

Electric field control of magnetism

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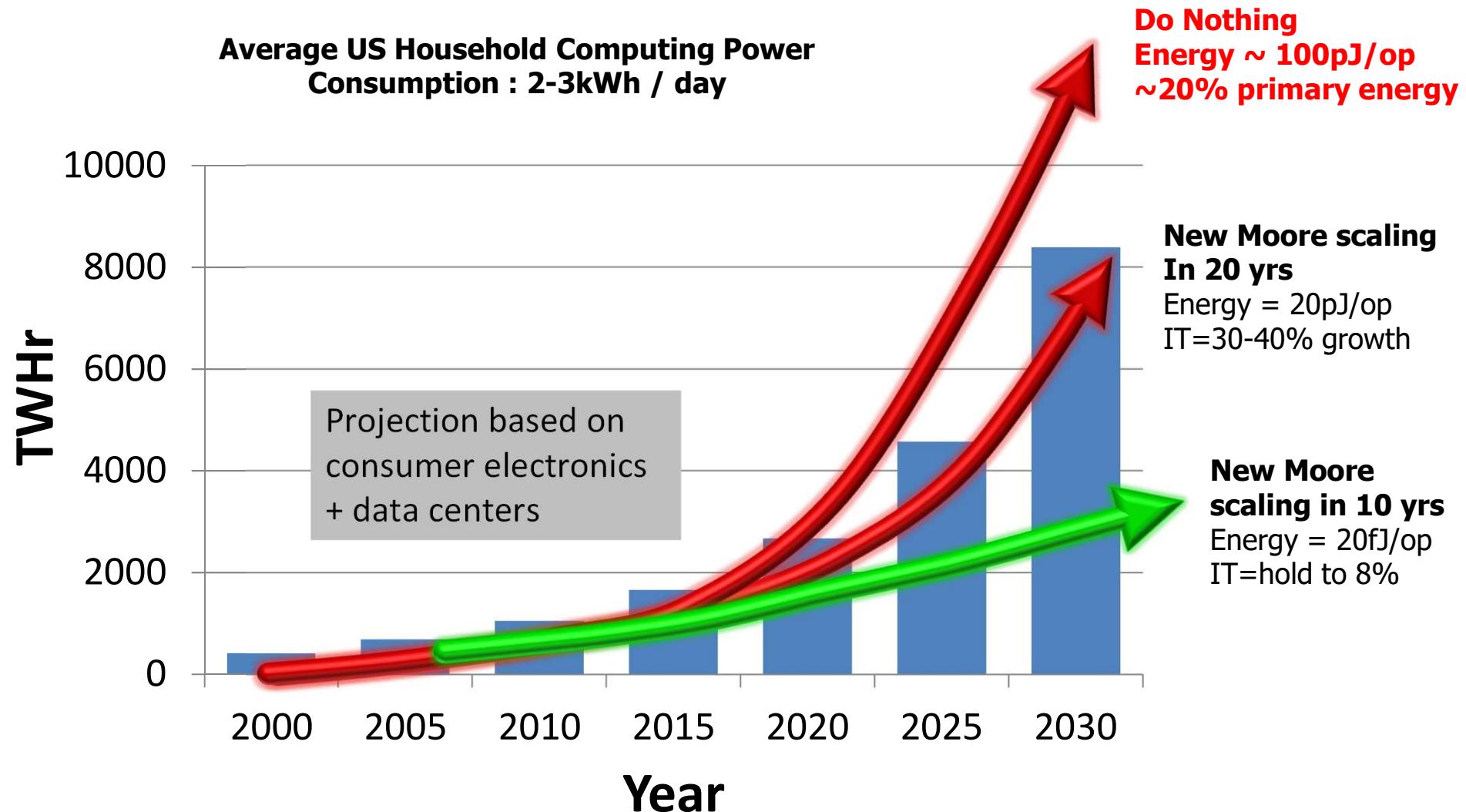
“Electrical control of magnetism by electric field and current-induced torques”

A. Fert, F. Casanova, V. Garcia, R. Ramesh and M. Bibes

Rev. Mod. Phys. 95, 015005 (2024) ; ArXiv 2311.11724

Power consumption of ICT systems

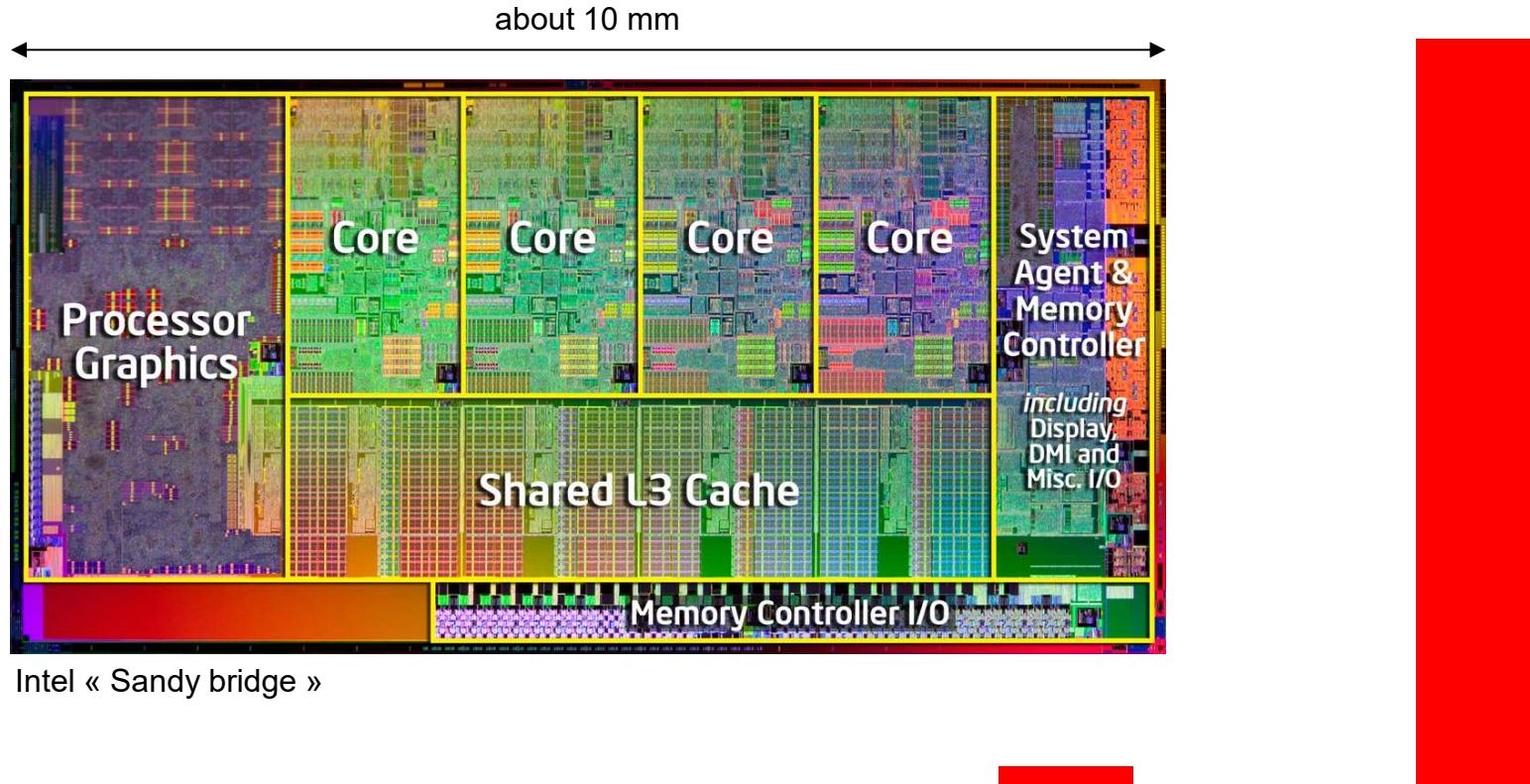
Global Semiconductor market size ~ \$2 trillion by 2030



www.alliancetrustinvestments.com/sri-hub/posts/Energy-efficient-data-centres

Courtesy R. Ramesh

Power needs in microprocessors



Intel « Sandy bridge »

