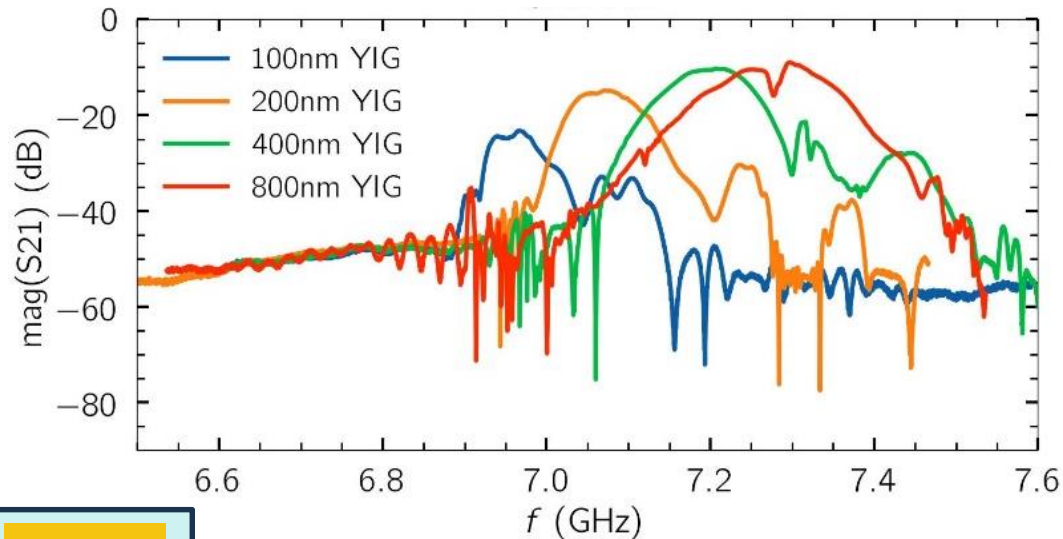


Surface mode hybridizes with modes quantized over the film thickness  
(Perpendicular standing spin wave modes, PSSW) => “hybridization gaps” in the spectrum



$$P = -20\text{dBm}$$

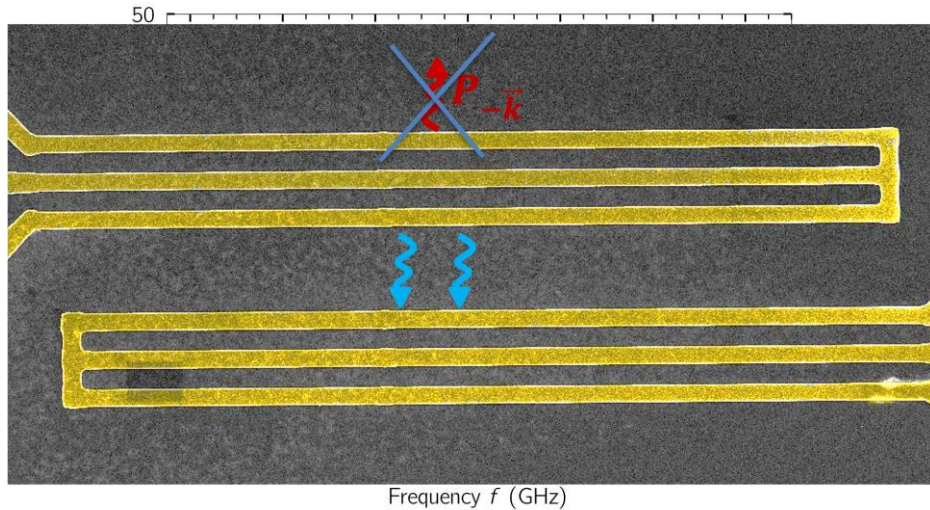
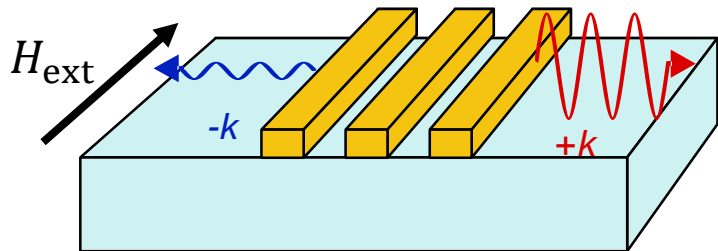
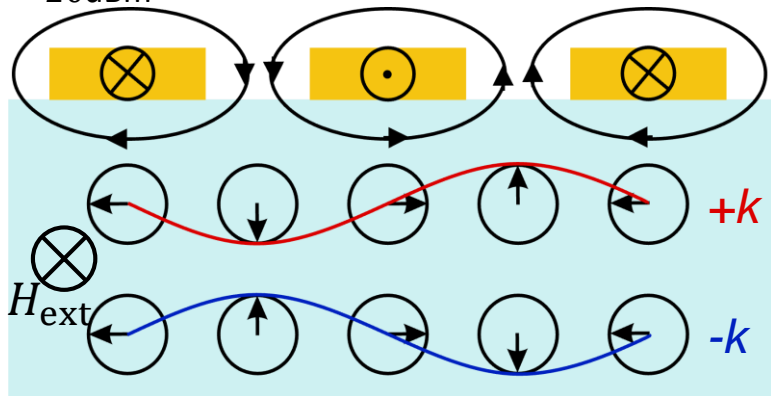
$$B_{\text{ext}} = 166\text{mT}$$



➤ Steeper dispersion relation increases bandwidth for same antenna

# Isolation / Non-reciprocal excitation of the surface wave ( $k \perp M$ )

$P = -20\text{dBm}$



- Directed flow of energy: lower losses
- Magnonic delay line is also a RF isolator

## Non-reciprocity doesn't always have the same reason

**Non-reciprocal excitation** caused by the *break of time reversal symmetry* + excitation by *both antenna field components* (in-plane and out-of-plane)

=> **NOT** connected to the spatial non-reciprocity of the Surface Wave (MSSW)

Most antennas excite only the Surface Wave using both components, that's why the other modes (Backward Volume or Forward Volume) are (usually) excited reciprocally.

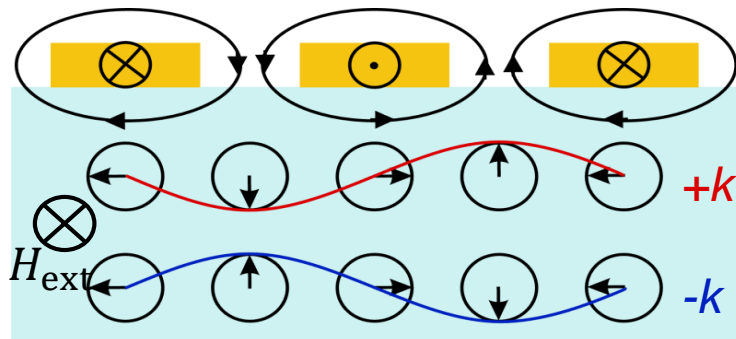
~~Non-reciprocal dynamic mode profile~~

~~$$m_x \sim \exp(-kx)$$~~

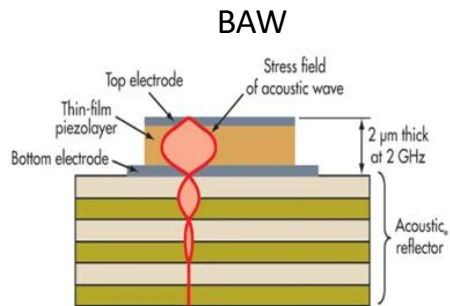
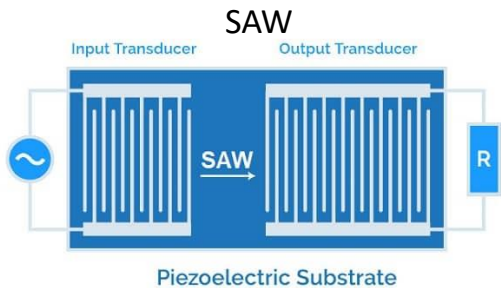


~~$\otimes B_0$~~

“Helicity” of magnetization precession

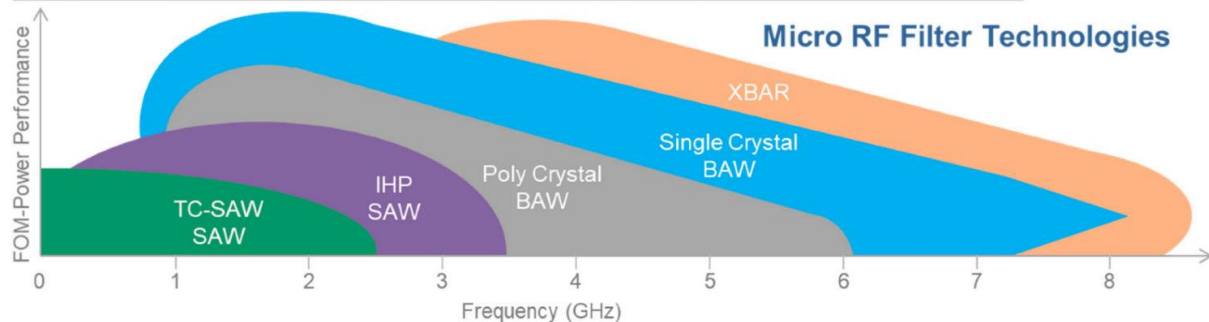


Miniaturized filters/phase shifters/time delay units in today's RF applications: **Surface and Bulk Acoustic Wave** systems



