

# Magnetic moments, dipoles and fields

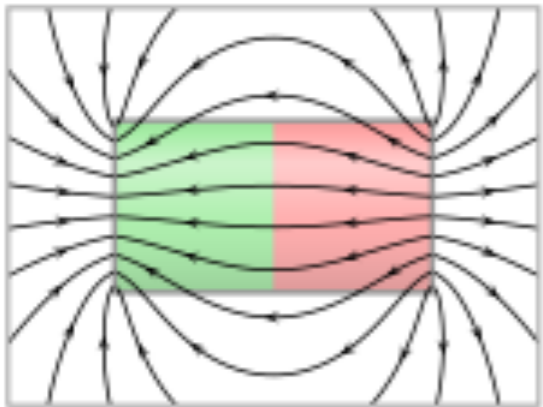
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ESM 2024

# Overview



Origin of magnetic moments



Magnetic fields and demagnetising factors

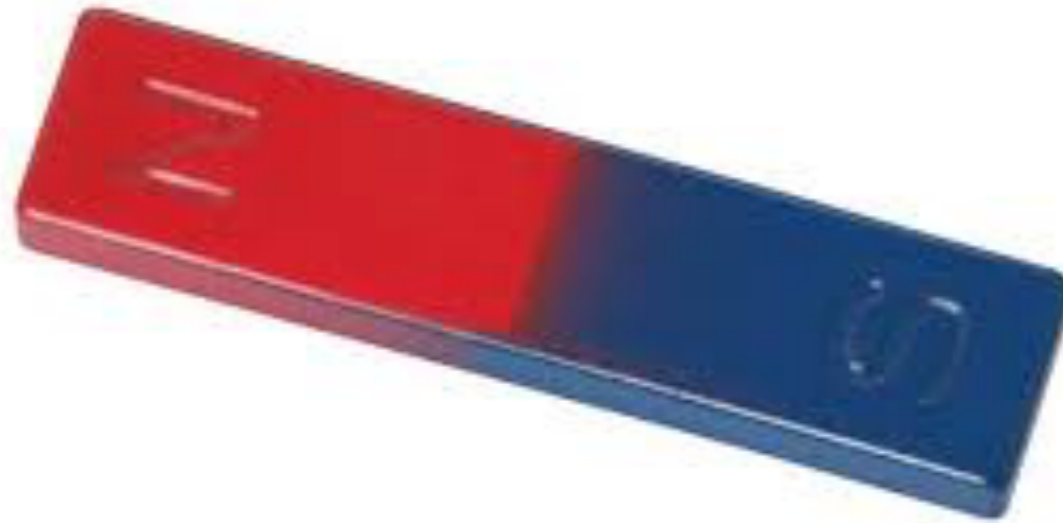
Units in magnetism

# Useful References

- J. M. D. Coey; Magnetism and Magnetic Materials. Cambridge University Press (2010) 614 pp
- Stephen Blundell Magnetism in Condensed Matter, Oxford 2001
- D. C. Jiles An Introduction to Magnetism and Magnetic Materials, CRC Press 480 pp
- J. D. Jackson Classical Electrodynamics 3rd ed, Wiley, New York 1998

# Magnetic moments

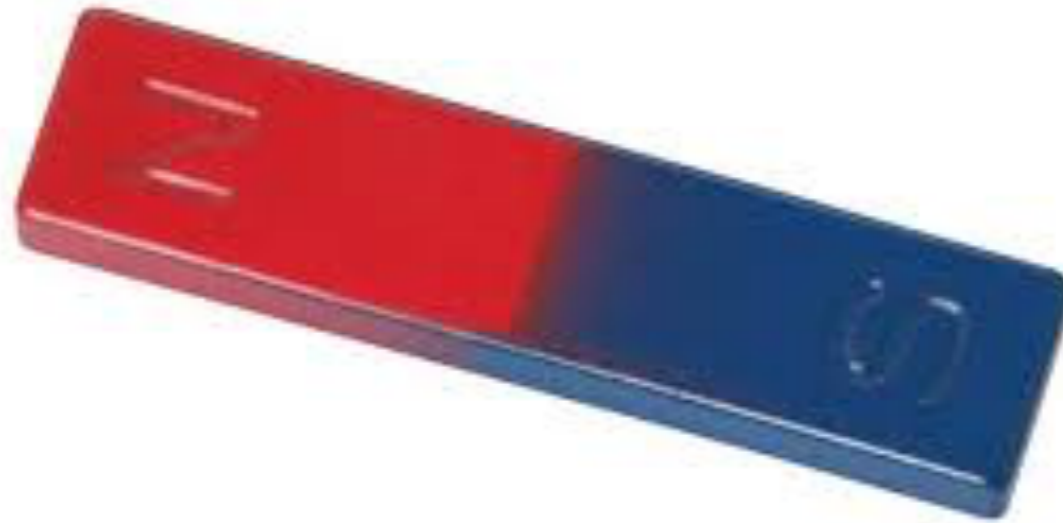
# What is a magnet?



*“A magnet is a material or object that produces a magnetic field”*

**Wikipedia**

# What is a magnet?

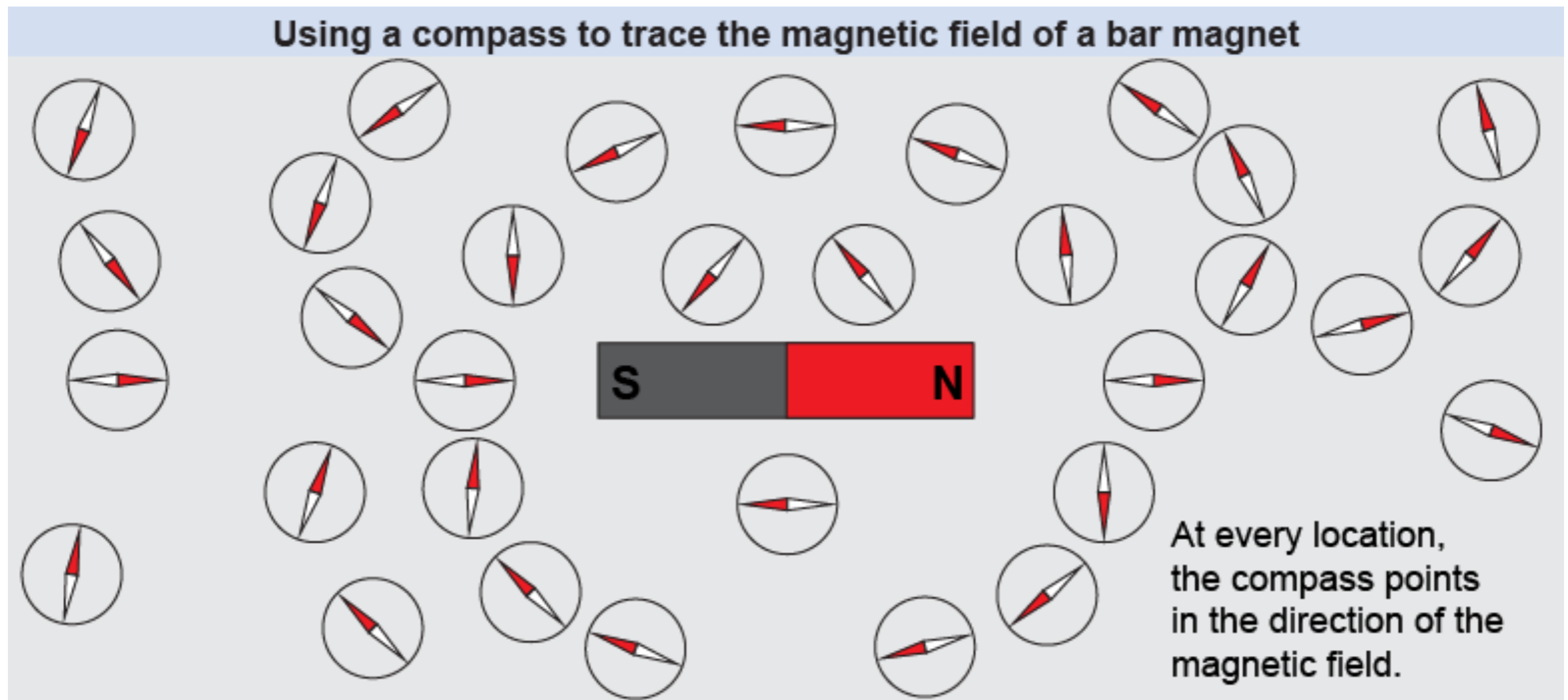


*“A magnet is a material or **object** that produces a magnetic field”*

**Wikipedia**

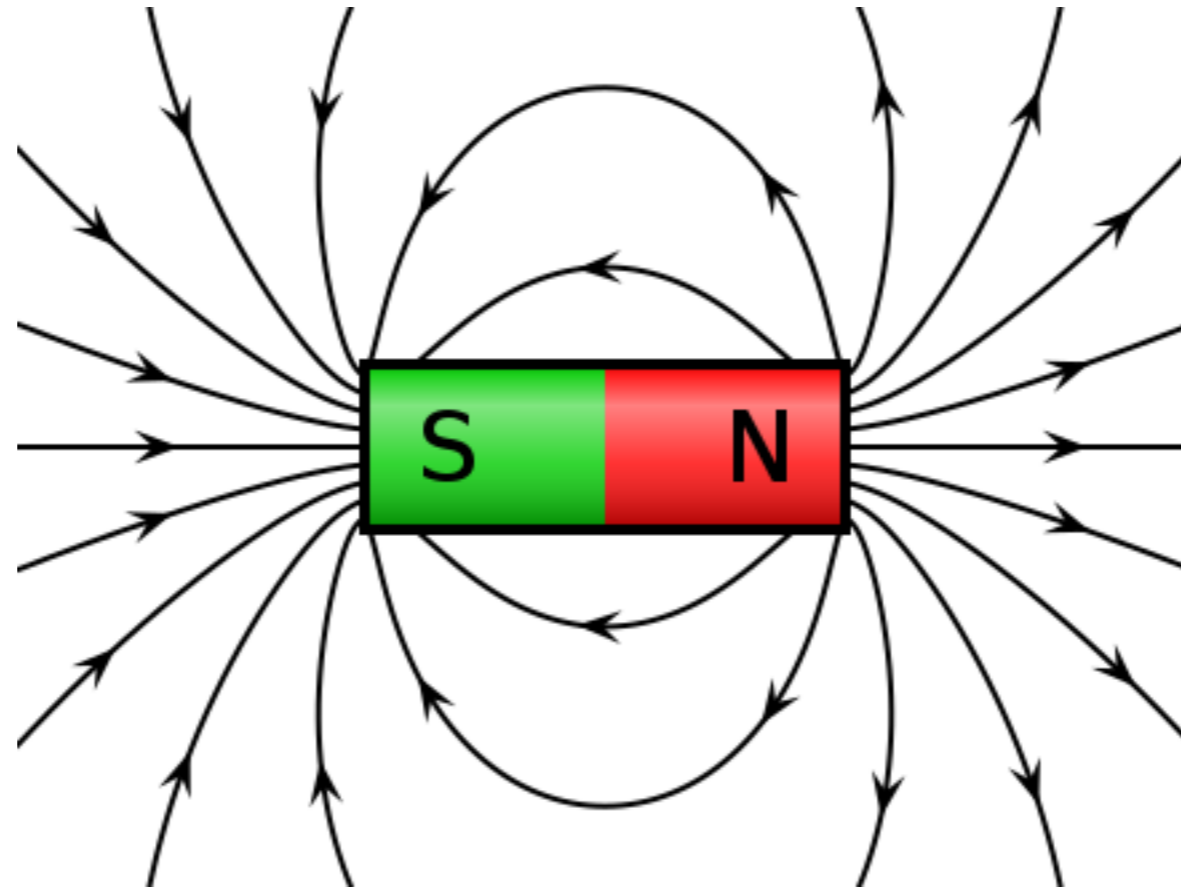
# What is a magnetic field?