Read one 64-bit
number in SRAM

14 pJ

Multiply two 64-bit
numbers

50 pJ

Move one 64-bit
number 10 mm away

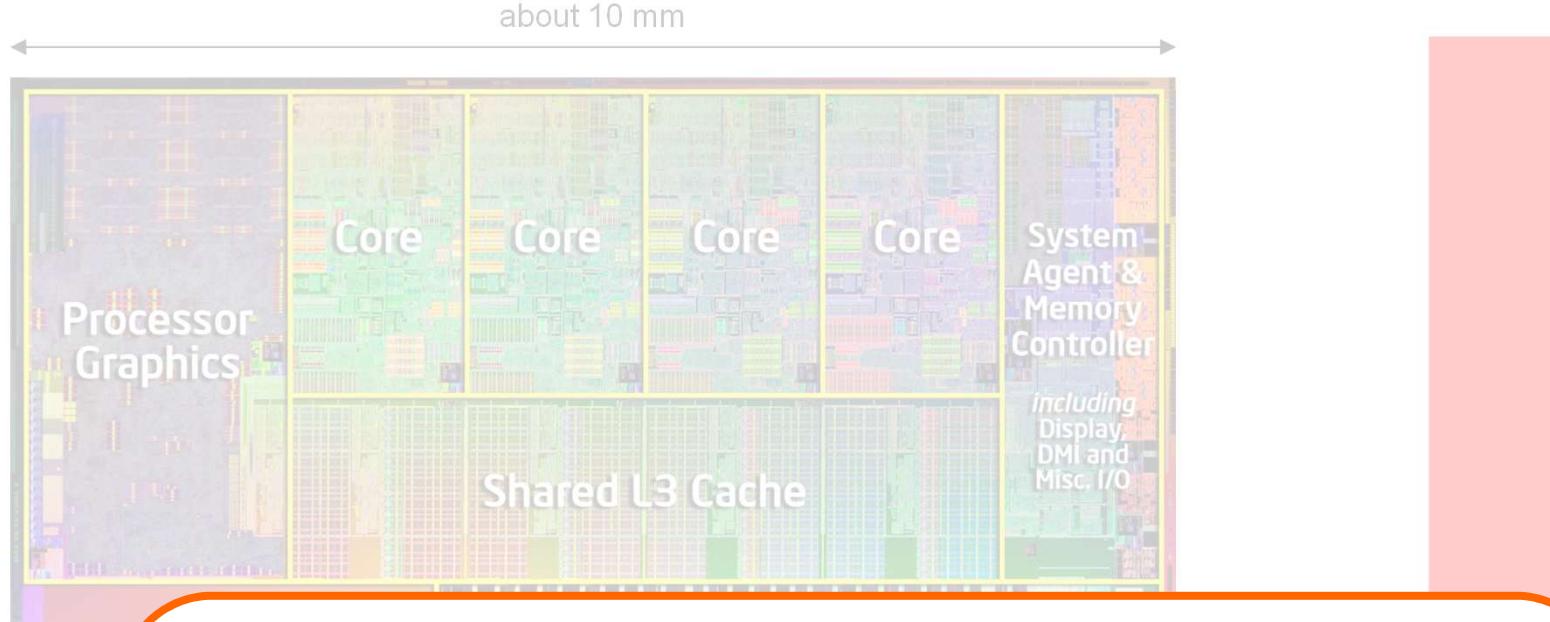
300 pJ

Move one 64-bit number
from external RAM

10000 pJ

GPUs and the future of parallel computing, W.J. Dally et al, IEEE Micro (2011)

Power needs in microprocessors



Data transfer on-chip consumes most of the power !

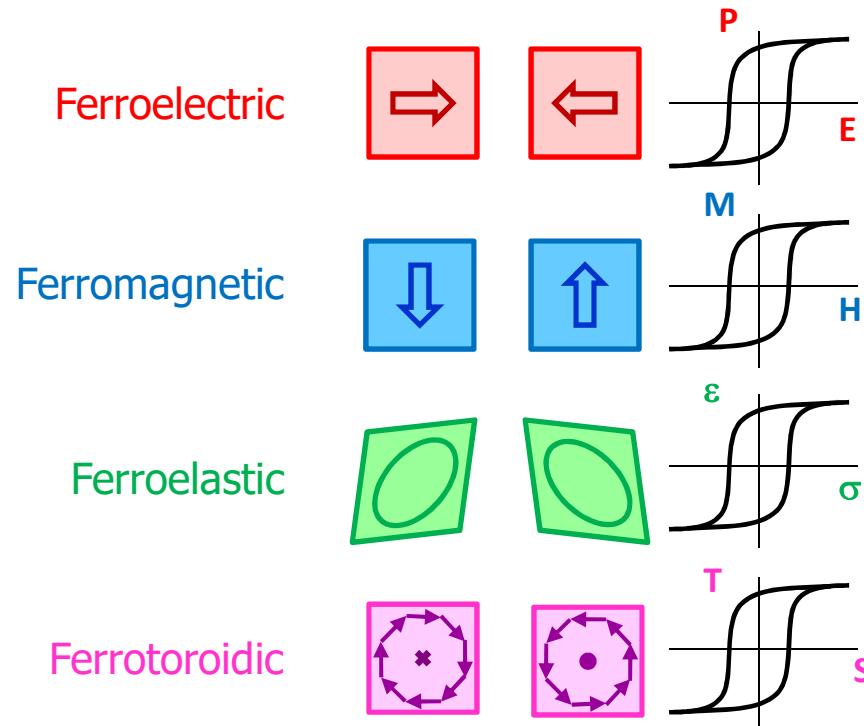
- « Bring memory into logic »
- Embed memory elements in the logic units
- Ferroic materials can bring solutions for beyond CMOS electronics

GPUs and the future of parallel computing, W.J. Dally et al, IEEE Micro (2011)

- ○ Ferroic orders, multiferroics and magnetoelectric coupling
- Approaches for the electric-field control of magnetism
- Electric-field control of spin-charge interconversion
- Low-power spin-based devices for logic-in-memory