BAW filter structure

## Market / Application



Chen, P.; Li, G.; Zhu, Z., *Micromachines* 2022, 13, 656, <https://doi.org/10.3390/mi13050656>

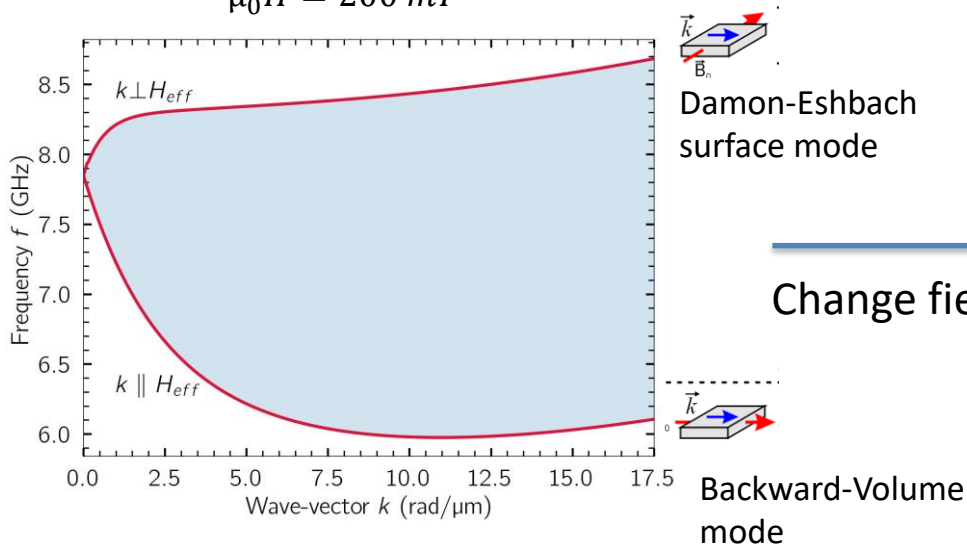
SAW /BAW:

- **No tunability**
- **No solutions above 10 GHz**

## 2. Tuning magnonic devices

**Example: Dispersion relations of dipolar-exchange spin-waves** in an 800 nm thick film of Yttrium Iron Garnet (YIG) with in-plane magnetisation:

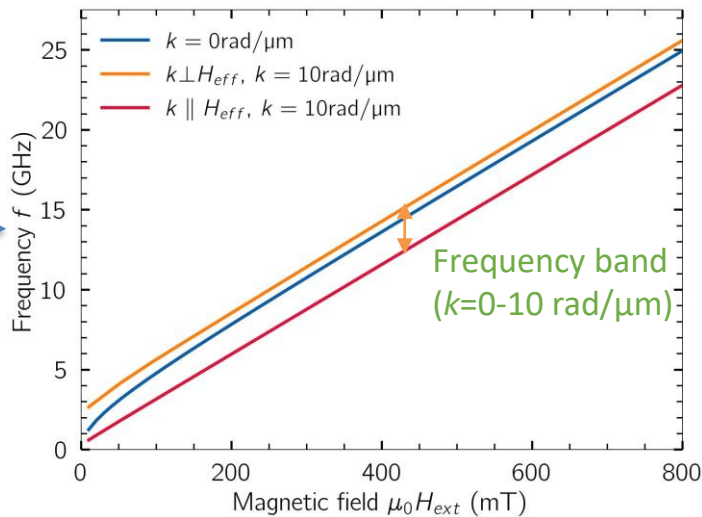
$$\mu_0 H = 200 \text{ mT}$$



Previous lectures:

*Dispersion relation* of magnons depends on propagation direction, saturation magnetization, film thickness, magnetic field.....

### Magnetic field dependence of frequencies

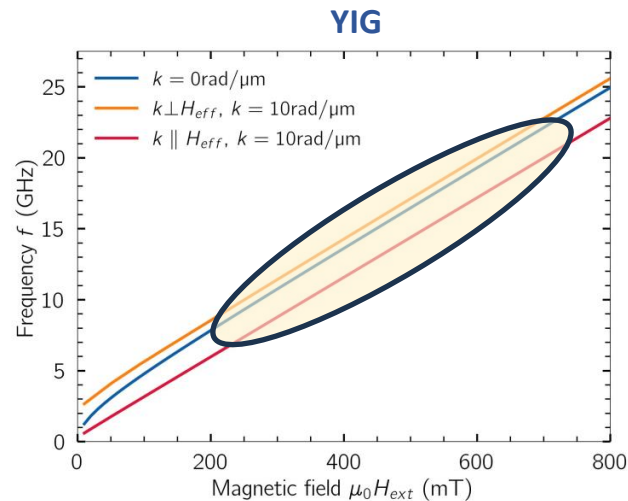
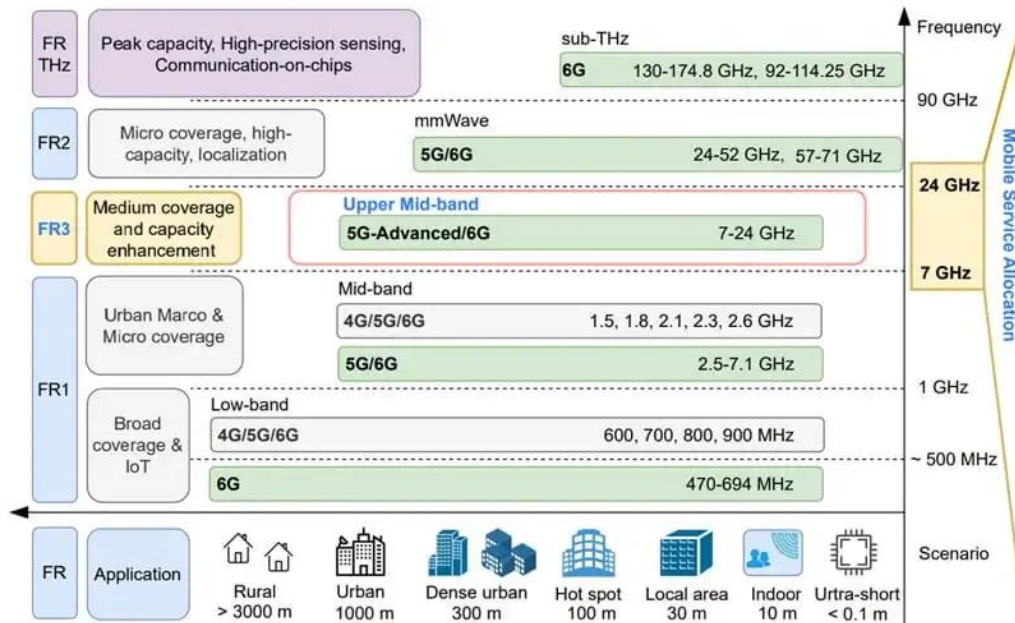


$$f_{FMR} = \gamma \sqrt{\mu_0 H (\mu_0 H + \mu_0 M_s)}$$

(effective) magnetic field      saturation magnetisation

$$\gamma \approx 28 \frac{\text{GHz}}{\text{T}}, \mu_0 H = 0 - 1 \text{ T}, \mu_0 M_s = 0.175 \text{ T (YIG)} (2.2 \text{ T for Fe)}$$

# Frequency bands available for magnonics



Important for cost-efficient 6G

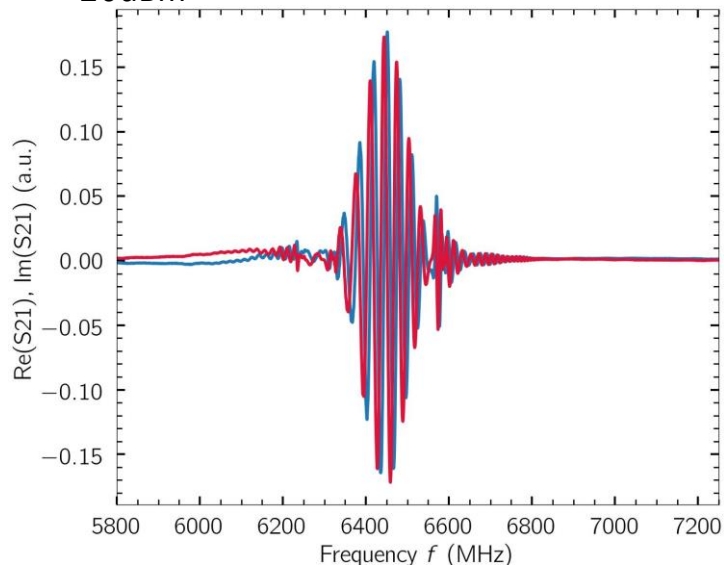
=> magnonic applications based on YIG are well suited for FR3 (7-24 GHz), no miniaturized solutions on the market