- An invisible vector field that interacts with other **magnets**



# What is a magnetic field?

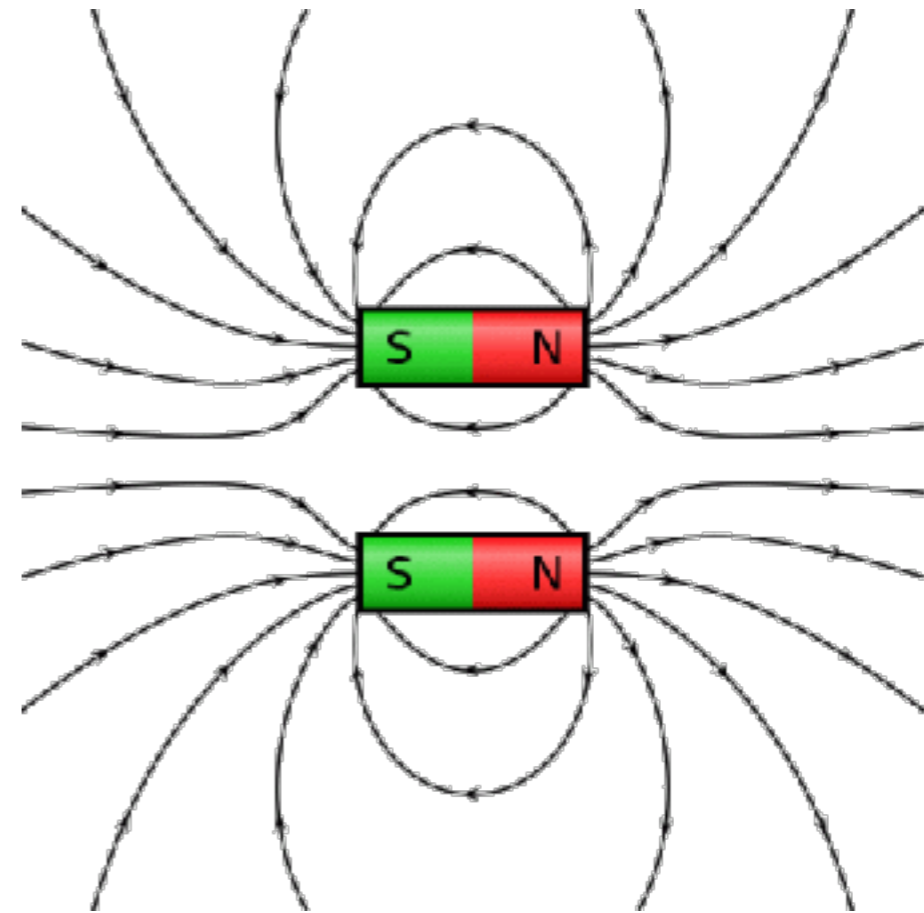
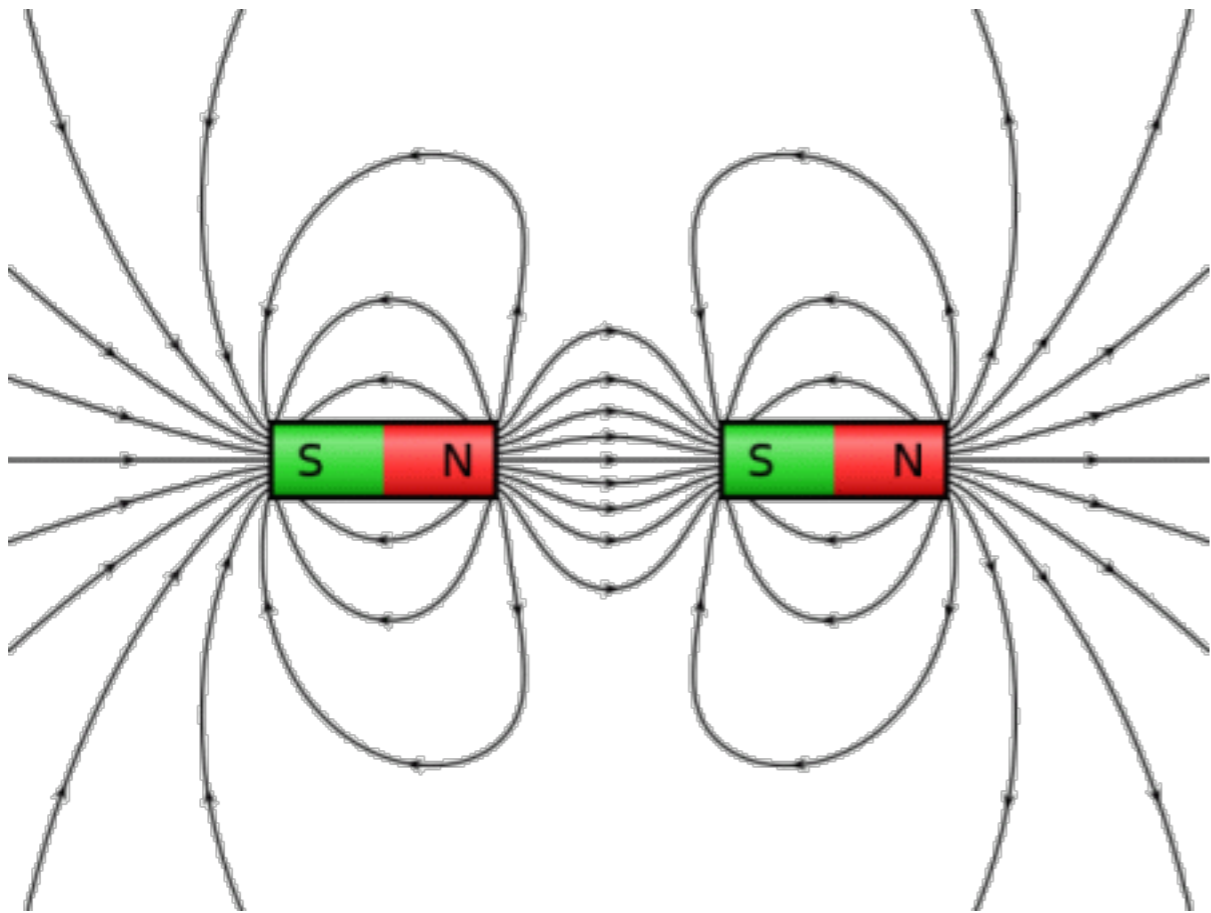
- An invisible vector field that interacts with other **magnets**





# What is a magnetic field?

- An invisible vector field that interacts with other **magnets**

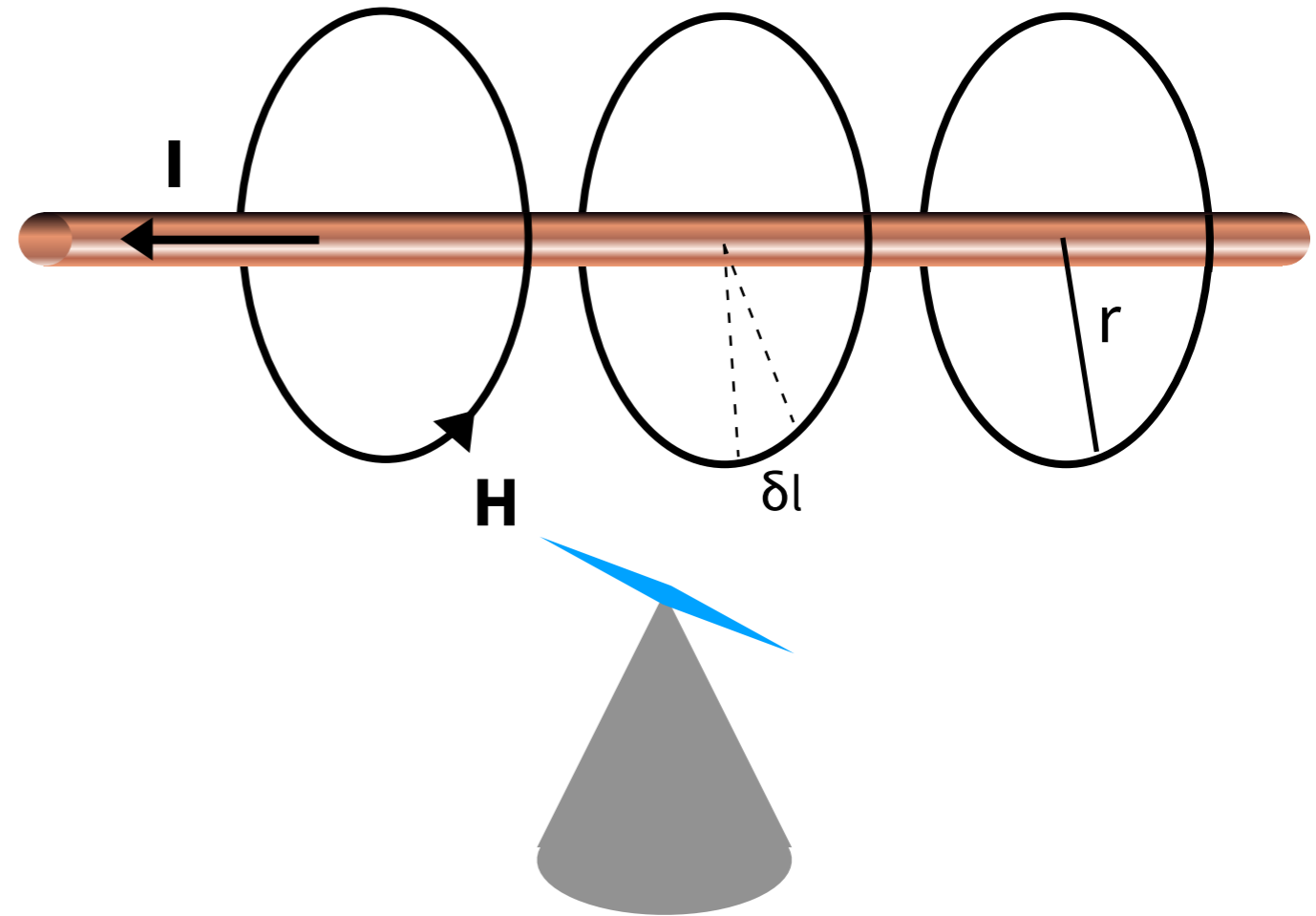


# Magnetic field, Øersted 1820

- Øersted discovered in 1820 that a current carrying wire was able to rotate a compass needle
- Current and field are related by Ampere's Law

$$I = \oint \mathbf{H} dl$$

- Example for  $I = 1\text{A}$ , integral around the loop is  $2\pi r$ ,  $r = 2\text{ mm}$   
 $H \sim 80\text{ A/m}$
- Earth's magnetic field  $\sim 40\text{ A/m}$