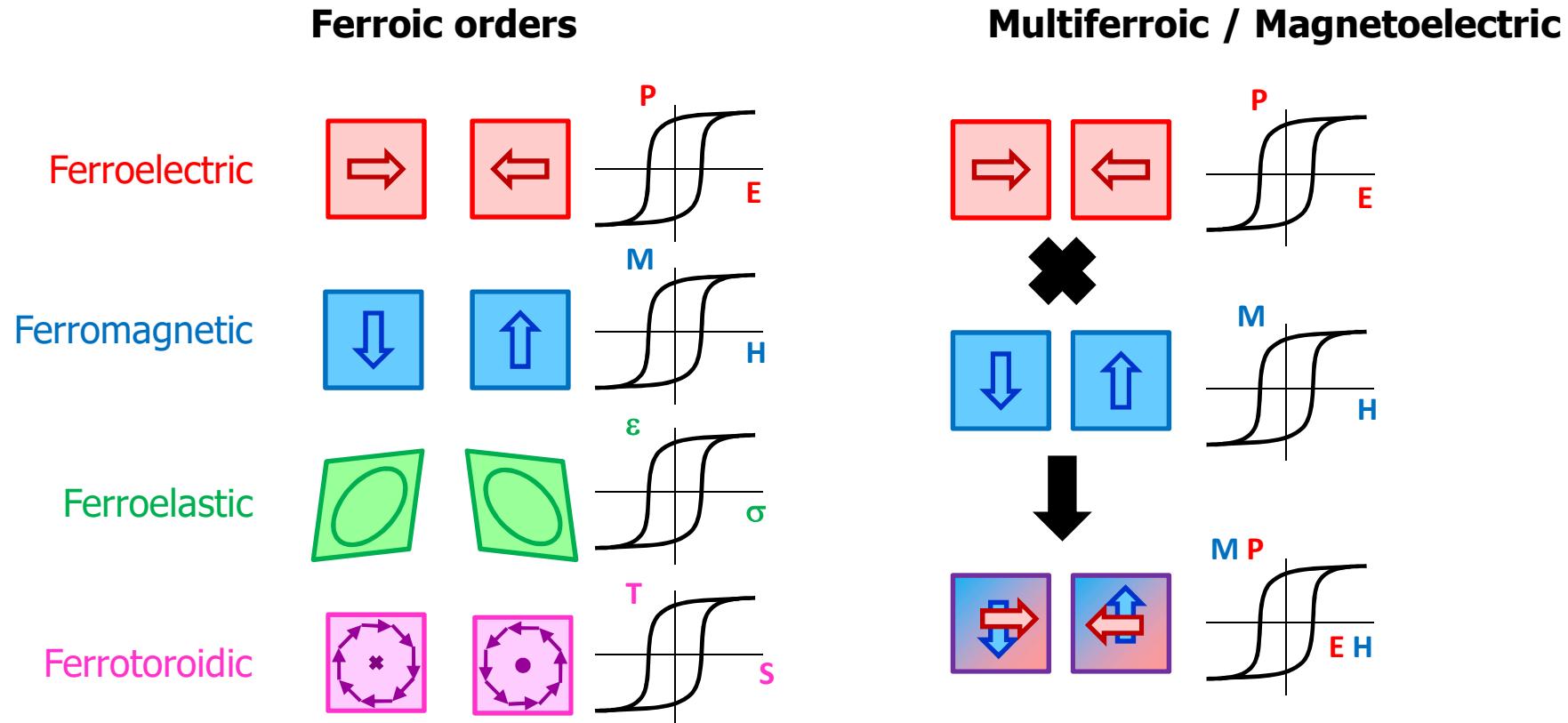
Introduction to ferroic orders

Ferroic orders



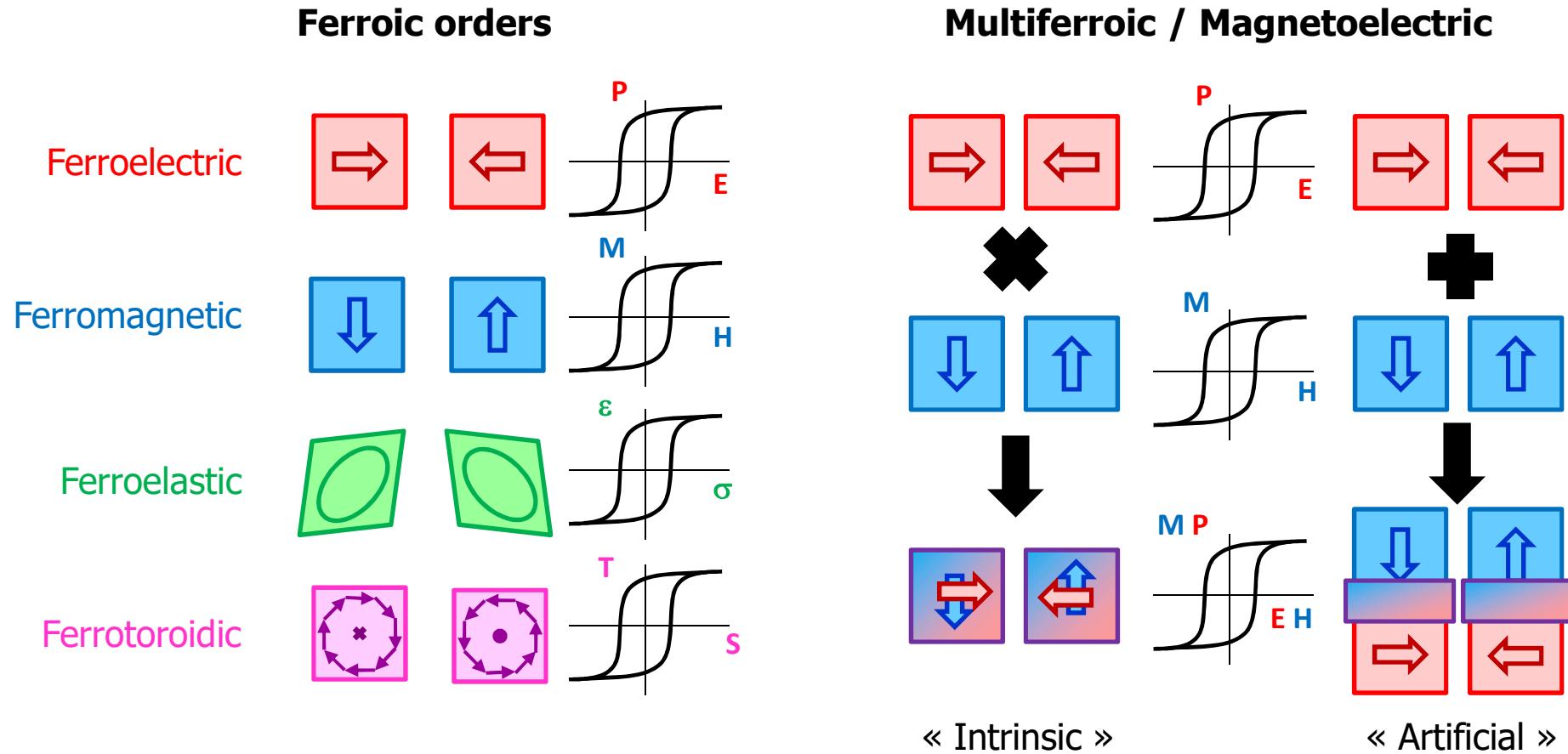
- Hysteretic dependence of order parameter : good for data storage

Introduction to ferroic orders



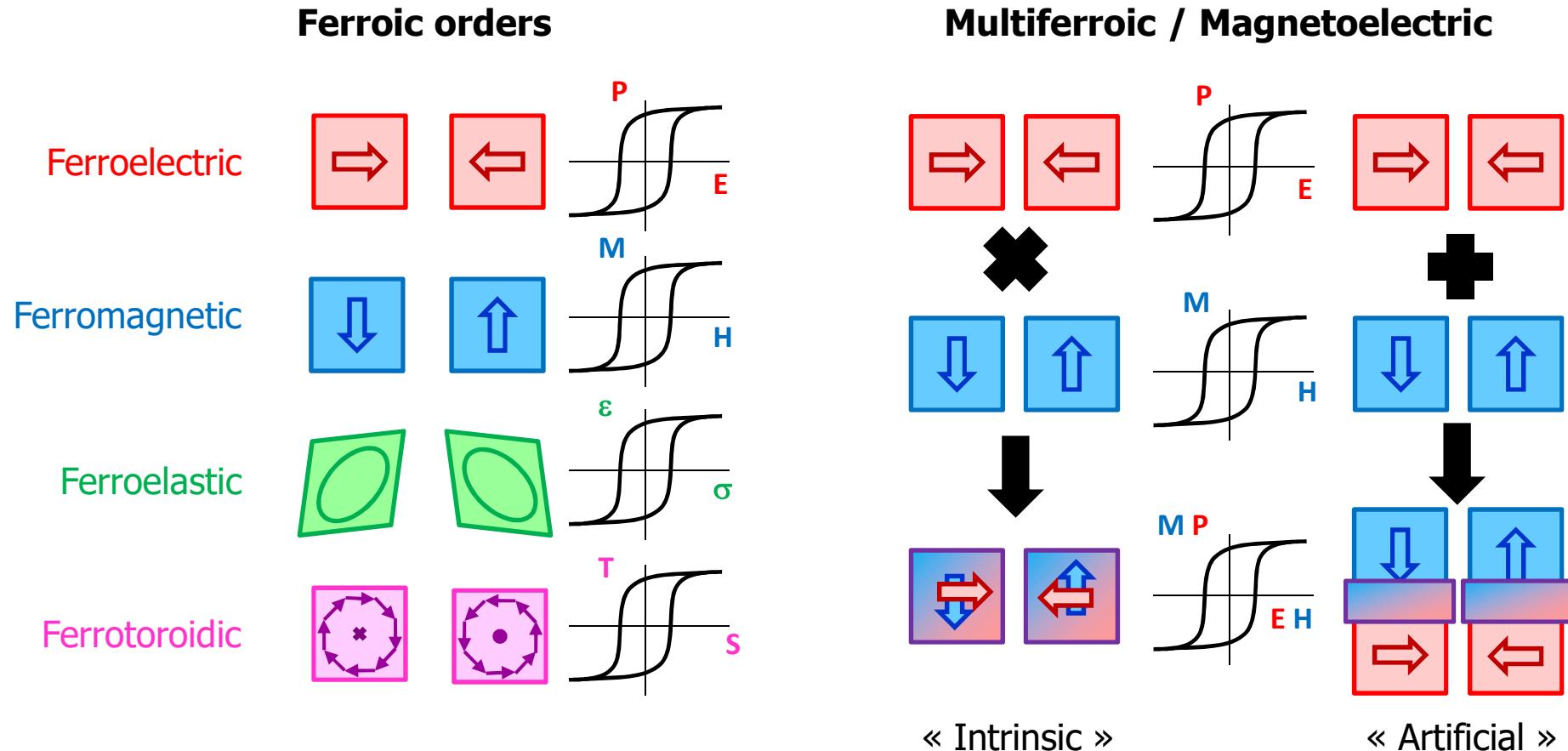
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Introduction to ferroic orders



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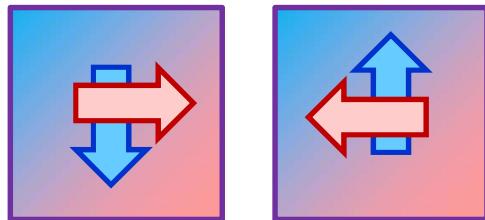
Introduction to ferroic orders



- Hysteretic dependence of order parameter : good for data storage
- Multiple order parameters : increased storage density
- Coupled orders : enhanced flexibility for data writing

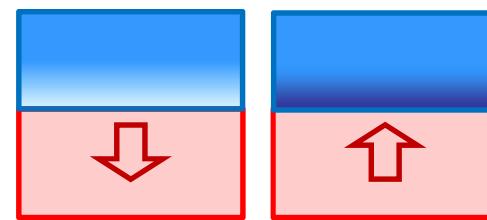
MB, Nature Mater. 11, 354 (2012)

Intrinsic magnetoelectric



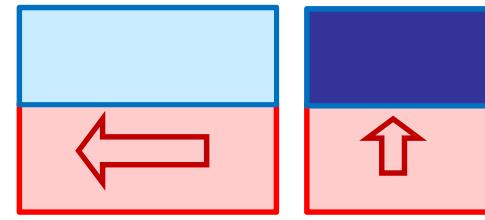
Use single-phase multiferroic material

Field-effect



Combine strong ferroelectric with carrier-mediated ferromagnet

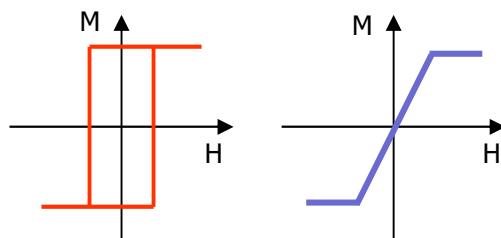
Strain-driven



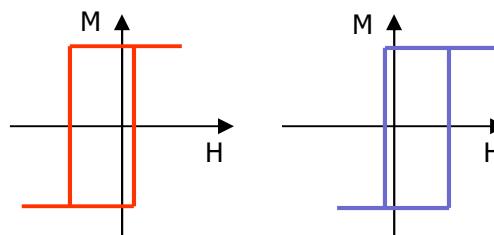
Combine piezoelectric or ferroelectric/ferroelastic with magnetostrictive ferromagnet