## Tunable signal transmission in time domain

400nm YIG

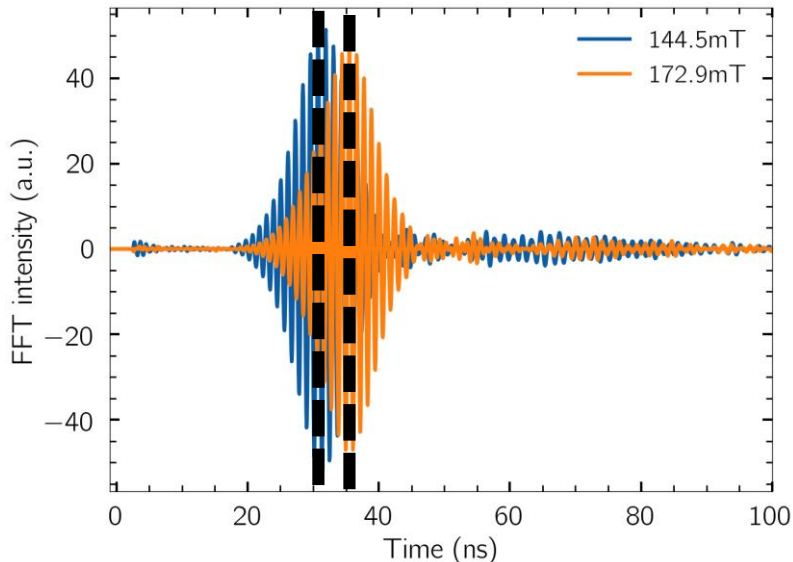
 $P = -20\text{dBm}$ 

Frequency domain

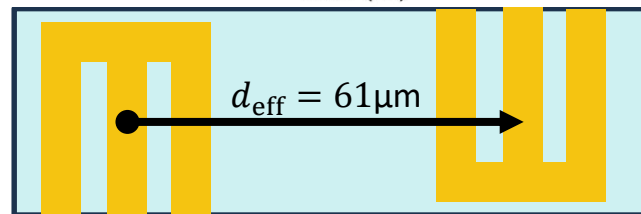


FFT

Time domain

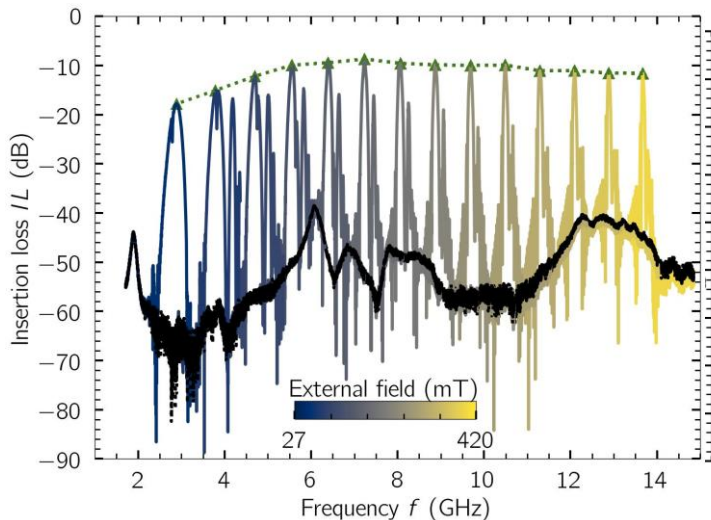


➤ Induced time delay by changing external field



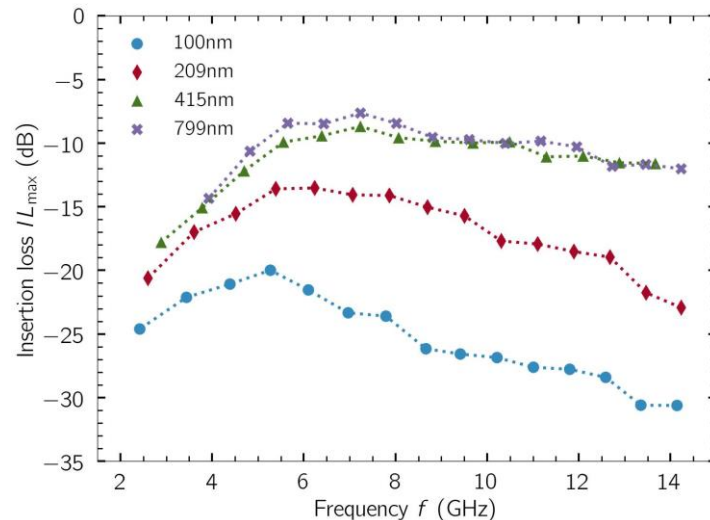
# Frequency dependence

## Field dependence



- Signal over 30dB larger than background
  - No background removal needed

## Thickness dependence



- Insertion losses **<10dB** achieved
- Lower losses with optimization possible
  - Tunable operation over large frequency range



# Commerical magnonic components

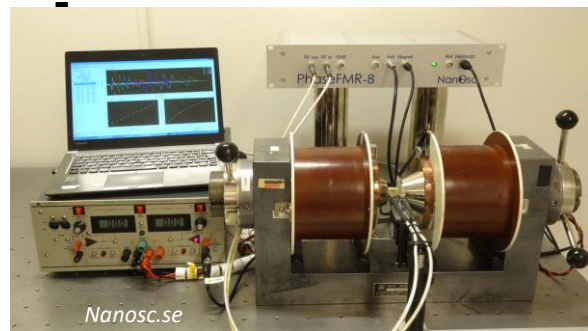


Yttrium-Iron-Garnet tuned **RF oscillators**  
(YTO) produced by Rohde & Schwarz.

=> Bulky, power-hungry due charge current based  
magnetic fields

### Strengths of Magnonics

- Tunability via magnetic fields
- Low power loss inside the magnonic domain



Electromagnet

*Standard experimental approach:*

- Use an **electromagnet** to provide the bias field

=> **Large energy consumption (Joule heating) outside the magnonic domain**

Challenge to create **energy-efficient** devices:

How to **provide magnetic fields (static and dynamic)** in an energy efficient, compact and scalable manner?

## Magnonic systems for signal processing

### Linear systems

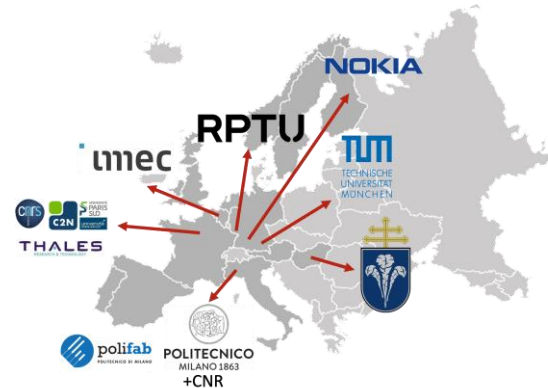
Linear SW propagation

- RF time-delay units
- RF filters, isolators tunable via the applied magnetic fields ( $H_1, H_2$ )

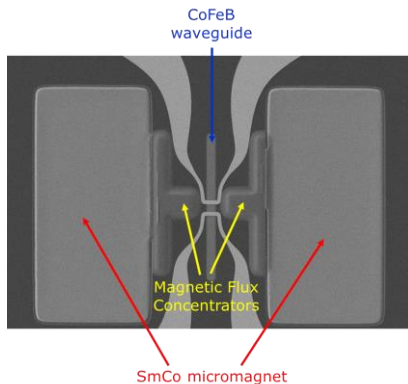
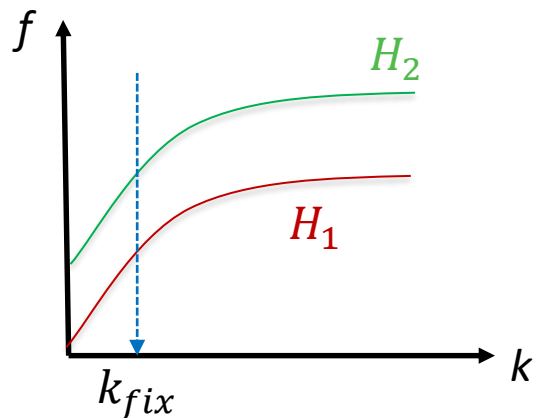


Funded by the European Union

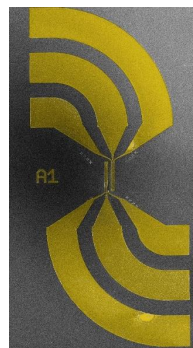
[www.mandmems.eu](http://www.mandmems.eu)



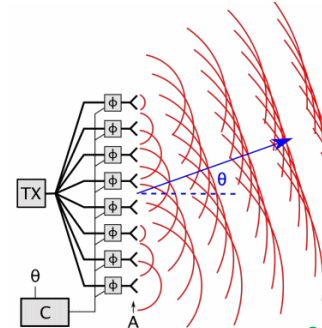
### Standalone magnonic device



### High frequency SW delay lines



### Tunable miniaturized time delay units

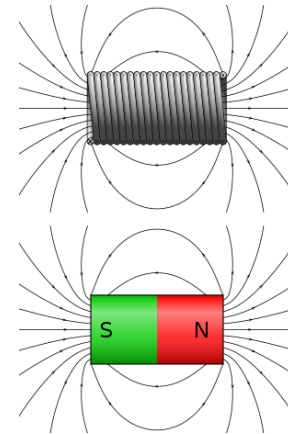


for 5G, 6G

**Replace electric currents** as sources for magnetic fields to maximize energy efficiency and scalability while maintaining tunability of the devices.