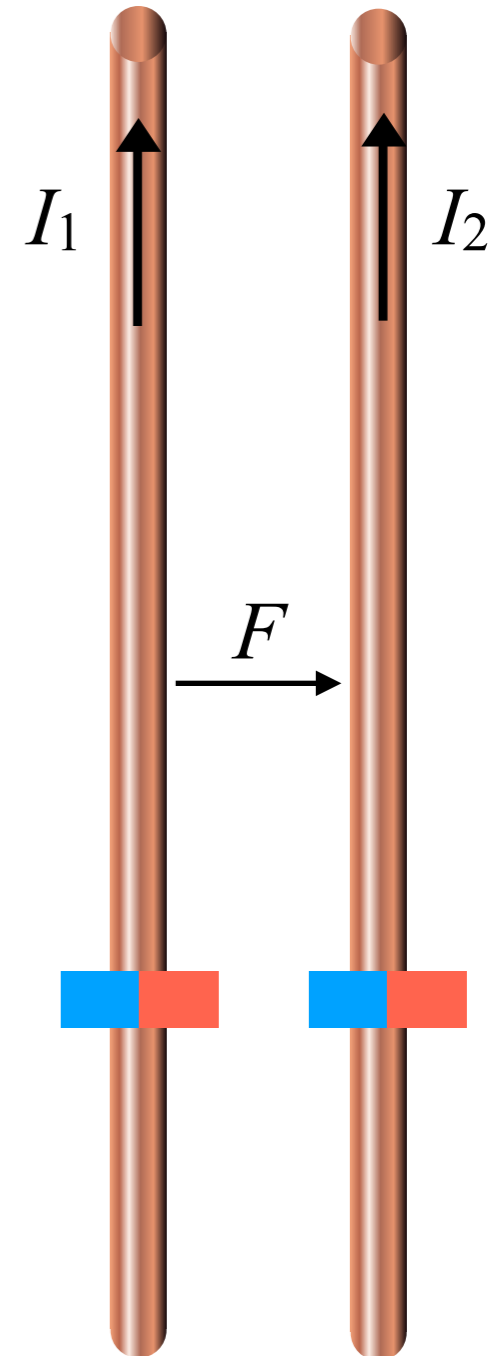


# Interaction of two current-carrying wires, Ampere 1825

- Two current carrying wires (one longer than the other) are attracted to each other for parallel current, and repel for anti-parallel current.

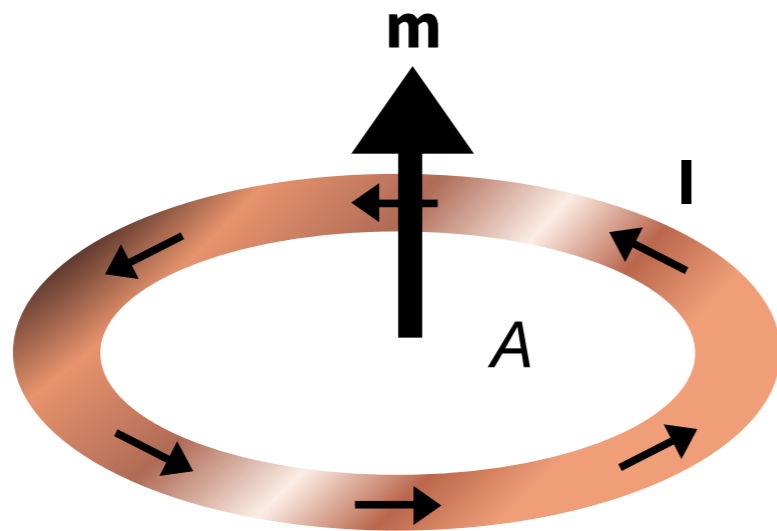
$$\frac{F}{l} = \frac{\mu_0}{2\pi} \frac{I_1 I_2}{r}$$

- The parallel wires “look like” magnets in the **perpendicular** direction
- Weird but central to electromagnetism (E and B fields in light)
- Different from electrostatics as this is a **dynamic** effect from the motion of charge



# Equivalence of currents and magnetic moments

- So currents look like magnets... do magnets look like currents?



$$\mathbf{m} = I_{\perp} A$$

- Can express a current loop as an effective moment, ie a source of magnetic field
- What kind of currents do we need compared to typical magnetic fields?

# Comparison of current magnitudes and magnets

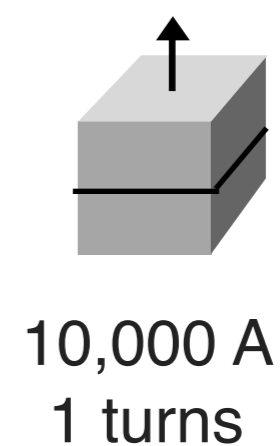
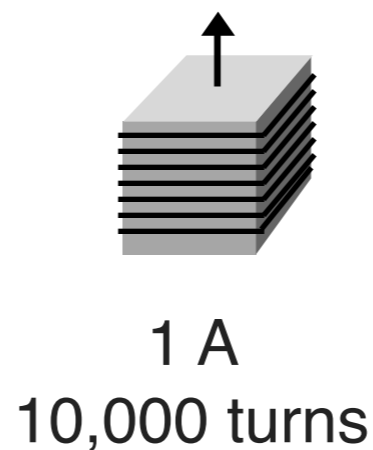
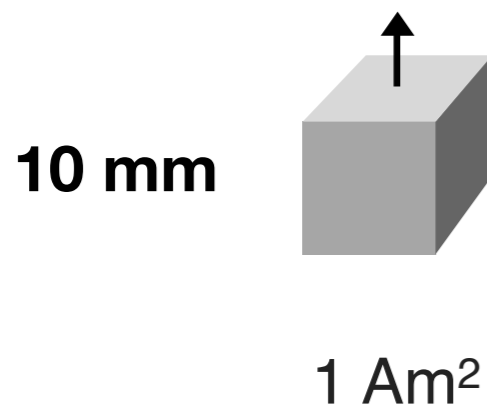
- Using the equivalence of current loops and magnetic moments we can compare the effective currents for a typical small magnet
- Moment given by for a single loop and a solenoid respectively, where  $n$  is the number of turns of the coil

$$\mathbf{m} = I_{\perp} A$$

$$\mathbf{m} = n I_{\perp} A$$



- For a small magnet



- At small sizes, magnets generate much larger fields -> applications in motors

# Difference between magnetic moment and magnetisation

- Magnetic moment is specific to the sample (bigger magnet, bigger field)

0.027 Am<sup>2</sup>



**3 mm**

1 Am<sup>2</sup>



**10 mm**

15.6 Am<sup>2</sup>



**25 mm**

- Magnetization is the **moment density**

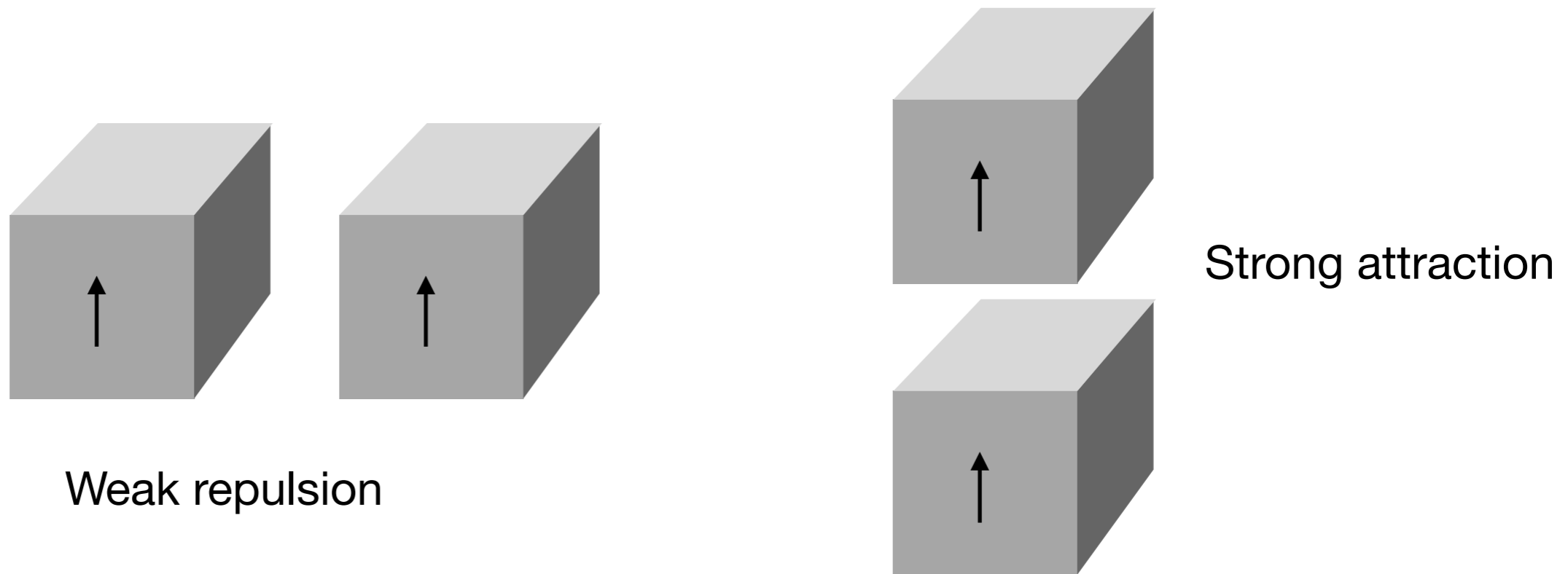
$$m = MV$$

Assume NdFeB  
 $M_s \sim 1 \text{ MA/m}$

- Magnetisation is a property of the **material**
- Moment is a property of a **magnet**
- Magnetisation is scale independent