Controlling magnetism with electric fields

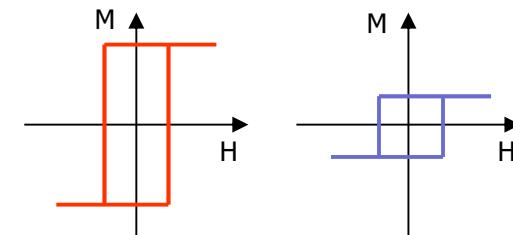
Magnetic anisotropy



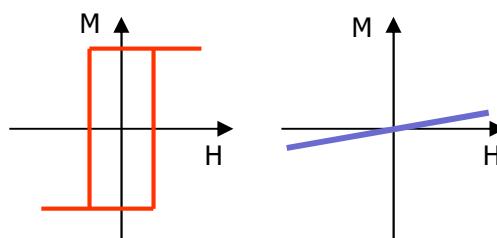
Exchange bias



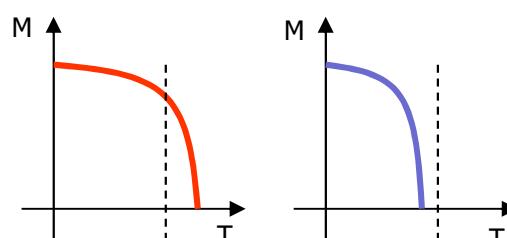
Magnetic moment



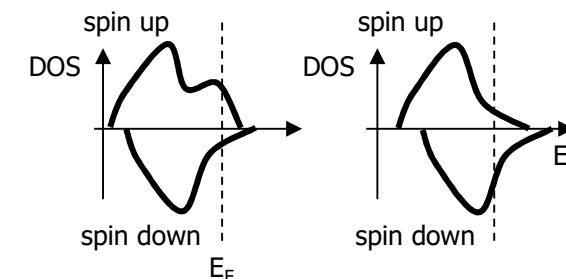
Magnetic order



Curie temperature



Spin polarization



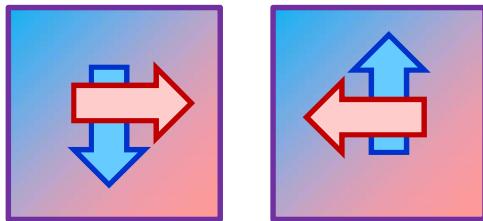
$P_{\text{spin}} > 0$

$P_{\text{spin}} < 0$

MB, Nature Mater 11, 354 (2012) & MB et al, Annu. Rev. Mater. Res. 44, 91 (2014)

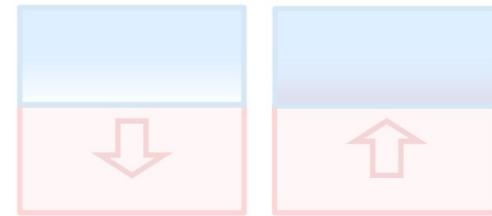
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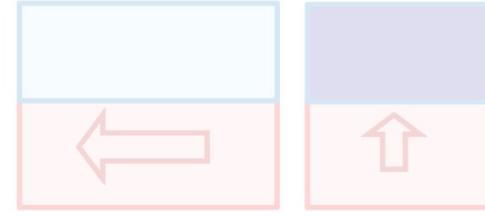
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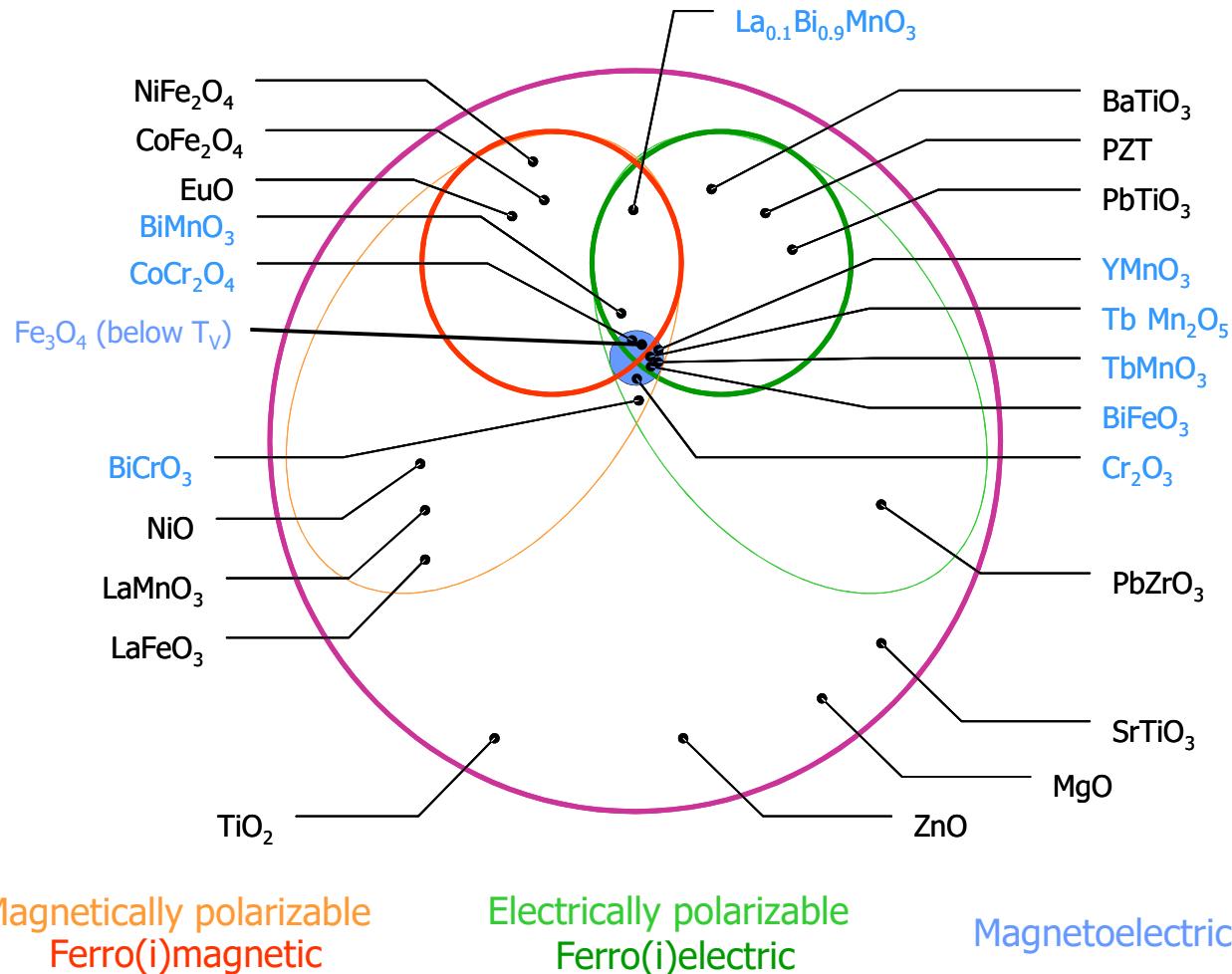
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Combine piezoelectric or ferroelectric/ferroelastic with magnetostrictive ferromagnet

E-field control of magnetism with intrinsic multiferroics

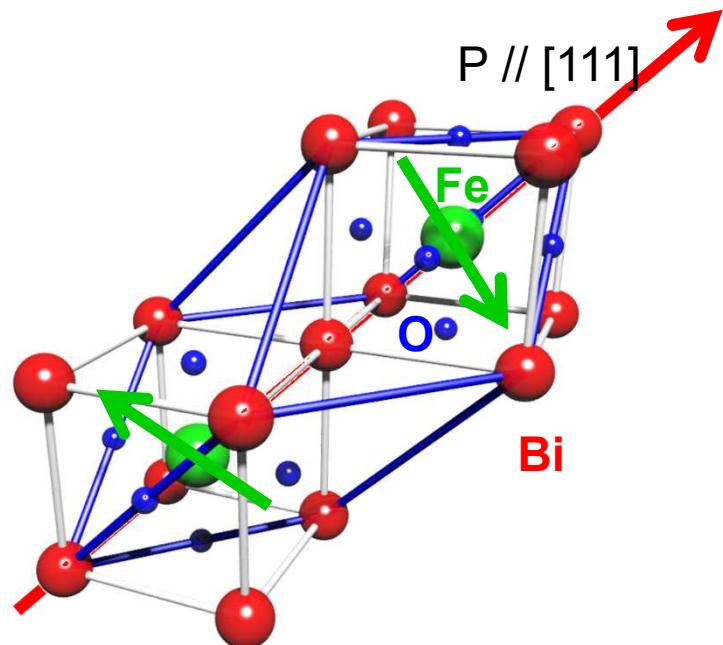
There are very few (room-temperature) multiferroics



H. Béa, MB et al, J. Phys.: Condens. Matter 20, 434221 (2008)

Derived from Eerenstein, Mathur and Scott, Nature 442, 759 (2006)

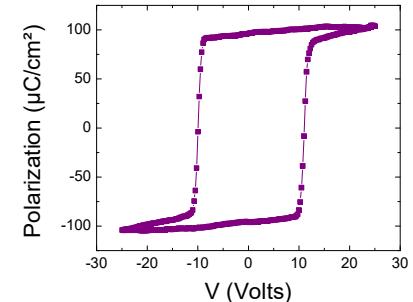
BiFeO_3 : a room-temperature multiferroic