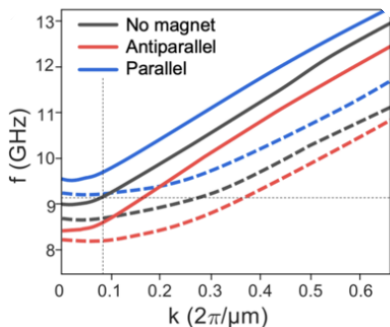
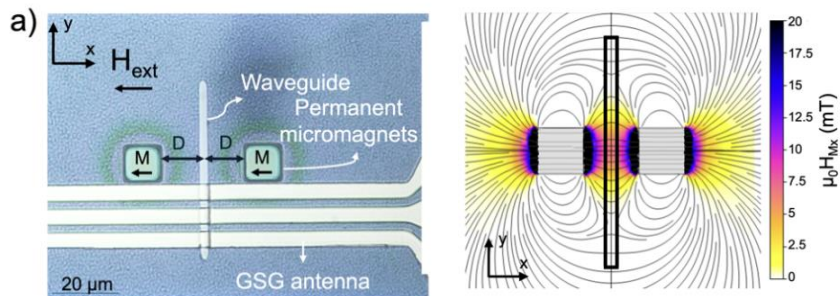
Two magnetic field sources are used in M&MEMS:

- **Permanent magnets**  
=> Tunability is provided by MEMS
- **Magneto-elastic fields**  
=> Strain is provided by surface acoustic waves (SAW)

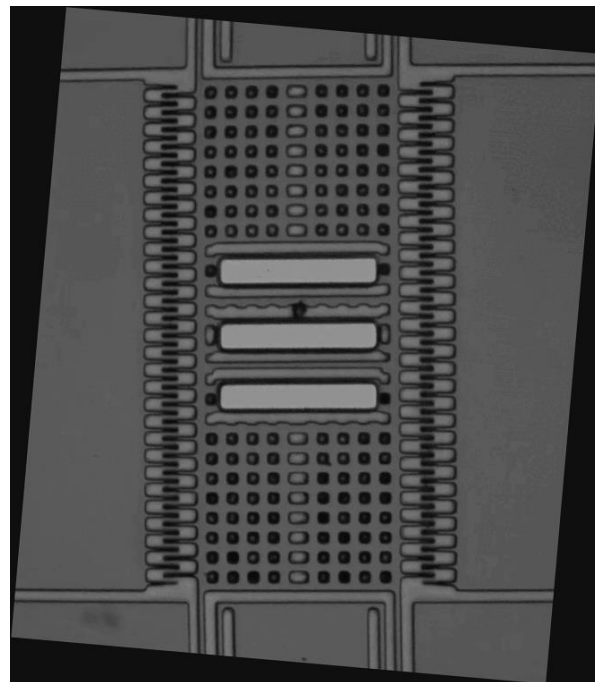


Both approaches allow for a **voltage-based** control and tuning of the magnonic devices with low energy consumption.

## SmCo permanent magnets

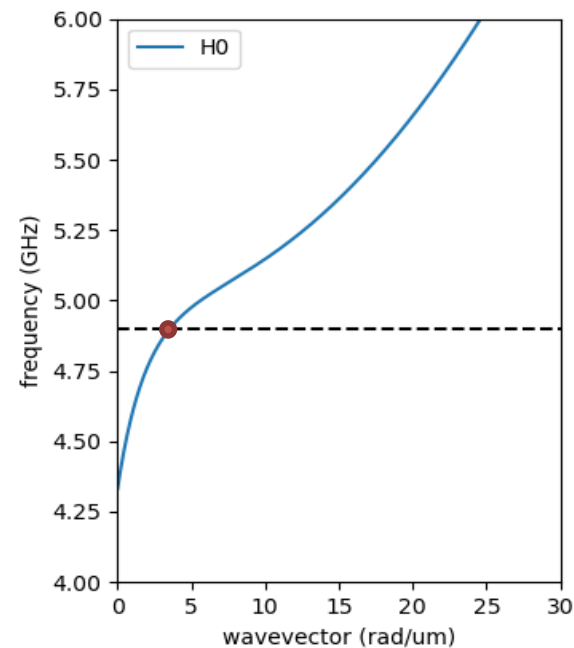
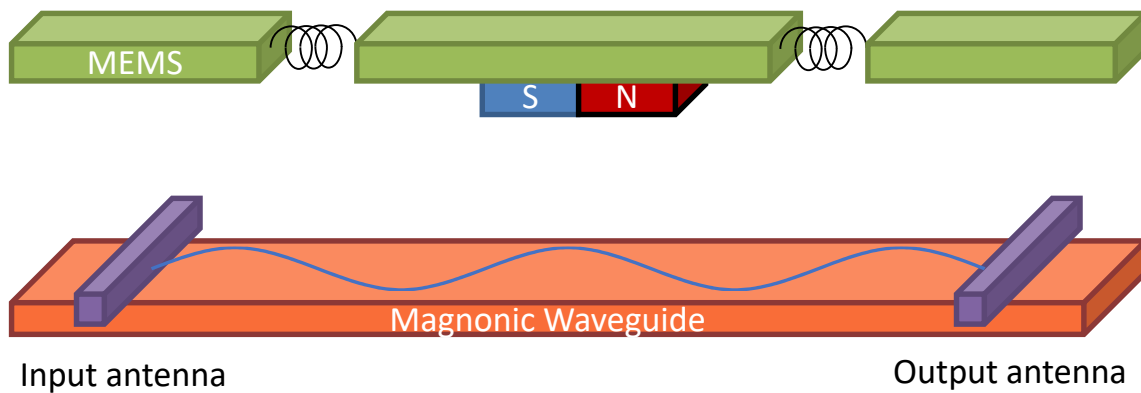


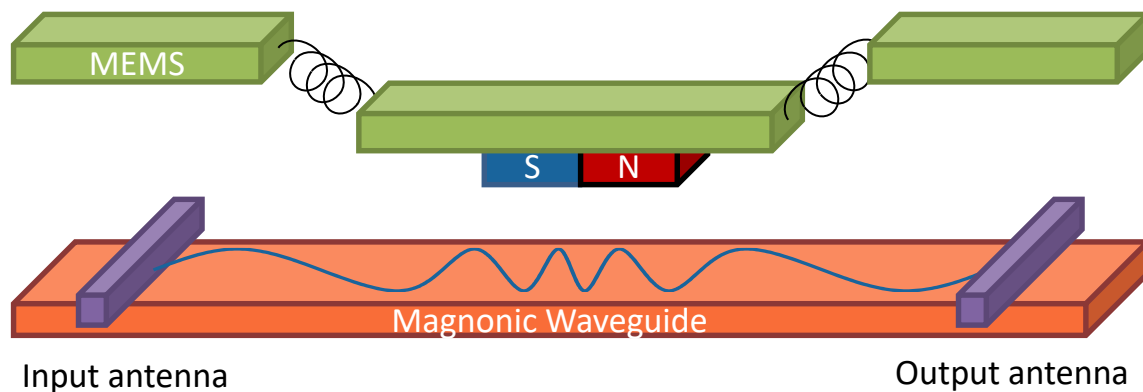
M. Cocconcelli et al., *Tuning Magnonic Devices with On-Chip Permanent Micromagnets*, arXiv 2406.03206 (2024).



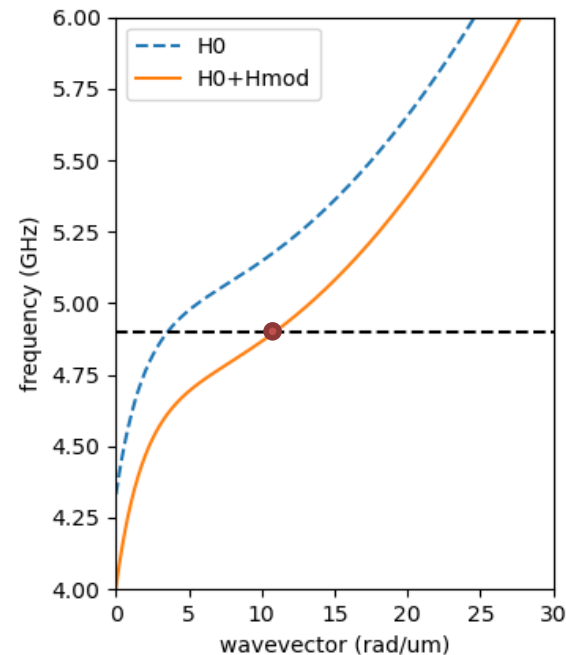
Permanent magnets on a laterally moving MEMS  
(Micro electronic mechanical system)