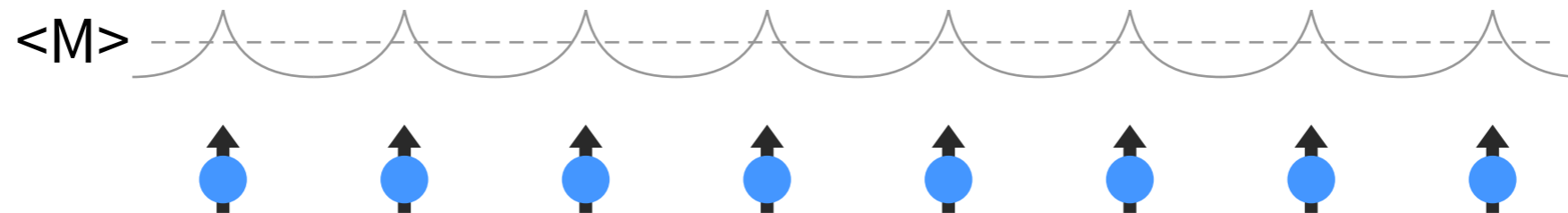
# Vectorial nature of magnetic moments

- A magnetic moment generates a field around it
- Interaction with non-magnets is weak
- Interaction with magnets is stronger but orientation dependent



# Physical origin of magnetization and magnetic moment

- At the atomic scale the magnetic moments fluctuate strongly in time and space due to the electrons 'orbiting' nuclei



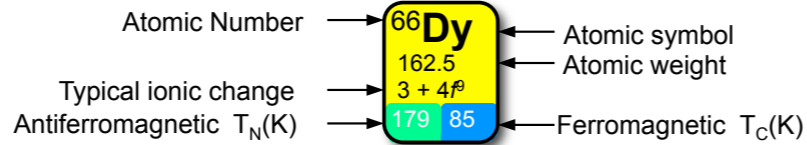
- Use a continuous medium approximation to calculate an average magnetisation  $\langle M \rangle$  (moment/volume)
- Avoids all the horrible details of fluctuating moments and can treat magnetism on a **continuum** level
- Good approximation for ferromagnets for volumes much larger than the atomic volume



# Which elements are magnetic

## Magnetic Periodic Table

1 H 1.00																	2 He 4.00
3 Li 6.94 1 + 2s <sup>0</sup>	4 Be 9.01 2 + 2s <sup>0</sup>											5 B 10.81	6 C 12.01	7 N 14.01	8 O 16.00 35	9 F 19.00	10 Ne 20.18
11 Na 22.99 1 + 3s <sup>0</sup>	12 Mg 24.21 2 + 3s <sup>0</sup>											13 Al 26.98 3 + 2p <sup>6</sup>	14 Si 28.09	15 P 30.97	16 S 32.07	17 Cl 35.45	18 Ar 39.95
19 K 38.21 1 + 4s <sup>0</sup>	20 Ca 40.08 2 + 4s <sup>0</sup>	21 Sc 44.96 3 + 3d <sup>0</sup>	22 Ti 47.88 4 + 3d <sup>0</sup>	23 V 50.94 3 + 3d <sup>2</sup>	24 Cr 52.00 3 + 3d <sup>5</sup> 312	25 Mn 55.85 2 + 3d <sup>5</sup> 96	26 Fe 55.85 3 + 3d <sup>6</sup> 1043	27 Co 58.93 2 + 3d <sup>7</sup> 1390	28 Ni 58.69 2 + 3d <sup>8</sup> 629	29 Cu 63.55 2 + 3d <sup>9</sup>	30 Zn 65.39 2 + 3d <sup>10</sup>	31 Ga 69.72 3 + 3d <sup>10</sup>	32 Ge 72.61	33 As 74.92	34 Se 78.96	35 Br 79.90	36 Kr 83.80
37 Rb 85.47 1 + 5s <sup>0</sup>	38 Sr 87.62 2 + 5s <sup>0</sup>	39 Y 88.91 2 + 4d <sup>0</sup>	40 Zr 91.22 4 + 4d <sup>0</sup>	41 Nb 92.91 5 + 4d <sup>0</sup>	42 Mo 95.94 5 + 4d <sup>1</sup>	43 Tc 97.9	44 Ru 101.1 3 + 4d <sup>6</sup>	45 Rh 102.4 3 + 4d <sup>6</sup>	46 Pd 106.4 2 + 4d <sup>8</sup>	47 Ag 107.9 1 + 4d <sup>10</sup>	48 Cd 112.4 2 + 4d <sup>10</sup>	49 In 114.8 3 + 4d <sup>10</sup>	50 Sn 118.7 4 + 4d <sup>10</sup>	51 Sb 121.8	52 Te 127.6	53 I 126.9	54 Xe 83.80
55 Cs 132.9 1 + 6s <sup>0</sup>	56 Ba 137.3 2 + 6s <sup>0</sup>	57 La 138.9 3 + 4f <sup>0</sup>	72 Hf 178.5 4 + 5d <sup>0</sup>	73 Ta 180.9 5 + 5d <sup>0</sup>	74 W 183.8 6 + 5d <sup>0</sup>	75 Re 186.2 4 + 5d <sup>5</sup>	76 Os 190.2 3 + 5d <sup>6</sup>	77 Ir 192.2 4 + 5d <sup>6</sup>	78 Pt 195.1 2 + 5d <sup>8</sup>	79 Au 197.0 1 + 5d <sup>10</sup>	80 Hg 200.6 2 + 5d <sup>10</sup>	81 Tl 204.4 3 + 5d <sup>10</sup>	82 Pb 207.2 4 + 5d <sup>10</sup>	83 Bi 209.0	84 Po 209	85 At 210	86 Rn 222
87 Fr 223	88 Ra 226.0 2 + 7s <sup>0</sup>	89 Ac 227.0 3 + 5f <sup>0</sup>															
			58 Ce 140.1 4 + 4f <sup>0</sup> 13	59 Pr 140.9 3 + 4f <sup>2</sup>	60 Nd 144.2 3 + 4f <sup>3</sup> 19	61 Pm 145	62 Sm 150.4 3 + 4f <sup>6</sup> 105	63 Eu 152.0 2 + 4f <sup>7</sup> 90	64 Gd 157.3 3 + 4f <sup>7</sup> 292	65 Tb 158.9 3 + 4f <sup>9</sup> 229 221	66 Dy 162.5 3 + 4f <sup>9</sup> 179 85	67 Ho 164.9 3 + 4f <sup>10</sup> 132 20	68 Er 167.3 3 + 4f <sup>11</sup> 85 20	69 Tm 168.9 3 + 4f <sup>12</sup> 56	70 Yb 173.0 3 + 4f <sup>13</sup>	71 Lu 175.0 3 + 4f <sup>14</sup>	
			90 Th 232.0 4 + 5f <sup>0</sup>	91 Pa 231.0 5 + 5f <sup>0</sup>	92 U 238.0 4 + 5f <sup>2</sup>	93 Np 238.0 5 + 5f <sup>2</sup>	94 Pu 244	95 Am 243	96 Cm 247	97 Bk 247	98 Cf 251	99 Es 252	100 Fm 257	101 Md 258	102 No 259	103 Lr 260	

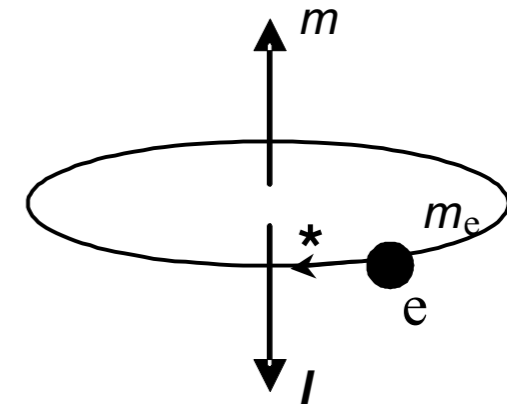


- Nonmetal
- Diamagnet
- Ferromagnet T<sub>C</sub> > 290K
- Metal
- Paramagnet
- Antiferromagnet with T<sub>N</sub> > 290K
- Radioactive
- BOLD** Magnetic atom
- Antiferromagnet/Ferromagnet with T<sub>N</sub>/T<sub>C</sub> < 290 K

From Coey

# Bohr magneton

- Can consider an electron 'orbiting' an atom
- A moving charge looks like a 'current', generating an effective magnetic moment
- In Bohr's quantum theory, orbital angular momentum  $\mathbf{l}$  is quantized in units of  $\hbar$ ;  $h$  is Planck's constant,  $6.6226 \times 10^{-34}$  Js;  $\hbar = h/2\pi = 1.055 \times 10^{-34}$  Js
- The orbital angular momentum is  $\mathbf{l} = m_e \mathbf{r} \wedge \mathbf{v}$
- It is the z-component of  $\mathbf{l}_z$  that is quantized in units of  $\hbar$ , taking a value  $m_l$ .  $m_l$  is a quantum number, an *integer* with no units. Eliminating  $r$  in the expression for  $m$
- $\mu_B$  is the Bohr magneton, the basic unit of atomic magnetism



$$m = IA = -\frac{evr}{2}$$

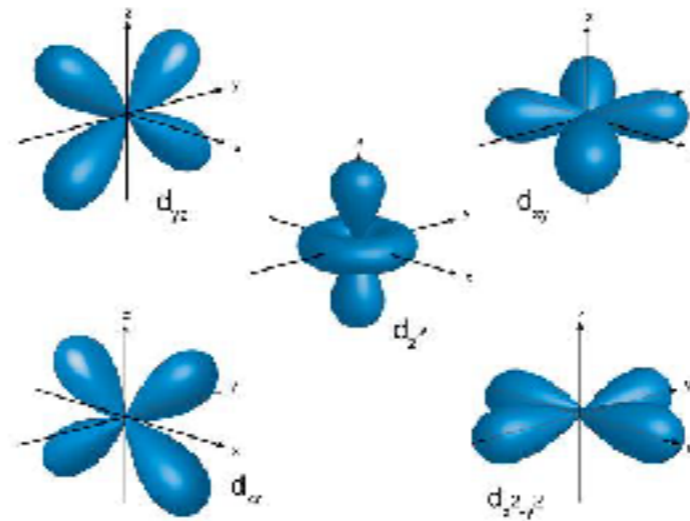
$$m = -\frac{e}{2m_e} \mathbf{l} = \frac{e\hbar}{2m_e} m_l = m_l \mu_B$$

$$\mu_B = \frac{e\hbar}{2m_e} = 9.274 \times 10^{-24} \text{ Am}^2 \text{ | JT}^{-1}$$

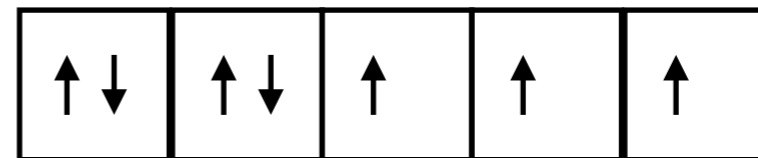
\* electrons travel in the opposite direction to currents