Ferroelectric properties

- Very high $T_C \approx 1100$ K
- Very large $P=100 \mu\text{C}/\text{cm}^2$

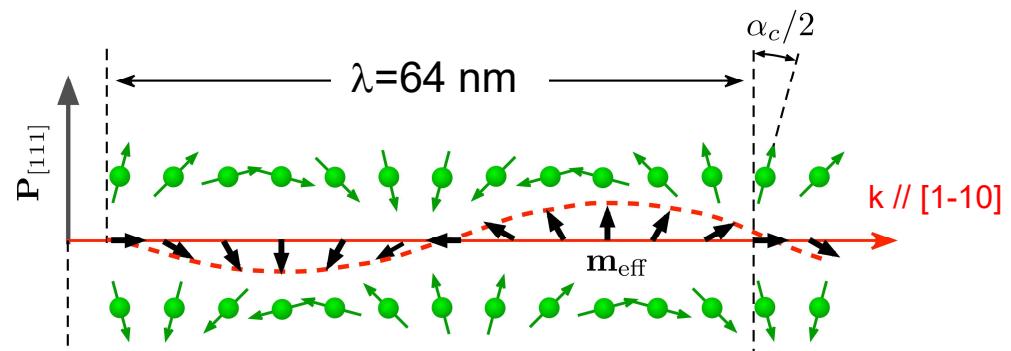
Fisher et al., J. Phys. C, 13, 1931 (1980)
Béa, MB et al, APL 93, 072091 (2008)



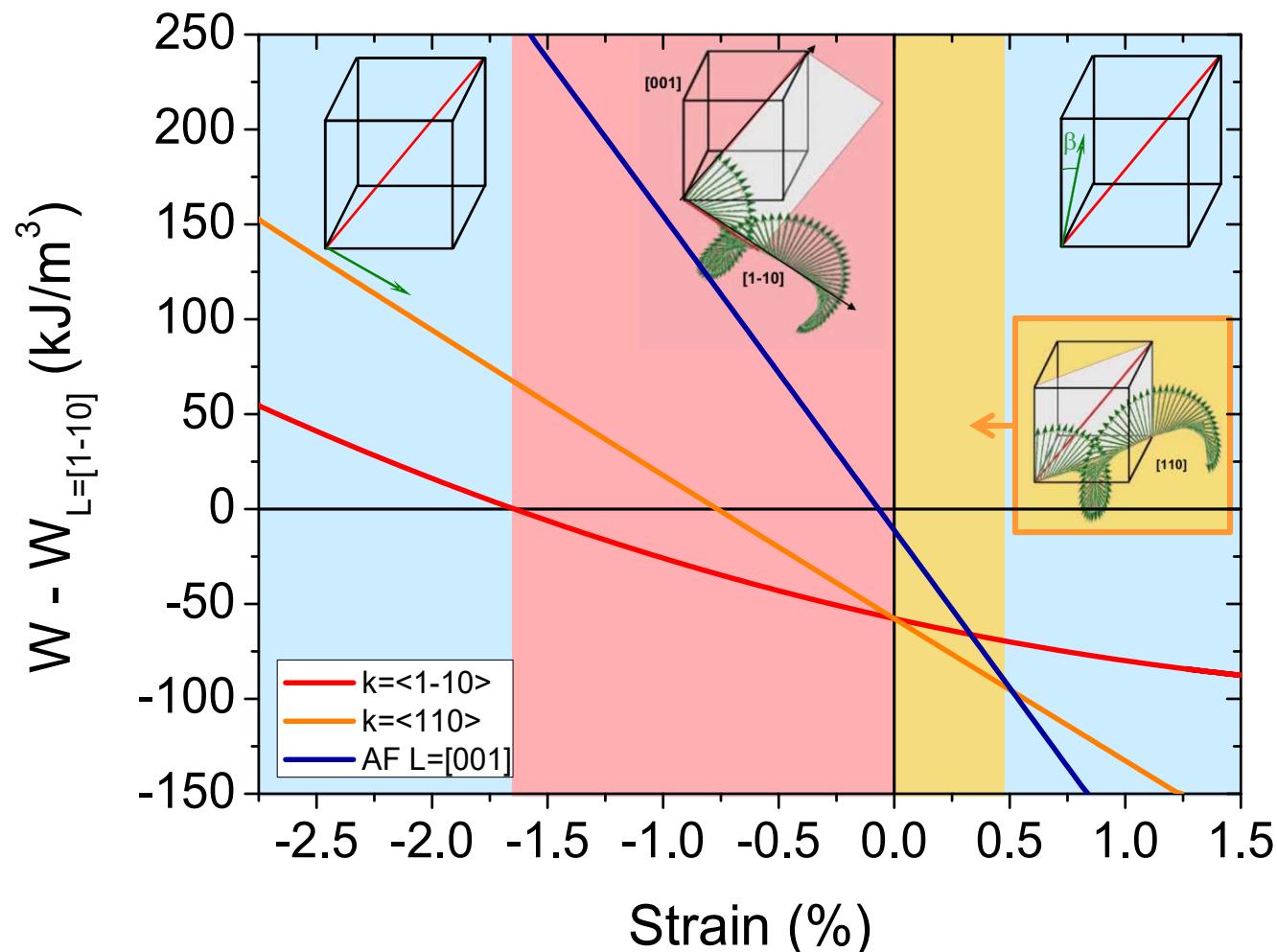
Magnetic properties

- G-type antiferromagnetic + cycloidal modulation ($\lambda=62$ nm)
- Weak moment with periodic modulation
- $T_N \approx 640$ K

Sosnowska et al., J. Phys. C, 15, 4835 (1982)



Influence of epitaxial strain on the magnetic properties of BiFeO₃

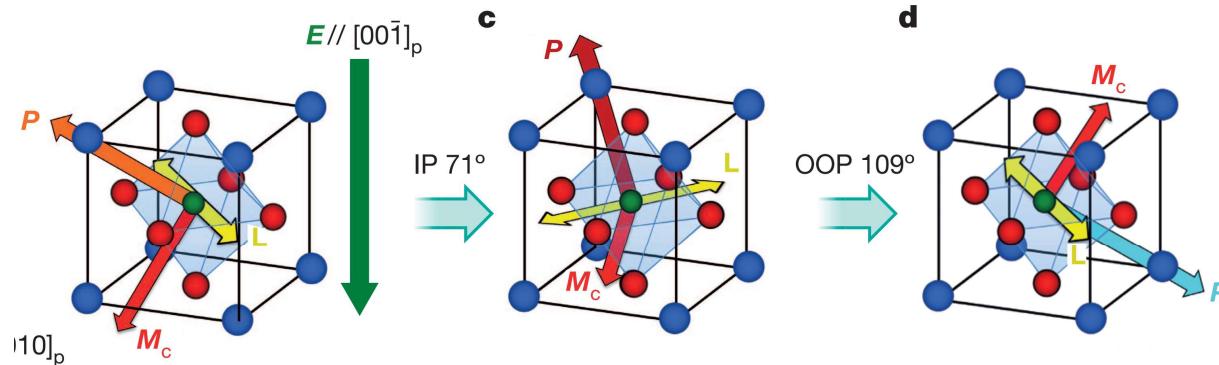


- Cycloidal state is destabilized by strain-induced (magnetoelastic) anisotropy
- Weak-FM state at high tensile or compressive state } Mössbauer spectroscopy + theory
- New cycloid stabilized at low tensile strain

D. Sando, MB et al, Nature Mater. 12, 641 (2013)

E-field induced magnetization switching with BiFeO_3 thin films

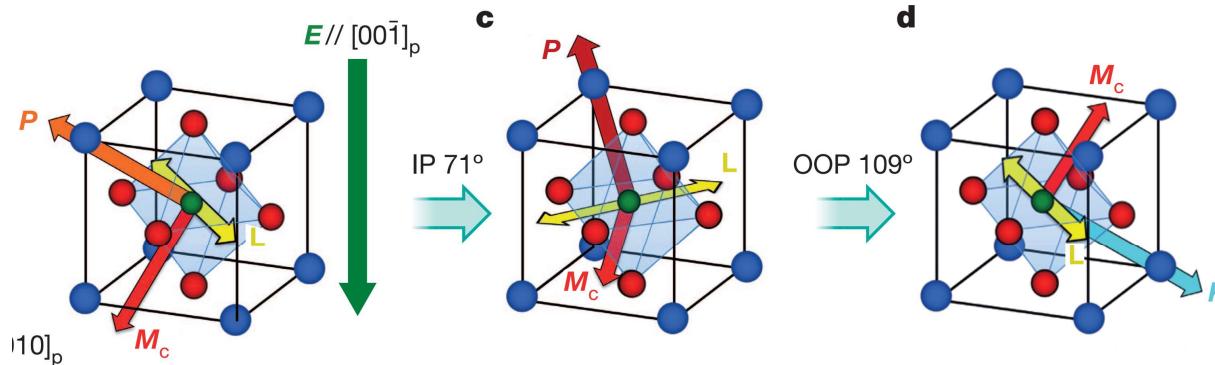
Sequential switching of P promotes switching of weak M