## Device concept: Phase / time delay unit

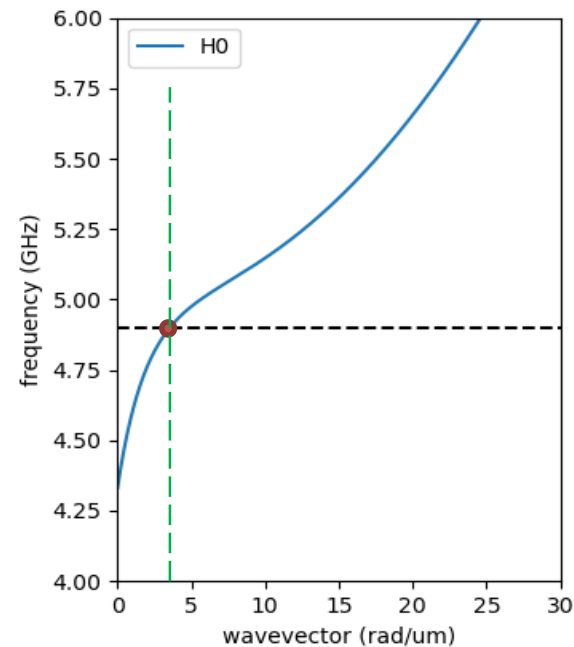
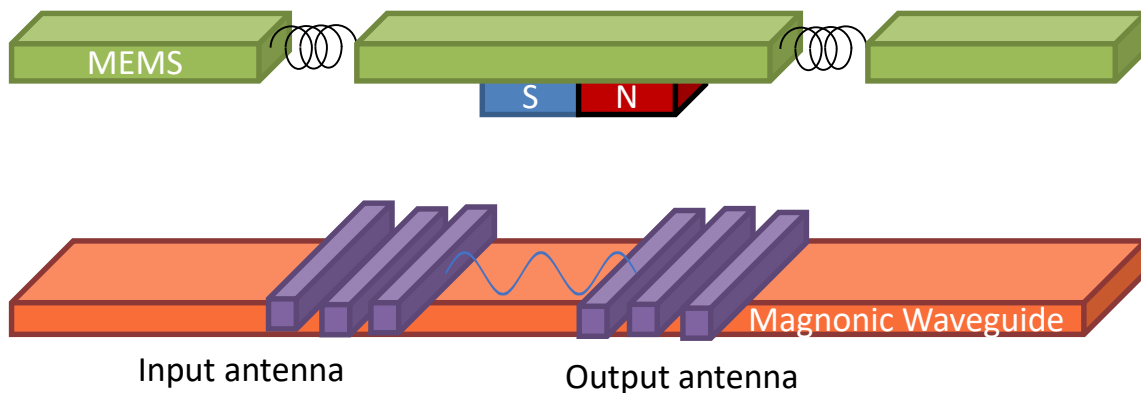




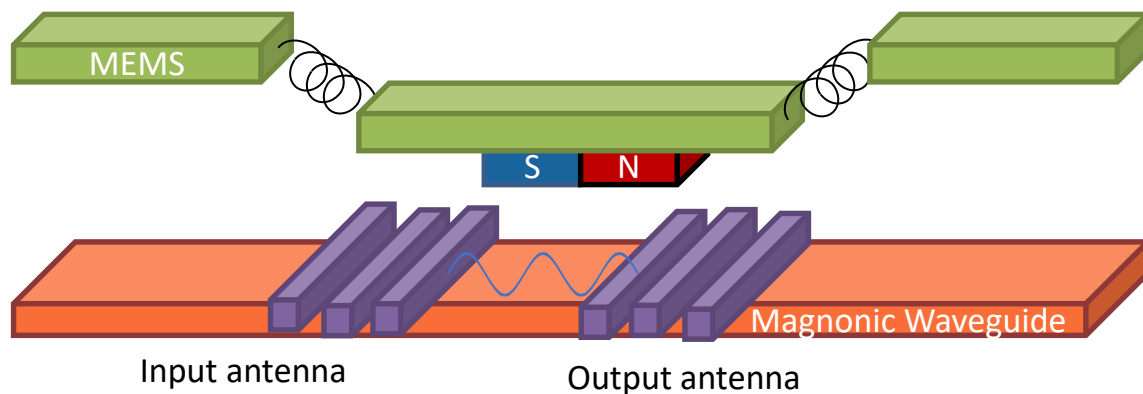
- Modulation of spin-wave wavelength between the antennas  
=> **phase** and **time delay** are tuned
- **No influence on the input/output antenna** => constant amplitude



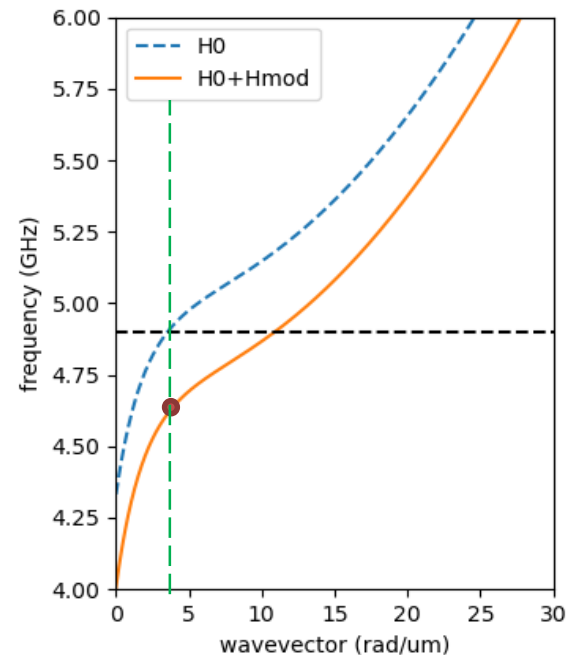
## Device concept: Filter







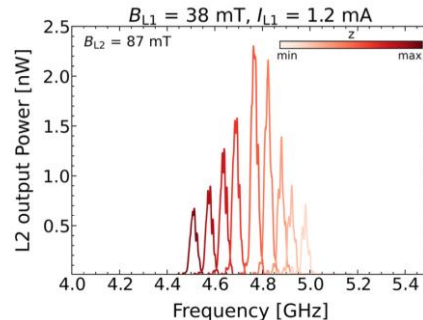
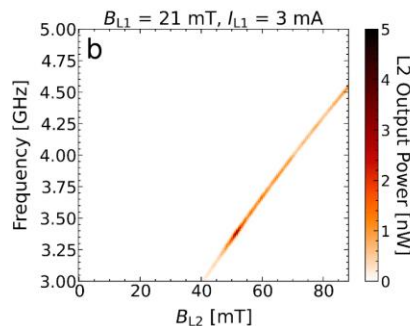
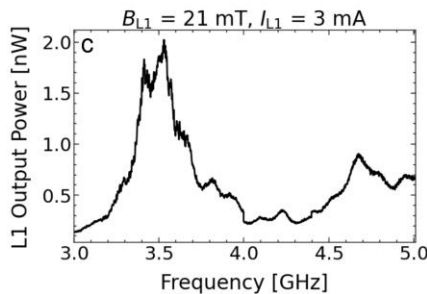
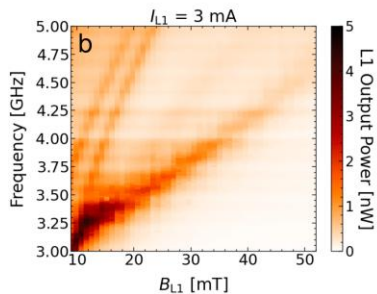
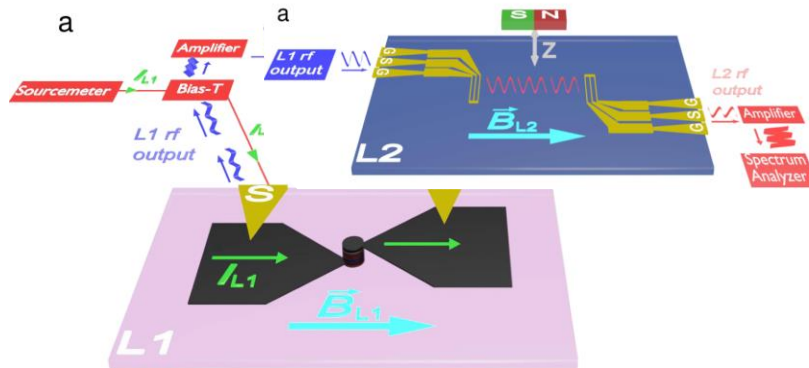
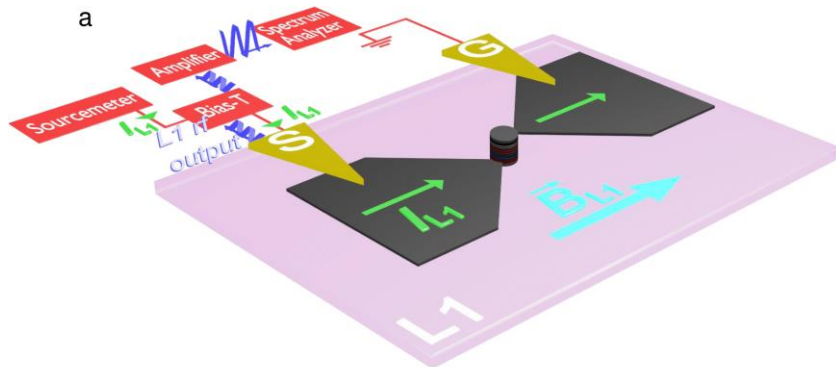
- Modulation of spin-wave resonance frequency **at the antennas**  
=> transmission frequency is shifted