# Non-integer magnetic moments

- Transition metal magnets tend to have non-integer magnetic moments, eg Fe  $\sim 2.2 \mu_B$ , Co  $\sim 1.72 \mu_B$ , Ni  $\sim 0.6 \mu_B$
- If electrons carry quanta of angular momentum, how is this possible?
- Classic explanation is itinerant magnetism: electrons are delocalised and form bands
- First principles calculations reveal a non-integer magnetic moment quite localised to the atom
- Effect due to electrons hopping between different  $d$ -orbitals



		Fe (BCC)
Fe s	↑	0.31
	↓	0.32
p	↑	0.36
	↓	0.40
d	↑	4.37
	↓	<u>2.19</u>
Total	↑	5.07
	↓	2.93

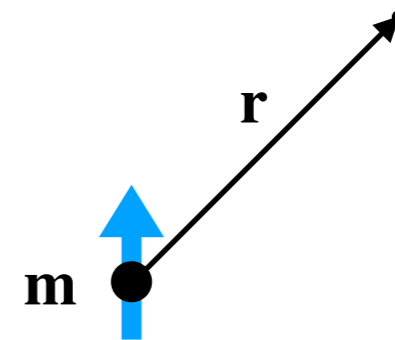


# Field from a dipole

- The magnetic induction (field) from a point dipole can be derived classically (see Jackson) and is given by

$$\mathbf{B} = \frac{\mu_0}{4\pi} \left( \frac{3(\mathbf{m} \cdot \hat{\mathbf{r}})\hat{\mathbf{r}} - \mathbf{m}}{|\mathbf{r}|^3} \right)$$

**B-field at any point  
from a point dipole**



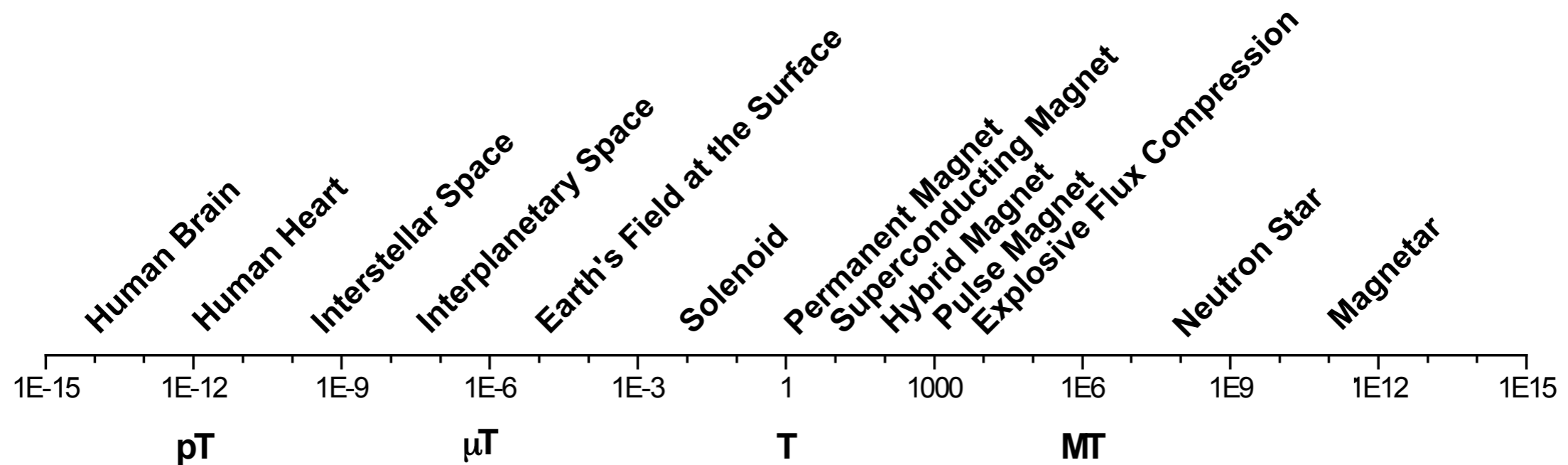
- Ignores any distribution of magnetic ‘charge’ at the dipole (need a multipole description)

J. D. Jackson, Classical electrodynamics (2nd ed.). New York: Wiley. (1975)

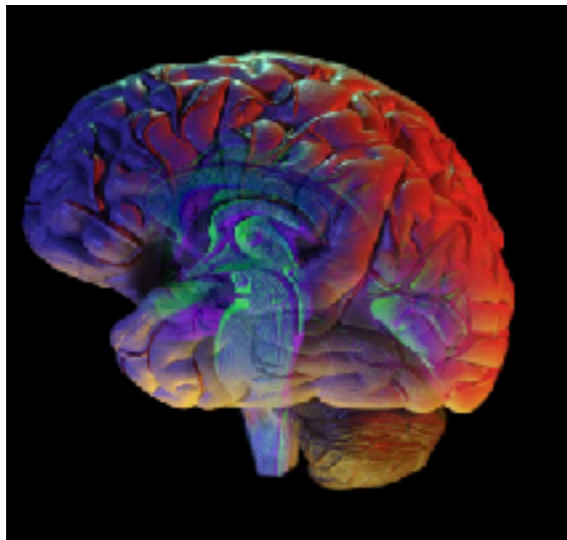
# **Magnetic fields and demagnetising factors**

# What ranges of magnetic fields exist?

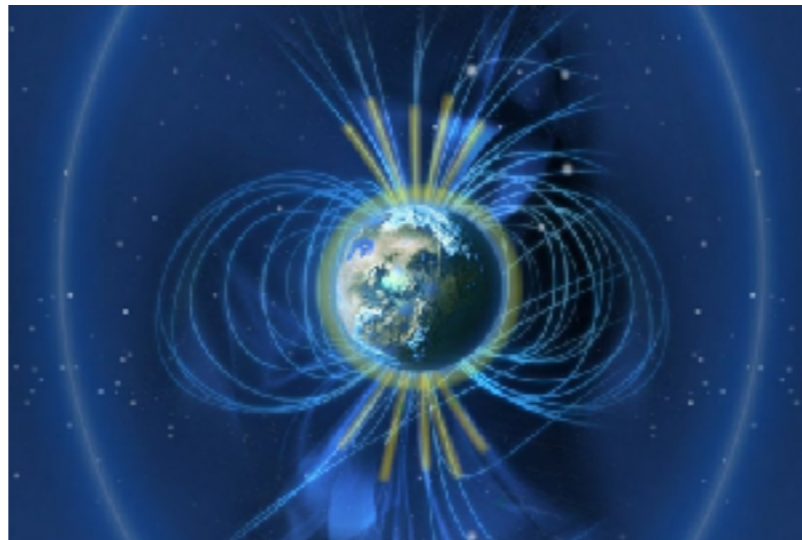
- Historically a terrestrial 1T field was considered 'large'
- Today that is not generally true
  - Recording Media coercivity  $\sim 1\text{T}$
  - MRI  $\sim 5\text{T}$



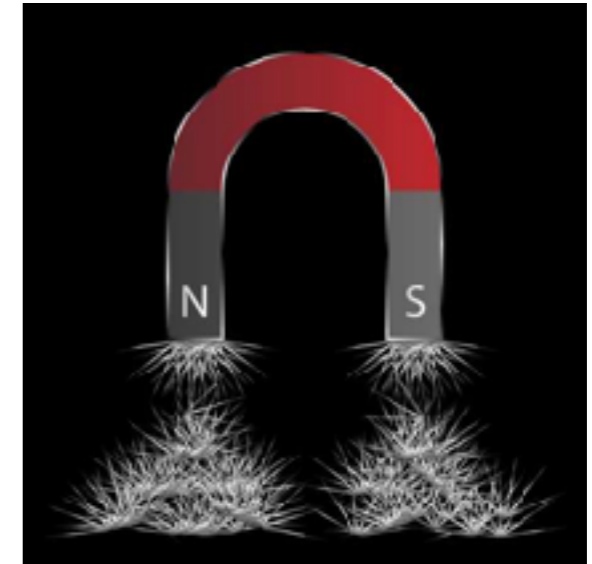
# Typical values of magnetic fields



Human Brain 1 fT



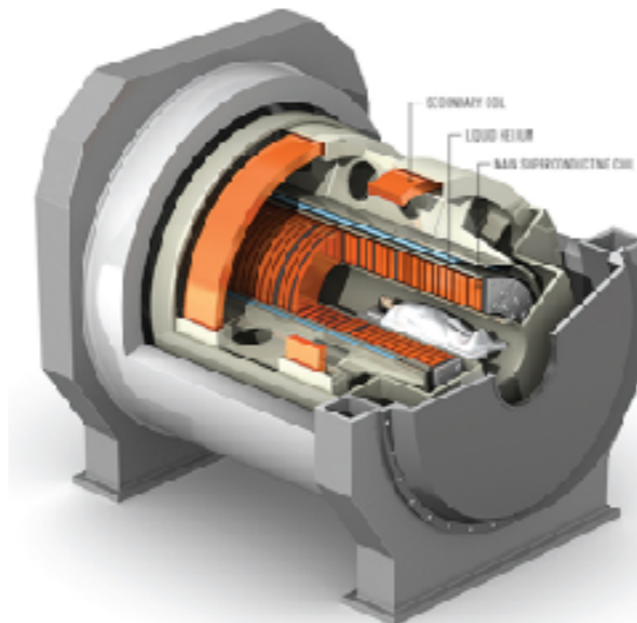
Earth 50  $\mu$ T



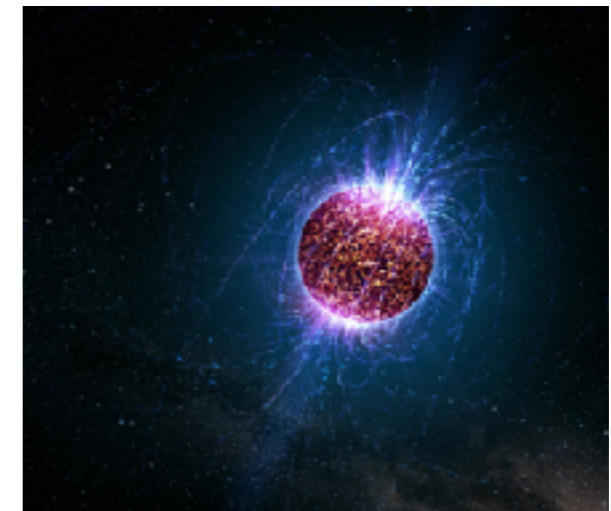
Permanent Magnet 0.5-1T



Electromagnet 1T



Superconducting magnet 10 T



Magnetar  $10^{12}$  T

# Magnetic fields in free space

- Two definitions of magnetic field

Magnetic Field  $\mathbf{H}$  [A/m]

Magnetic flux density  $\mathbf{B}$  [T]