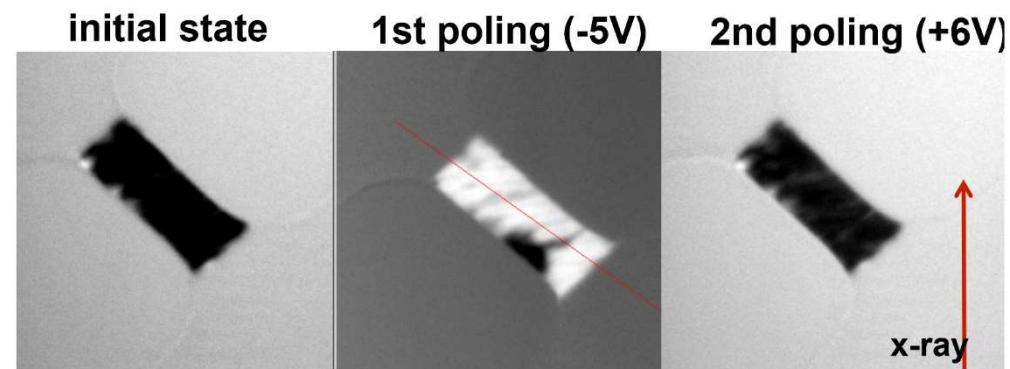
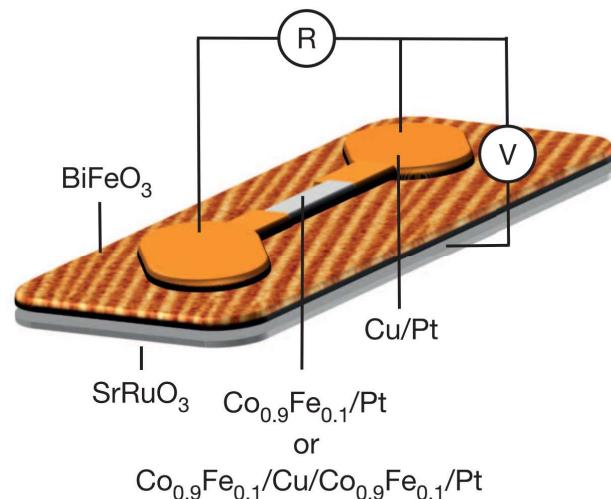
Heron et al, Nature 370, 516 (2014)

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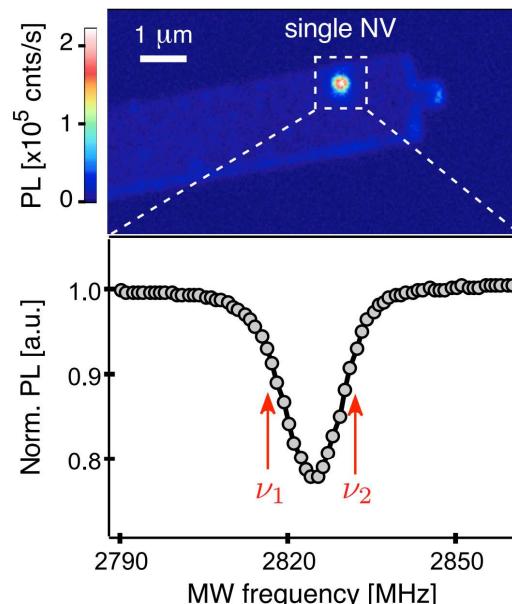
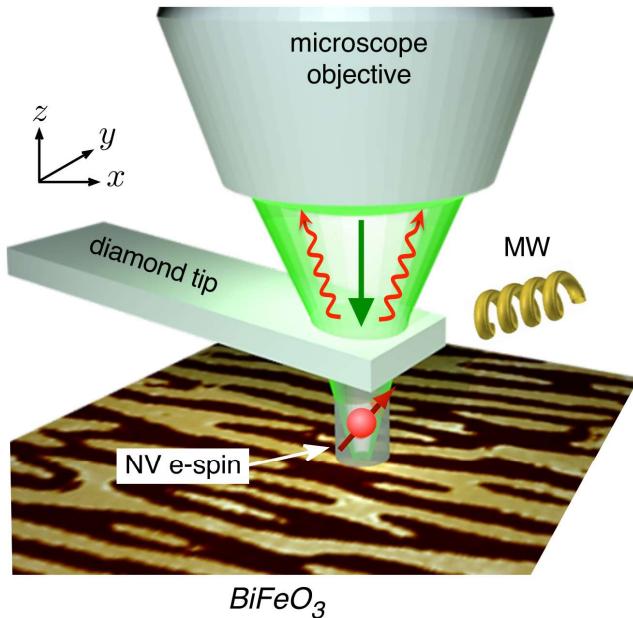


Heron et al, Nature 370, 516 (2014)



- Application of out-of-plane voltage to BFO film promotes local switching of magnetization in Co film grown on top of BFO

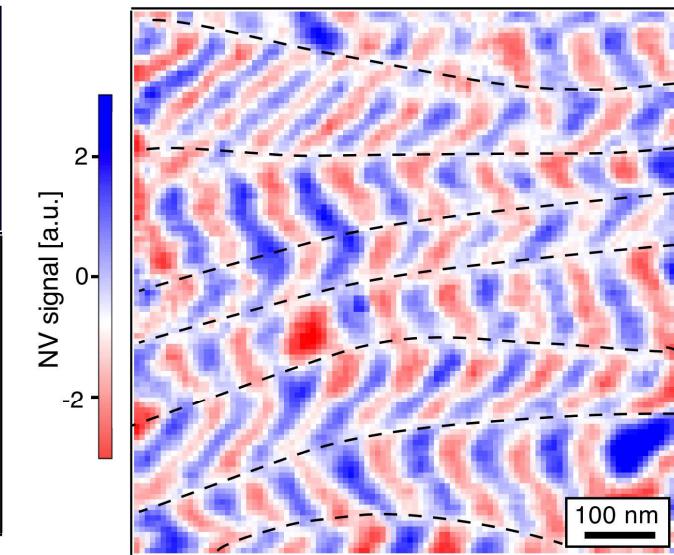
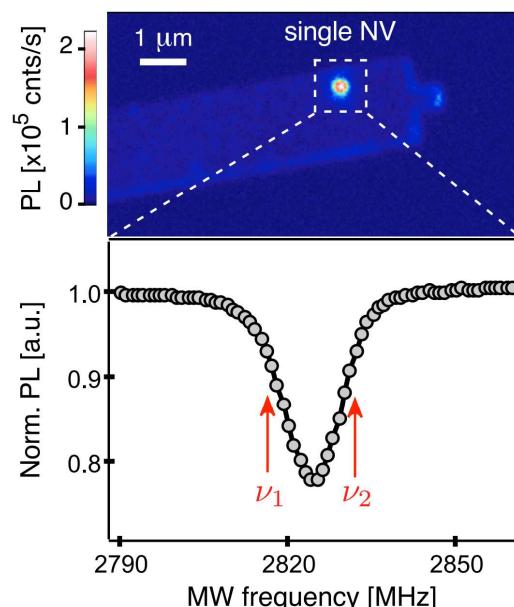
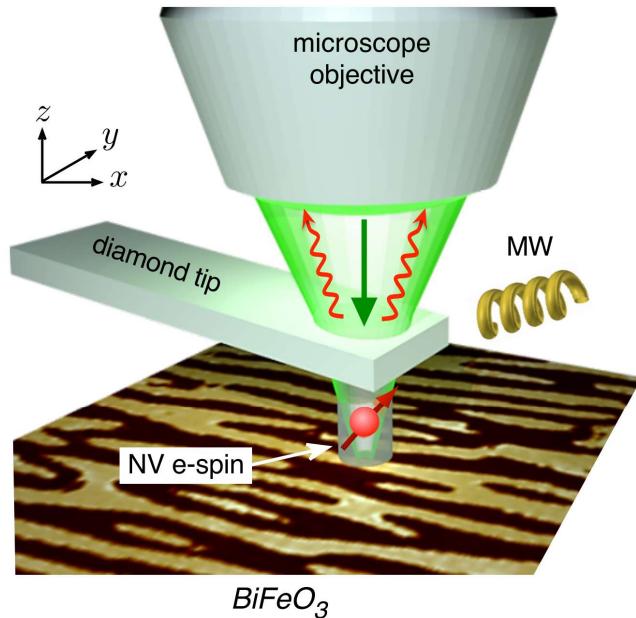
Direct magnetic imaging of BFO with NV center microscopy



- Atomic-sized detection volume
- Quantitative and vectorial
- No magnetic back-action onto sample

Collab. V. Jacques, U. Montpellier and
P. Maletinsky, Basel

Direct magnetic imaging of BFO with NV center microscopy



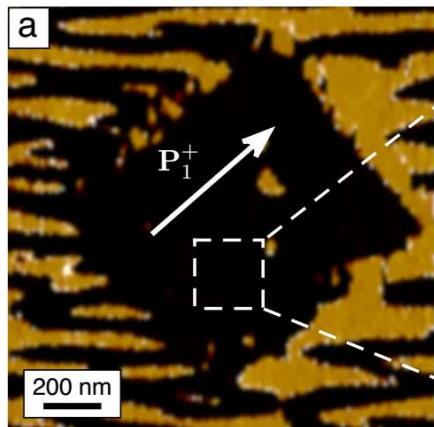
- Atomic-sized detection volume
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- No magnetic back-action onto sample