# Example: Filter the signal from an STNO

## STNO filtered by magnonic filter

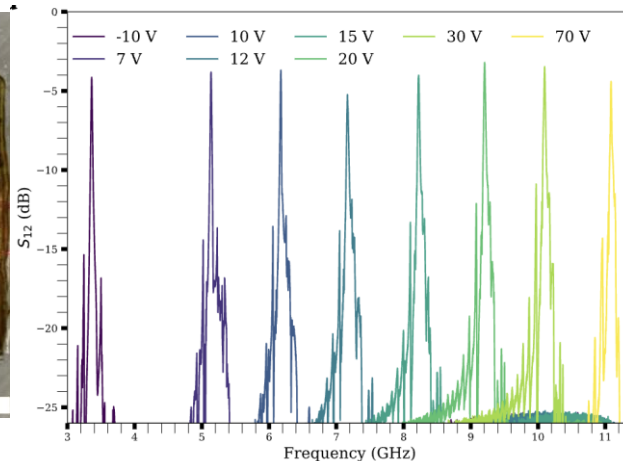
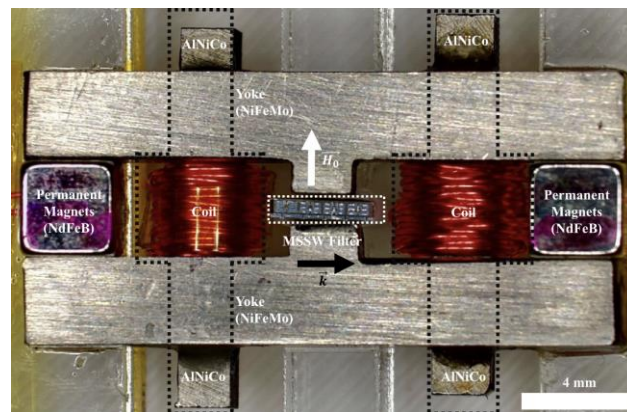
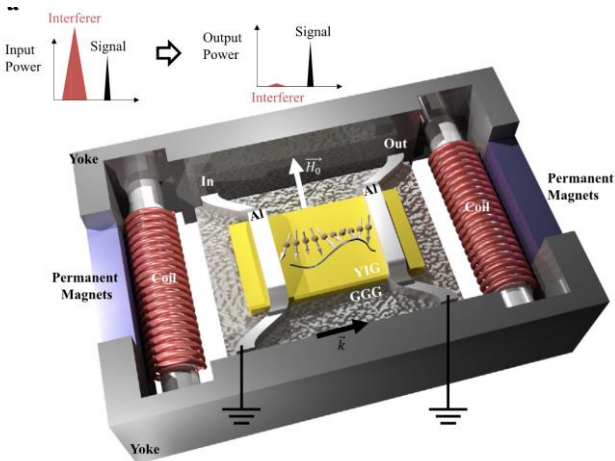
### Only STNO



STNO from U. Ebels

Filtered spectrum of STNO as function of magnet position

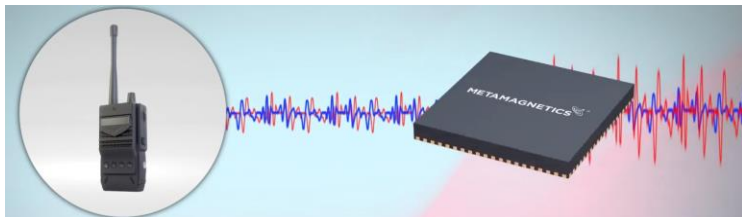
## Zero-power reconfigurable magnonic RF filter



Short current pulses magnetize the AlNiCo ( $\Rightarrow$  see lecture of P. Tozman) magnets partially  $\Rightarrow$  depending on the pulse voltage, different bias fields/resonance frequencies are reached

X. Du, et al., *Frequency Tunable Magnetostatic Wave Filters with Zero Static Power Magnetic Biasing Circuitry*, Nat. Commun. **15**, 3582 (2024).

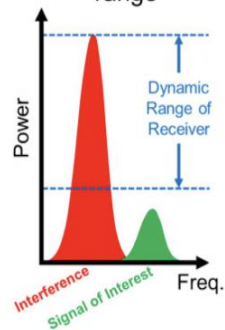
You can buy it:



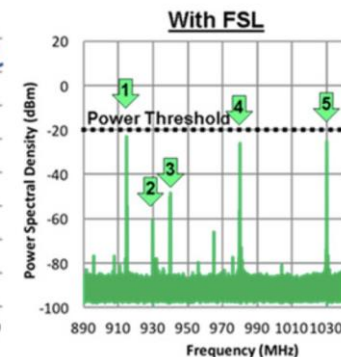
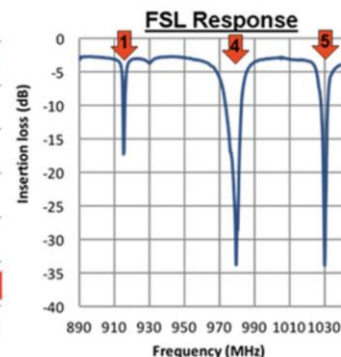
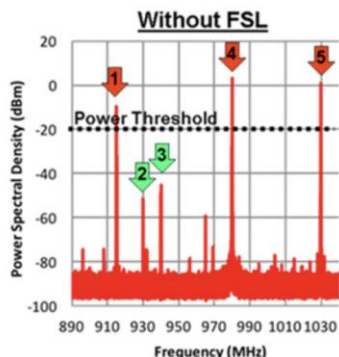
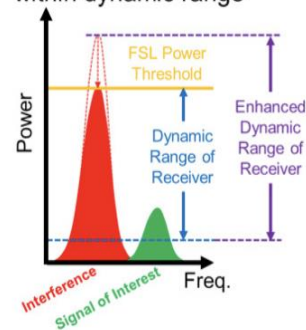
**Auto-tune filters** to mitigate electromagnetic interference (Metamagnetics Inc.)

⇒ Highly specialized, nonlinear spin-waves to limit signal power in every frequency channel (“Frequency selective limiters”, FSL)

**Without FSL:** Small signal outside dynamic range



**With FSL:** Interferer is attenuated, small signal within dynamic range



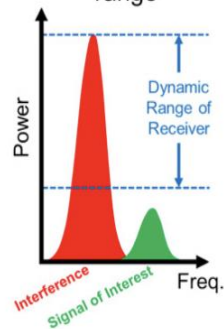
S. M. Gillette, et al. 2018 IEEE MTT-S Int. Microw. Work. Ser. Adv. Mater. Process. RF THz Appl. (IMWS-AMP) 00, 1 (2018).

How does this nonlinear limiter work?  
They don't tell you ;-)

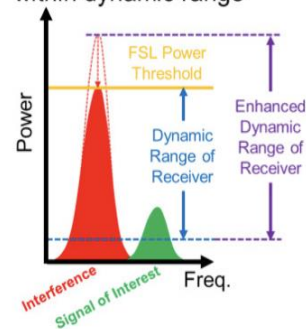
My guess:

- 1) If the RF frequency component gets too strong (the “attacking/jamming” signal), a **magnon instability** is triggered because the threshold for this process is triggered. The energy is flowing then from the initial mode to modes at other frequencies, which dissipate in several channels.
- 2) For large spin-wave powers, the **magnon impedance** is changed due to the nonlinearity. This leads to an **increase of RF reflection**, decreasing the transmitted signal.

**Without FSL:** Small signal outside dynamic range



**With FSL:** Interferer is attenuated, small signal within dynamic range



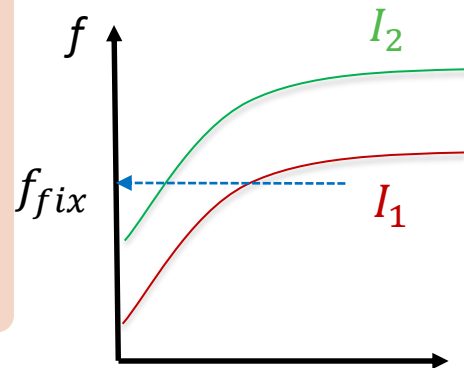
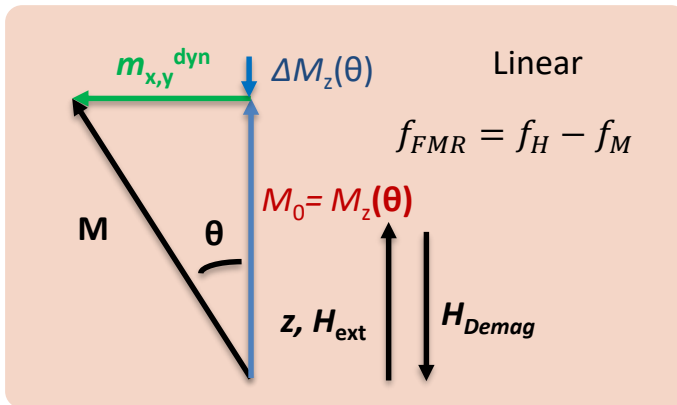
# Nonlinear devices based on the nonlinear magnon shift

Nonlinear shift:

SW amplitude

$$f(k) = \underbrace{f_0(k)}_{\text{linear}} + \underbrace{T_k |c_k|^2}_{\text{nonlinear}}$$