- When talking about generated magnetic fields in free space, they express the exact same physical phenomenon, and are related by

$$\mathbf{B} = \mu_0 \mathbf{H}$$

- $\mu_0 = 4\pi \cdot 10^{-7}$  H/m is the permeability of free space
- The difference between H-field and B-field is a common point of confusion, but only when considering a magnetic medium  $\mathbf{B} = \mu_0(\mathbf{H} + \mathbf{M})$
- **B**-field component arising from **applied H**-field is exactly  $\mathbf{B} = \mu_0 \mathbf{H}$



# Magnetic fields in media

- The actual B-field in response to media is generally more complex

$$\mathbf{B} = \mu_0 (\mathbf{M} + \mathbf{H})$$

- Or alternatively in terms of a relative permeability or susceptibility

$$\mathbf{B} = \mu_0 \mu_r \mathbf{H} = \mu_0 (1 + \chi) \mathbf{H}$$

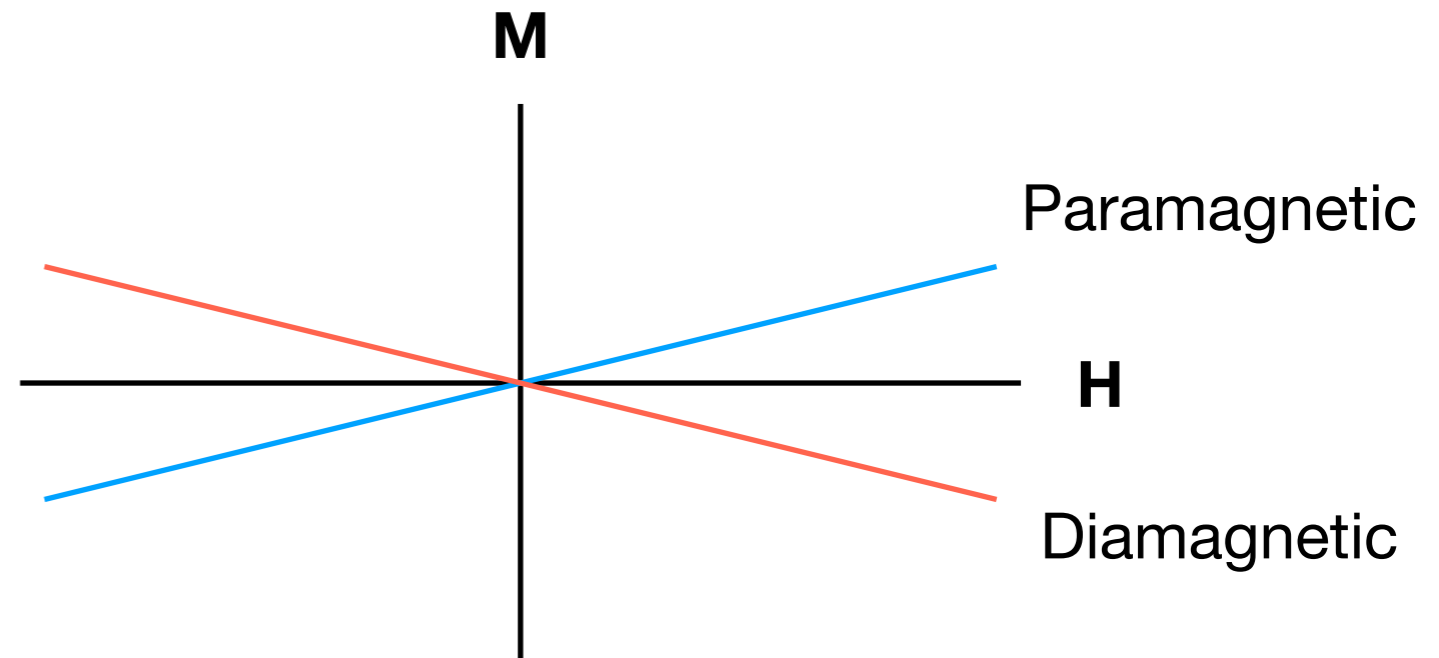
- where the susceptibility gives the full magnetic response, or limit of small fields (initial susceptibility)

$$\chi = \mathbf{M}(\mathbf{H}) = \left. \frac{d\mathbf{M}}{d\mathbf{H}} \right|_{H \rightarrow 0}$$

- Different media ave very different responses, ferromagnets highly non-linear

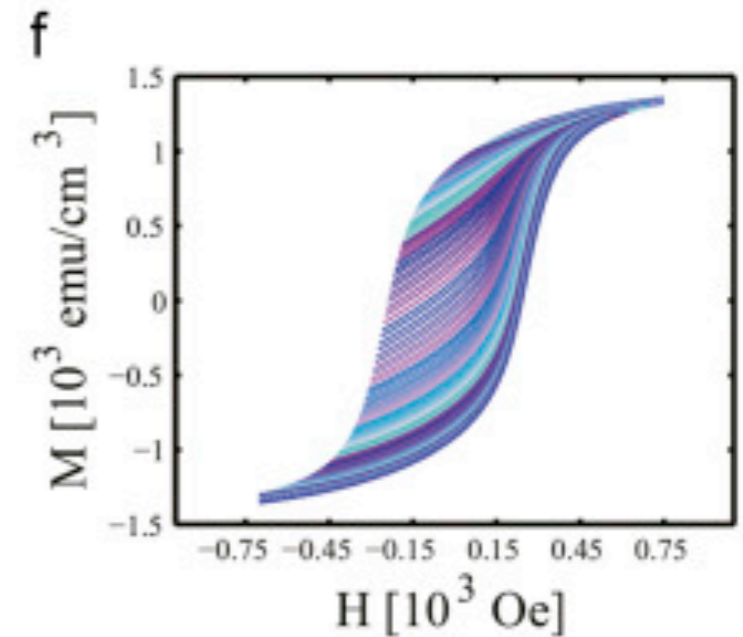
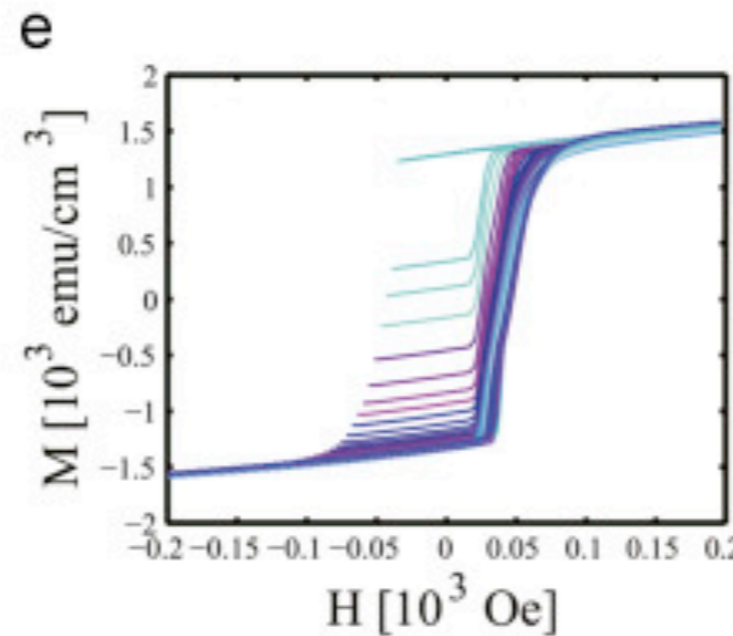
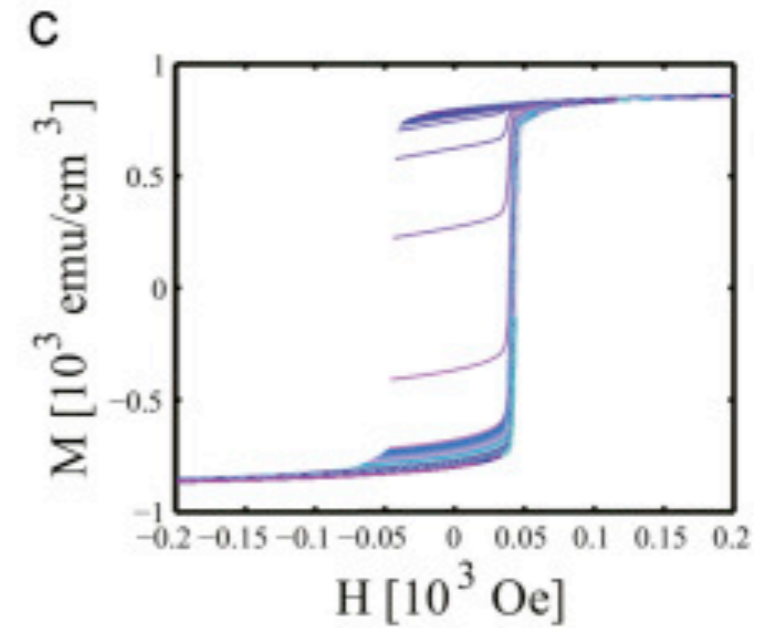
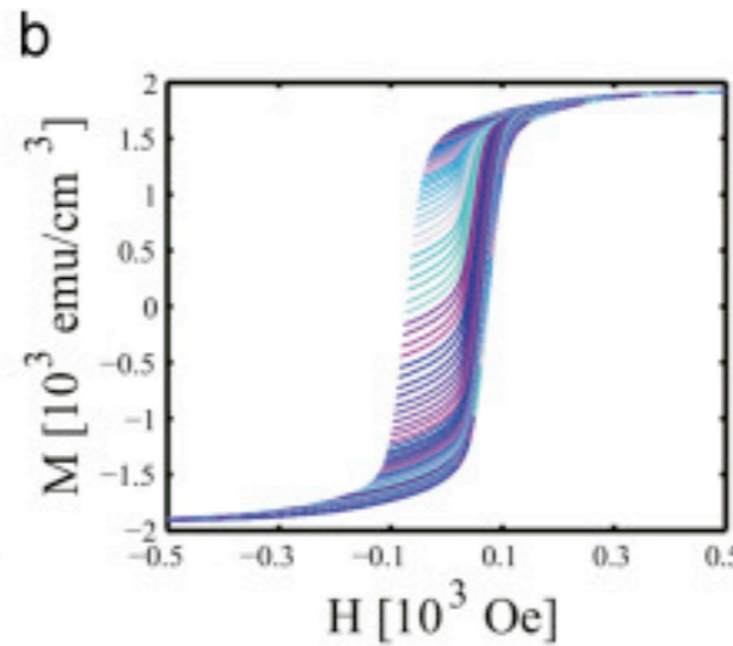
# Diamagnetism and Paramagnetism

- Diamagnets and paramagnets have a weak magnetic response ( $\chi \ll 1$ ),  $\sim 10^{-4} - 10^{-6}$
- Response typically isotropic with respect to the field
- Diamagnets repel external magnetic fields due to Larmor precession of bound electrons that induces a moment opposite to the applied field
- Paramagnets weakly align with an external field overcoming thermal fluctuations



# Ferromagnetism

- Complex and anisotropic behaviour of  $\mathbf{M}(\mathbf{H})$
- Definition of  $\chi = M/H$  is not very sensible in most cases
- Saturated case easier to deal with!