- Clear periodic contrast of magnetic origin

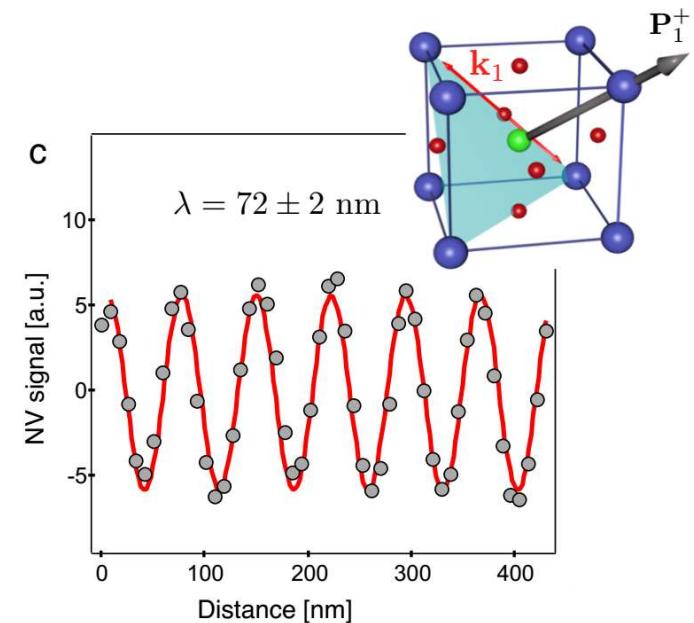
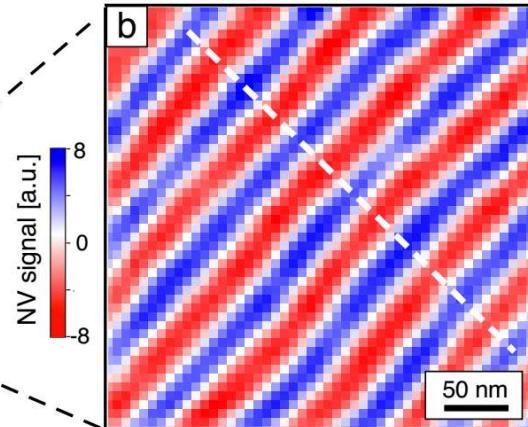
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P. Maletinsky, Basel

Magnetic imaging of single ferroelectric domains

Ferroelectric domains imaged by PFM



Magnetic contrast imaged by NV center microscopy

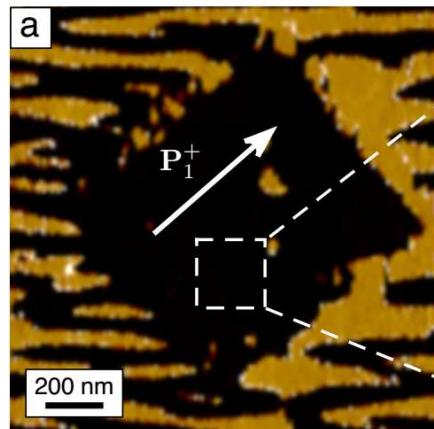


- Electric poling of BFO film : single ferroelectric and magnetic domain
- Periodic contrast with well-defined propagation direction and period
- Period is **72 nm**, slightly longer than in bulk (64 nm), likely due to epitaxial strain

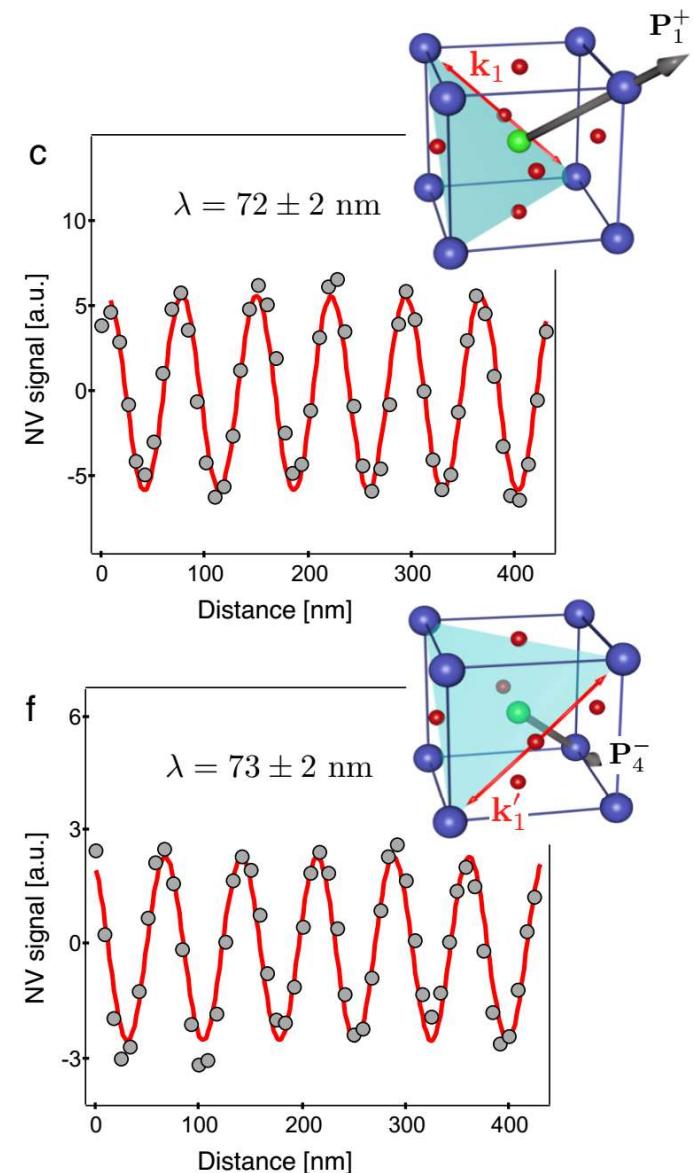
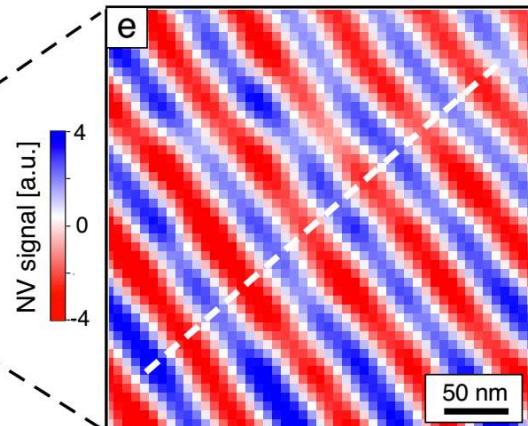
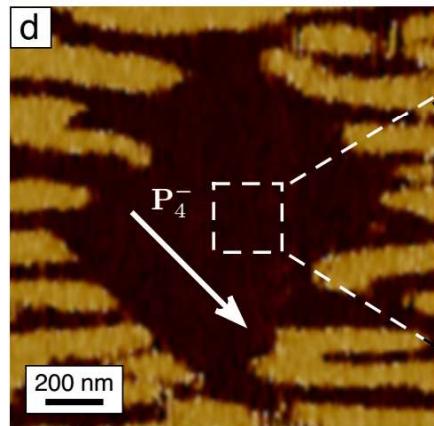
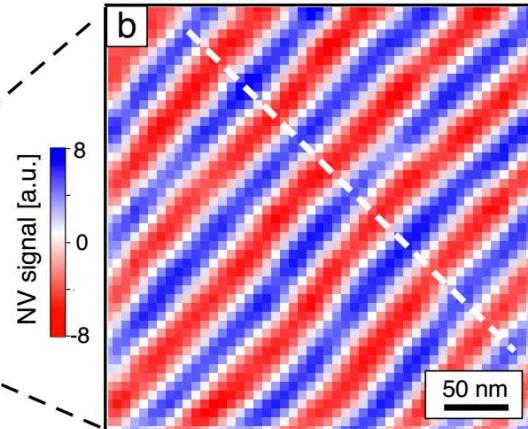
I. Gross, MB et al, Nature 549, 252 (2017)

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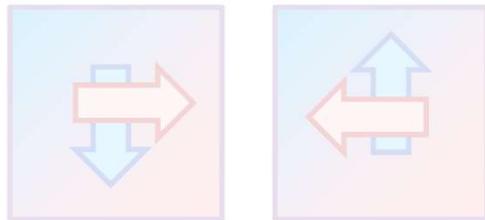
Magnetic contrast imaged by NV center microscopy



Electric-field control of magnetic texture / cycloidal propagation direction

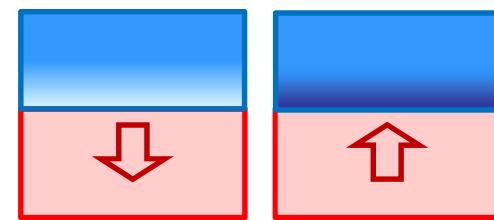
Different approaches for E-field control of magnetism