Out-of-plane magnetized film



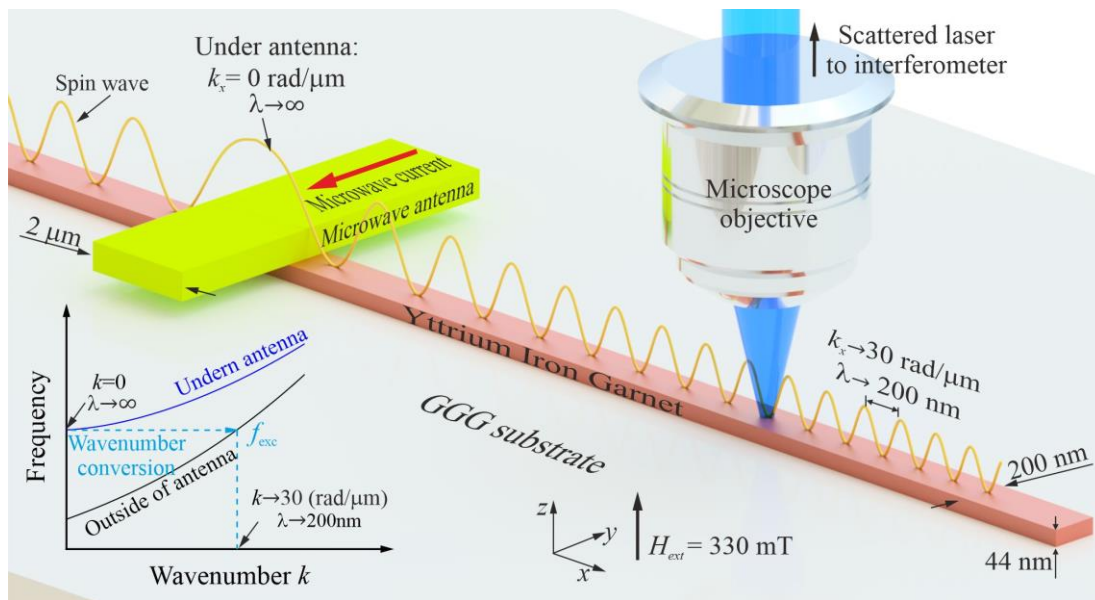
$I_2 \gg I_1$

Including anisotropies:  $T_k \propto M_{eff}$

YIG:  $M_{eff} \approx M_S$

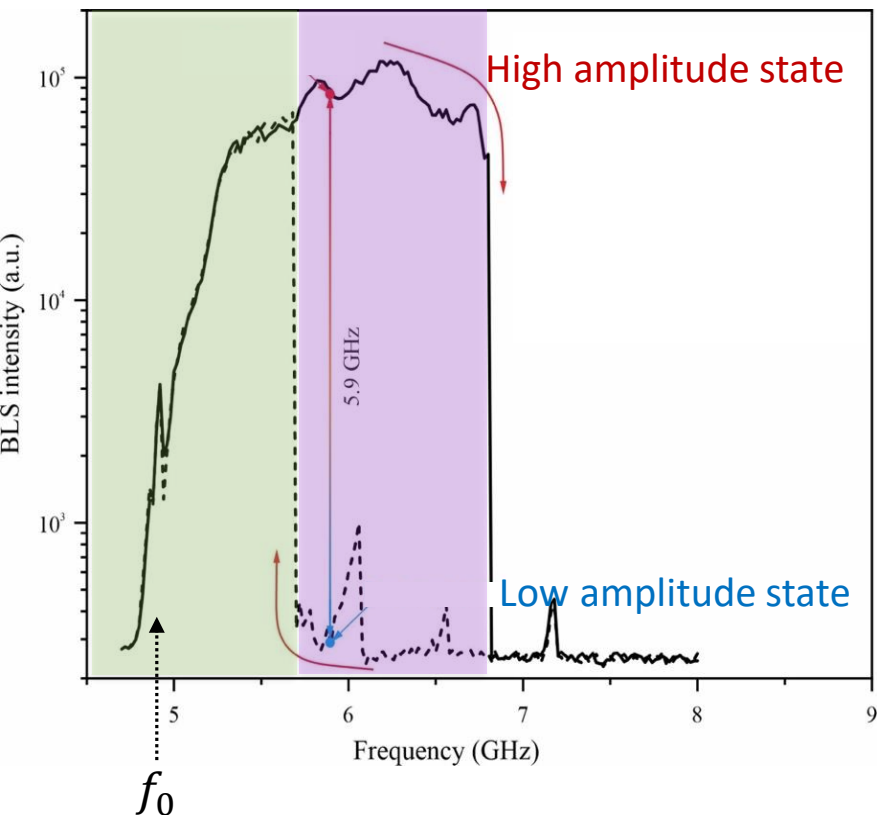
„ $T_k > 0$  if the system is saturated in a hard axis”

$M_{eff} = M_S - K_u$ ,  $K_u$ : Uniaxial anisotropy out-of-

Nonlinear magnonic system A: YIG waveguide in *Forward Volume Geometry*

- Pure YIG (out-of-plane saturated): Positive nonlinear shift coefficient  $T_k$  ( $M_{\text{eff}} > 0$ )
- Large antenna that excites only  $k \approx 0$  in the linear case

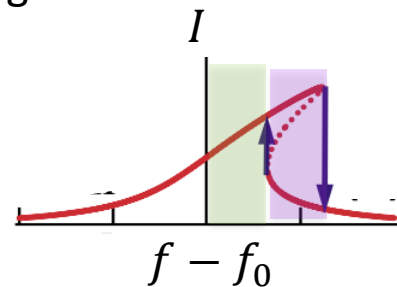
Q. Wang, PP, et al., Sci. Adv. 9, eadg4609 (2023).



- Very wide antenna: linear excitation only at  $f(k \approx 0) = f_0$ .
- Large MW-photon amplitude (high amplitude of dynamic magnetic fields)

Two nonlinear excitation regions:

- Single state foldover
- Bistable foldover

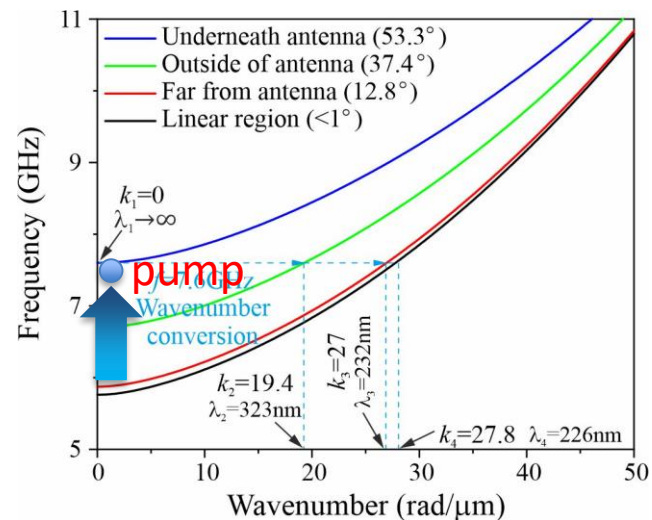
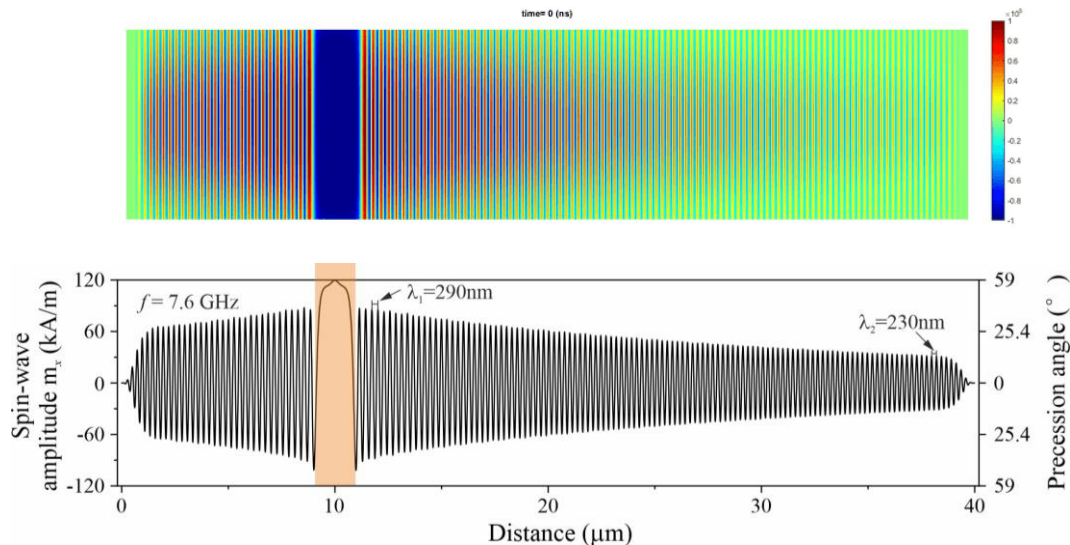


Relative size of the regions is strongly system and power dependent.

see: Q. Wang, PP, et al., *Sci. Adv.* **9**, eadg4609 (2023).