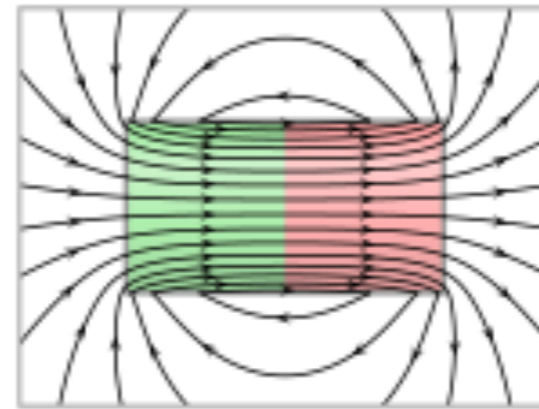
# Relation between $\mathbf{B}$ and $\mathbf{H}$ in a saturated material

- Magnetic field around a saturated magnet  
simple  $\mathbf{B} = \mu_0 \mathbf{H}$

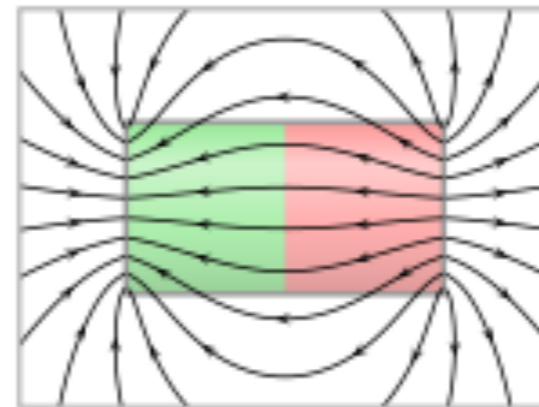
- What about inside the magnet?

$$\mathbf{B} = \mu_0 (\mathbf{M} + \mathbf{H})$$

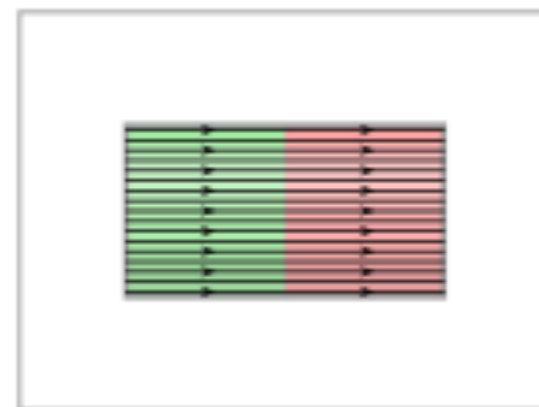
- Why do we care?
- In general magnetization processes are anisotropic and depend on sample shape



$\vec{B}$



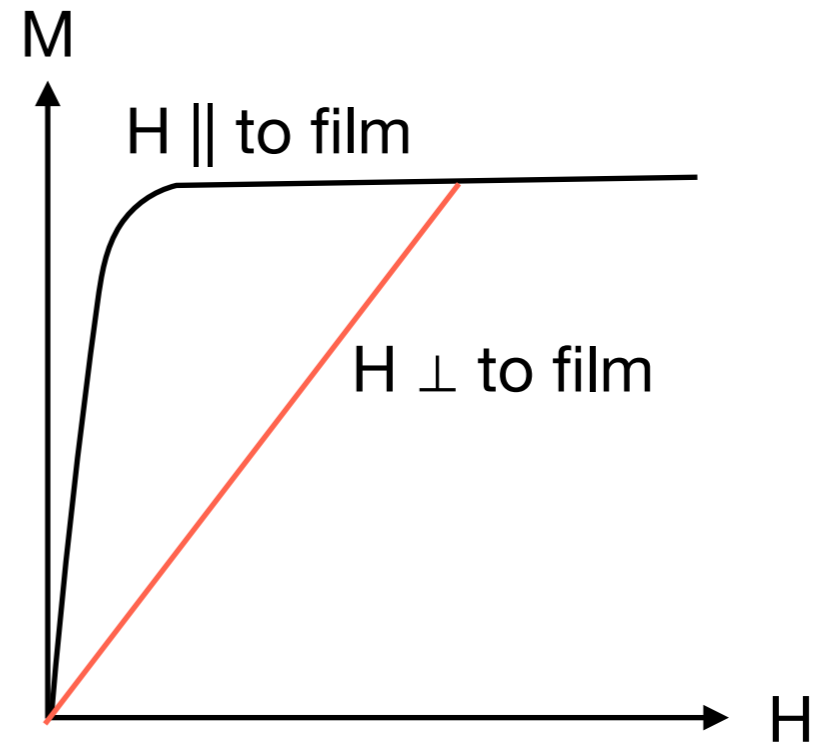
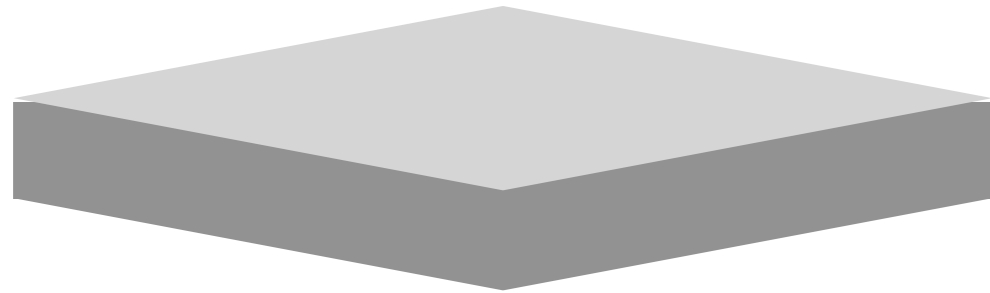
$\vec{H}$



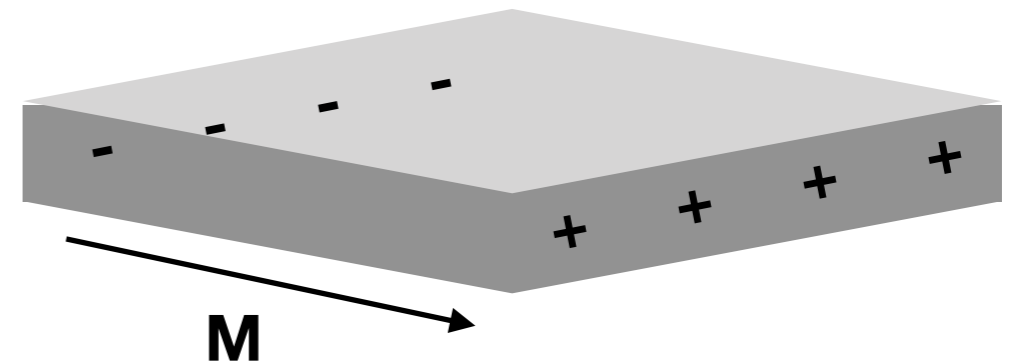
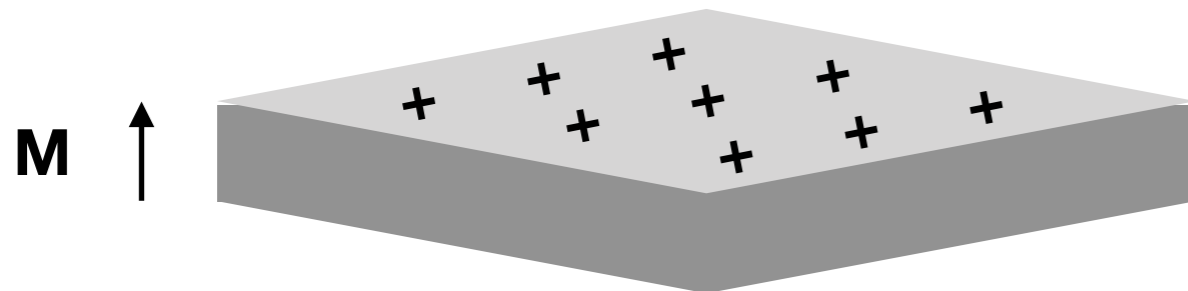
$\vec{M}$

# Example: thin magnetic film

- Much easier to magnetise in the plane than out-of-plane



- Origin is *demagnetising* field - aims to minimise surface charges



# Demagnetizing fields

- Local effective field inside the magnet depends on surface

$$\mathbf{H}(\mathbf{r}) = \frac{1}{4\pi} \left\{ - \int_V \frac{(\nabla \cdot \mathbf{M})(\mathbf{r} - \mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|^3} d^3r' + \int_S \frac{\mathbf{M} \cdot \mathbf{e}_n(\mathbf{r} - \mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|^3} d^2r' \right\}$$

- Since  $\mathbf{M}$  is uniform, first (bulk) term is zero
- For surface term,  $\mathbf{M} \cdot \mathbf{e}_n$  determines surface charge density, larger surface leads to larger field opposing magnetisation
- Leads to concept of a *demagnetising field*

$$\mathbf{H} = \mathbf{H}_{\text{app}} - \mathbf{H}_d$$

# Demagnetization Factor

- Calculating demagnetisation field is tedious (lots of boring and complicated integrals)
- Simplify - invent a “demagnetising factor” or “shape factor”  $N$

$$\mathbf{H}_d = -N\mathbf{M}$$

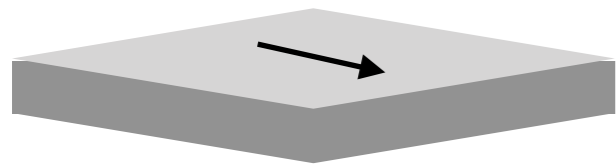
- Shape factor gives a constant of proportionality between the demagnetising field and shape
- Always between 0-1 and in general a tensor with trace 1

$$N_x + N_y + N_z = 1$$

- Known for simple geometric shapes (spheres, ellipsoids, rectangular prisms)
- Is usually calculated numerically for anything complicated