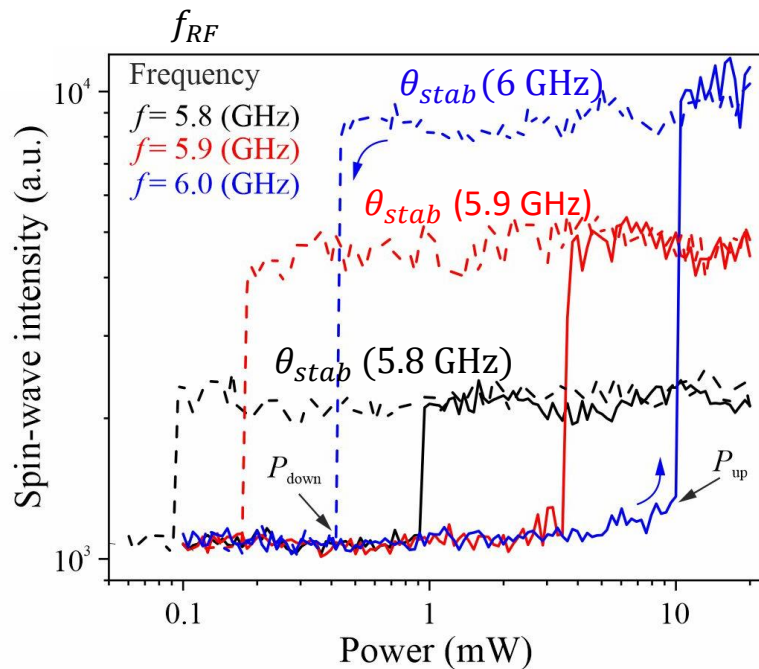


Micromagnetic simulations  
(confirmed by time resolve  $\mu$ BLS)

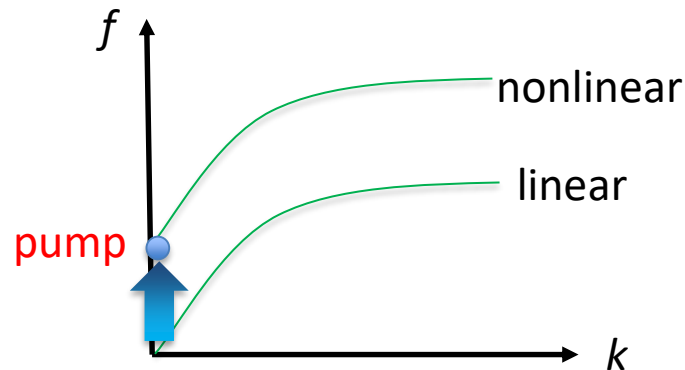


- Spin waves with **230 nm wavelength** (12 % of antenna width) are excited by the nonlinear transformation
- Nonlinear emission** is about **260% more efficient** compared to a linear transformation in a simple field gradient

Q. Wang, PP, et al., *Sci. Adv.* **9**, eadg4609 (2023).



Output spin-wave intensity is **independent** of the input power



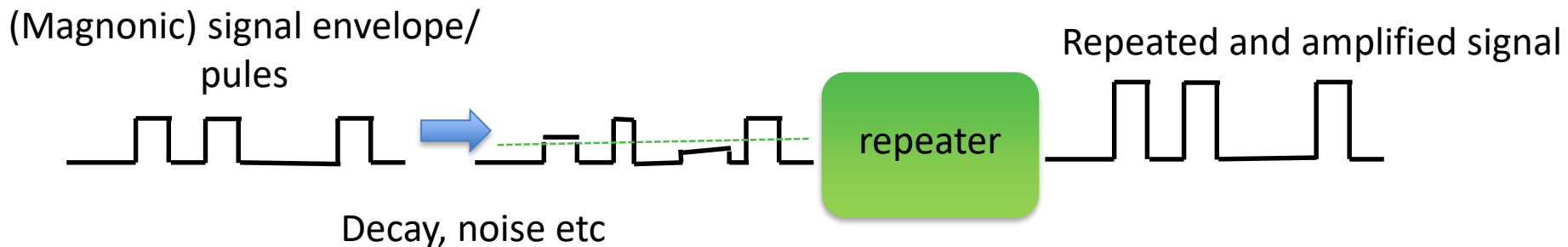
Nonlinearly shifted resonance frequency:

$$f(|c_k|^2) = f_{lin} + T_k \cdot |c_k|^2 \quad |c_k|^2 = 1 - \cos^2(\theta)$$

Condition of stable high amplitude state:

$$\Rightarrow \cos(\theta_{stab}) = 1 - (f_{RF} - f_0)/T_k$$

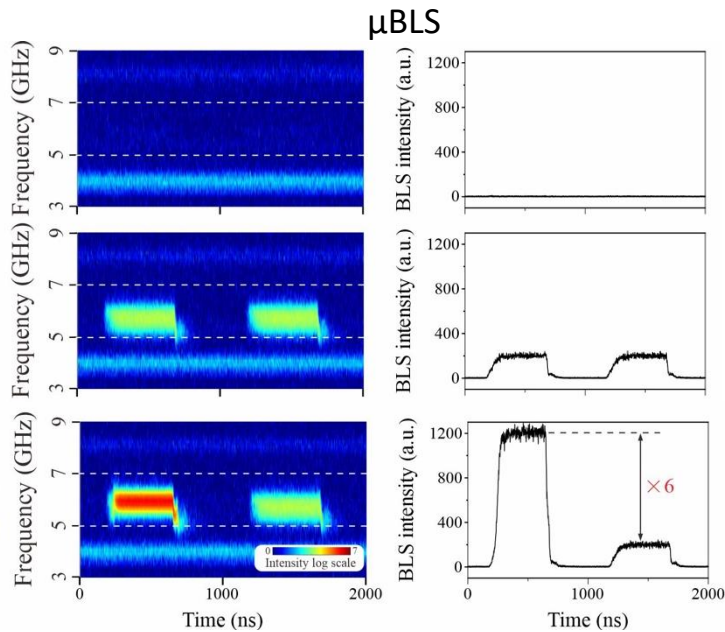
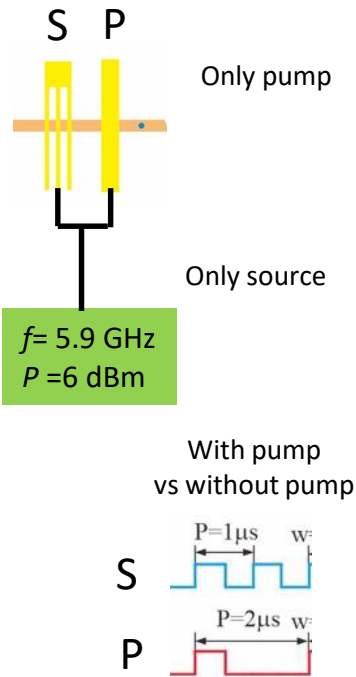
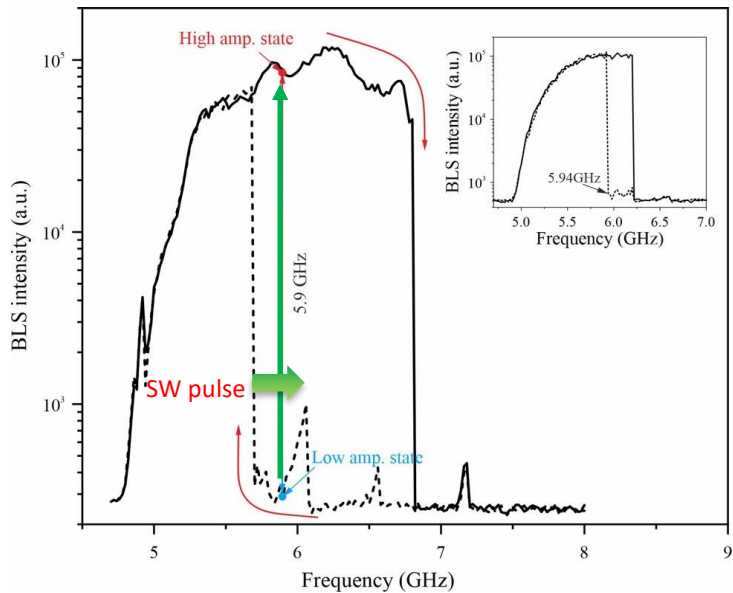
Q. Wang, PP, et al., *Sci. Adv.* **9**, eadg4609 (2023).



Q. Wang et al., *All-Magnonic Repeater Based on Bistability*, Nat. Commun. **15**, 7577 (2024).

# Magnonic repeater based on bistability

YIG waveguide, 1  $\mu\text{m}$  wide, Forward volume (330 mT)



“Trigger transition to high amplitude state using a SW packet.”

- Repeats spin-wave pulse with 6x intensity amplification

Q. Wang et al., *All-Magnonic Repeater Based on Bistability*, Nat. Commun. **15**, 7577 (2024).

- Spin waves can be used to **delay** and **filter Radio-Frequency signals**
- Spin-wave devices are **tunable** by the magnetic field  
=> replace several static devices by one flexible device
- Spin waves can act as **RF Isolators** due to their intrinsic non-reciprocities
- Due to **intrinsic nonlinearities**, SWs can play an important role for **advanced nonlinear RF processing** like signal-to-noise enhancement and frequency selective limiters.