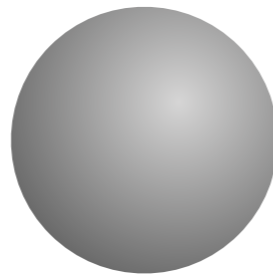
# Demagnetization factors for different shapes

$$N = 0$$



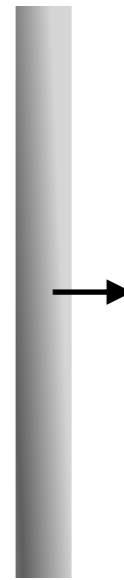
Infinite thin film

$$N = 1/3$$



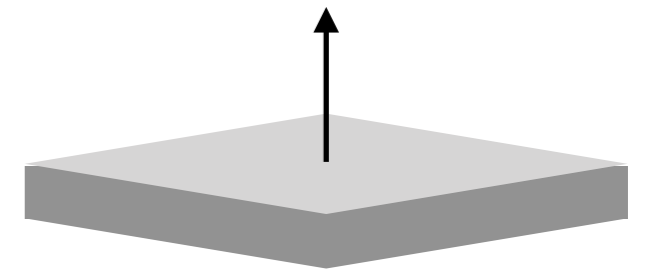
Sphere

$$N = 1/2$$



Infinitely long cylinder

$$N = 1$$



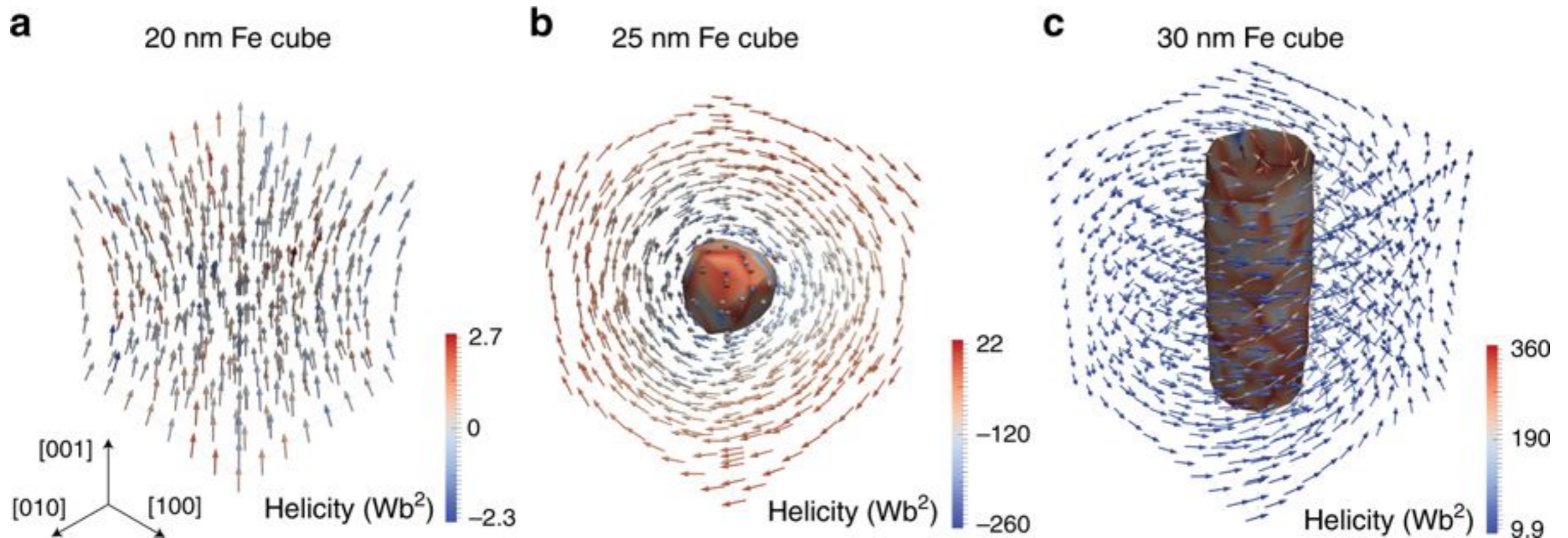
Short cylinder

Infinitely long cylinder



# Beware of non-uniformities

- In general magnetization is not uniform for other shapes



# Magnetic units

# Magnetism units

- The older Gaussian/cgs units are still common in the literature
- (Some) conversion factors between the different systems

Quantity	Symbol	Gaussian & cgs emu	Conversion factor	SI
Magnetic flux density	B	gauss (G)	$10^{-4}$	tesla (T)
Magnetic field strength	H	oersted (Oe)	$10^3/4\pi$	A/m
Magnetization	M	emu/cc	$10^3$	A/m, J/T/m <sup>3</sup>
Magnetic Moment	m	emu	$10^{-3}$	Am <sup>2</sup> , J/T
Permeability of free space	$\mu_0$	dimensionless	$4\pi \times 10^{-7}$	H/m, T <sup>2</sup> J <sup>-1</sup> m <sup>3</sup>

[http://www.ieeemagnetics.org/images/stories/magnetic\\_units.pdf](http://www.ieeemagnetics.org/images/stories/magnetic_units.pdf)

# Old units

- Redefinition of SI system in 2018 now makes the speed of light  $c$  and electronic charge  $e$  fixed constants.
- Now  $\mu_0$  is in principle a measurable quantity, defined from the fine structure constant  $\sim 1/137$

$$\begin{aligned}(h/e^2)_{\text{exp}} &= (\mu_0 c/2)_{\text{fixed}} \cdot (1/\alpha)_{\text{exp}} \\ (\mu_0)_{\text{exp}} &= (2h/ce^2)_{\text{fixed}} \cdot (\alpha)_{\text{exp}}\end{aligned}$$

- This breaks the previous convention fixing  $\mu_0$  as  $4\pi \cdot 10^{-7}$  H/m and thus compatibility between the SI units and old CGS units

Magnetics has been one of the scientific disciplines most resistant to adoption of the SI. With the revised SI, the “peaceful coexistence” of two systems of units [Silsbee 1962] is no longer feasible. The following recommendations warrant consideration.

- 1) Scholarly journals that publish articles in magnetics should require use of the SI and disallow EMU such as oersted, gauss, and “emu per cubic centimeter.” Authors who find the expression of magnetic field strength  $H$  in units of ampere per meter to be inconvenient could instead refer to  $\mu_0 H$  in units of tesla (or milli-, micro-, nano-, or picotesla). Similarly, magnetization  $M$  could be expressed as  $\mu_0 M$  or as magnetic polarization  $J$  in units of tesla or millitesla.
- 2) For the benefit of future generations of magneticians, professors should use SI in classroom instruction. Commercial instruments and magnetometers should be programmed to report measurement results in SI.
- 3) In writing equations, it is adequate to use phrases such as “where  $\mu_0$  is the permeability of vacuum” (or “the vacuum magnetic permeability” or “the permeability of free space” or “the magnetic constant”) without giving a numerical value. This follows typical usage when referring to the speed of light  $c$ , the Boltzmann constant  $k$ , or the Bohr magneton  $\mu_B$ .

# A recent trend to using teslas for everything

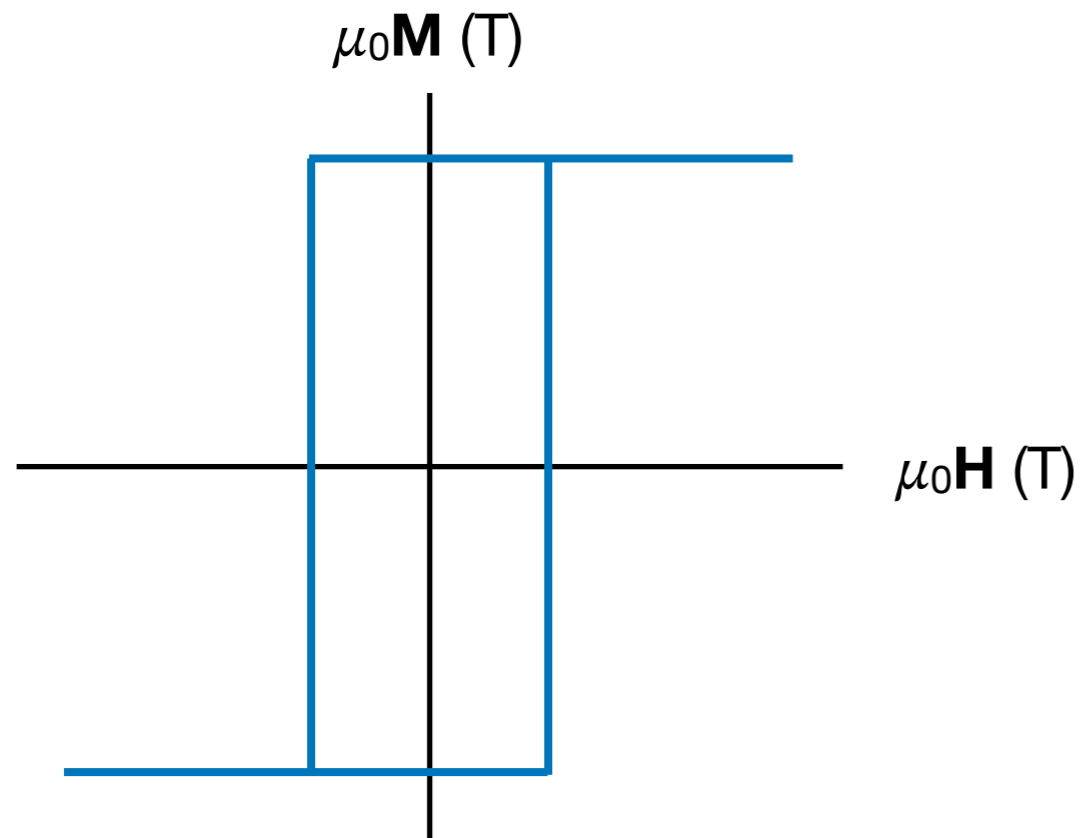
- $\mathbf{B}$ ,  $\mu_0\mathbf{M}$  and  $\mu_0\mathbf{H}$  are all defined in terms of magnetic field (intensity) in teslas (T)
- Started with Superconducting and Permanent magnet communities, probably due to avoidance of odd numerical conversions, dimensions and units
- Now common in the literature, theoretical and experimental
- Best way is to think about everything as current loop 'sources' of flux  $\mu_0\mathbf{M}$  and  $\mu_0\mathbf{H}$

$$\mathbf{B} = \mu_0 (\mathbf{M} + \mathbf{H})$$

$$\mathbf{B}_{\text{total}} = \mathbf{B}_{\text{magnetization}} + \mathbf{B}_{\text{applied}}$$

- This convention leads to oddities in hysteresis - what are the units of  $\mathbf{M} \cdot \mathbf{B}$ , both in Tesla??

# Making sense of M-B loops



- Not immediately obvious that this is useful  
- a single loop cycle should give units of energy (density)
- BUT - can easily extract the magnetization in sensible dimensions by dividing by  $\mu_0$

$$\mathbf{M}(\text{JT}^{-1}\text{m}^{-3}) \equiv \frac{\mu_0 \mathbf{M}}{\mu_0} \frac{(\text{T})}{(\text{T}^2 \text{J}^{-1} \text{m}^3)}$$

- In this case, a hysteresis cycle  $\text{Int}(\mathbf{M} \cdot \mathbf{B})$  has units of  $\text{J}/\text{m}^3$
- Same is true of  $\mathbf{B}_{\text{tot}}(\mathbf{H})$  loops but with inverted units

# Summary

- Magnetic moments and current loops behave equivalently
- Quantum mechanical origin of magnetic moments not too far from a classical current loop
- Magnetic fields are different inside and outside magnetic media
- Internal magnetic fields in magnets are generally complicated
- Units in magnetism are generally horrible, but always use SI
- Remembering that  $\mu_0$  has units of  $\text{T}^2 \text{J}^{-1} \text{m}^3$  will make you happy