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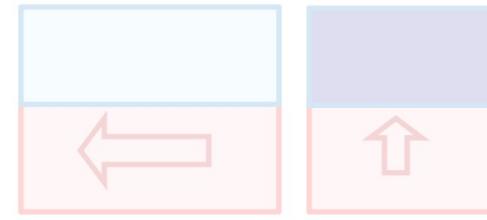
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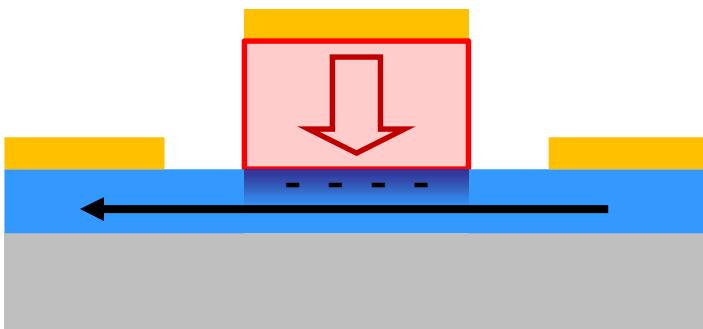
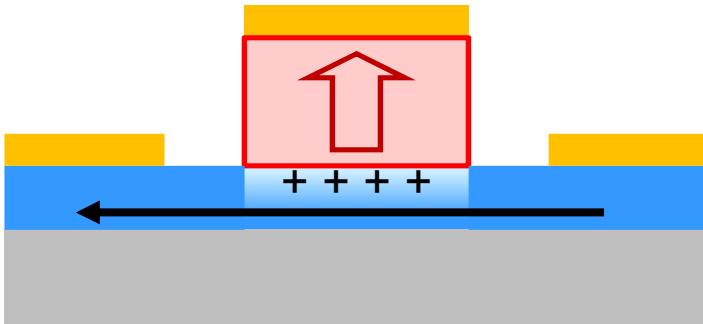
Combine strong ferroelectric with carrier-mediated ferromagnet

Strain-driven



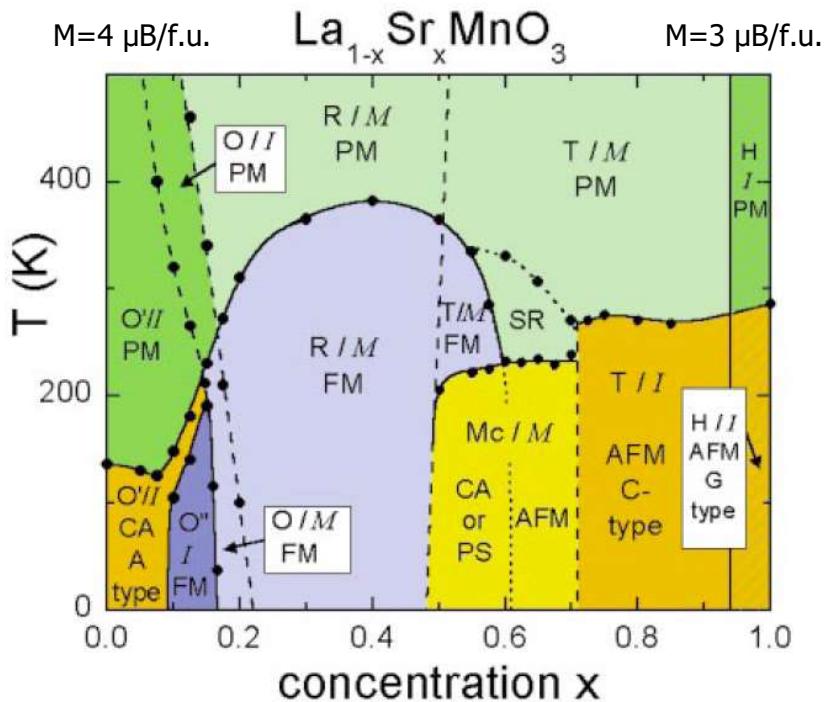
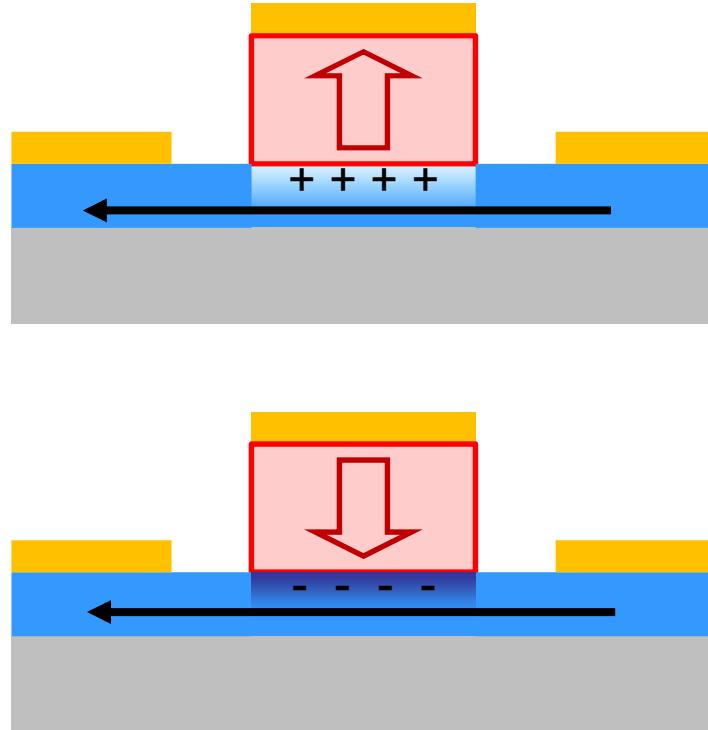
Combine piezoelectric or ferroelectric/ferroelastic with magnetostrictive ferromagnet

Field-effect control of magnetism



- Charge accumulation / depletion thanks to a dielectric or ferroelectric (non-volatile)
- If magnetism in channel material is (highly) sensitive to carrier density
- **Change of magnetic properties by electric field**
- Effect occurs over small distance, typically **Thomas Fermi screening length** (\AA for metals, nm for oxides)

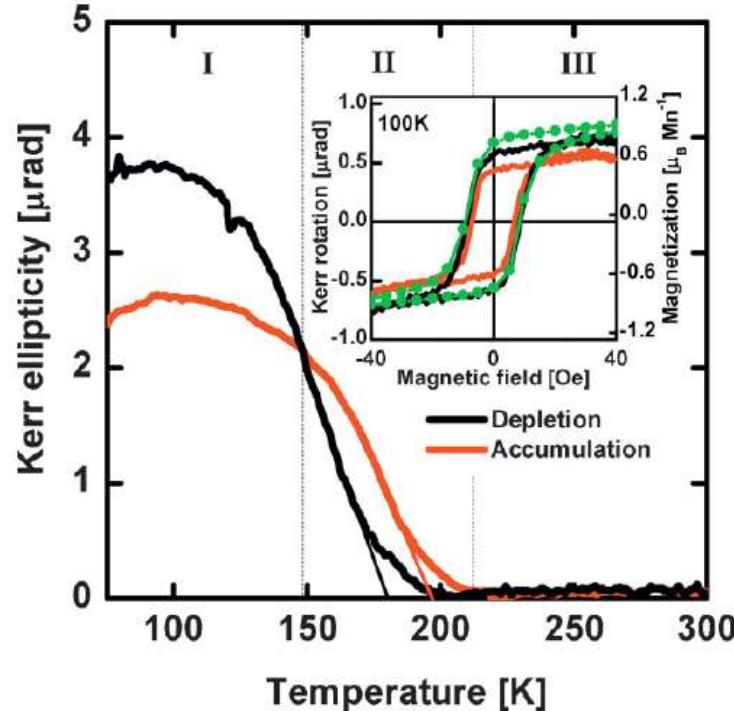
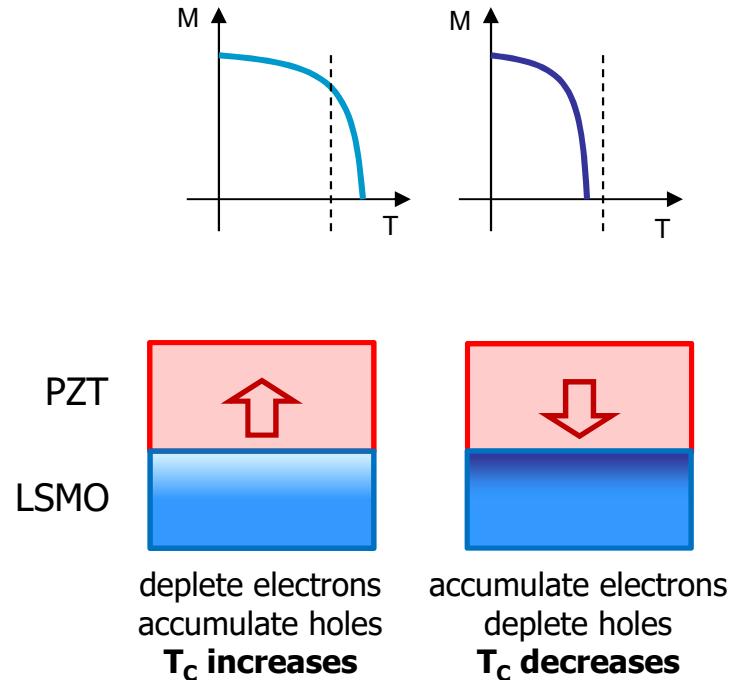
Field-effect control of magnetism



Hemberger et al., PRB, 66, 094410 (2002)

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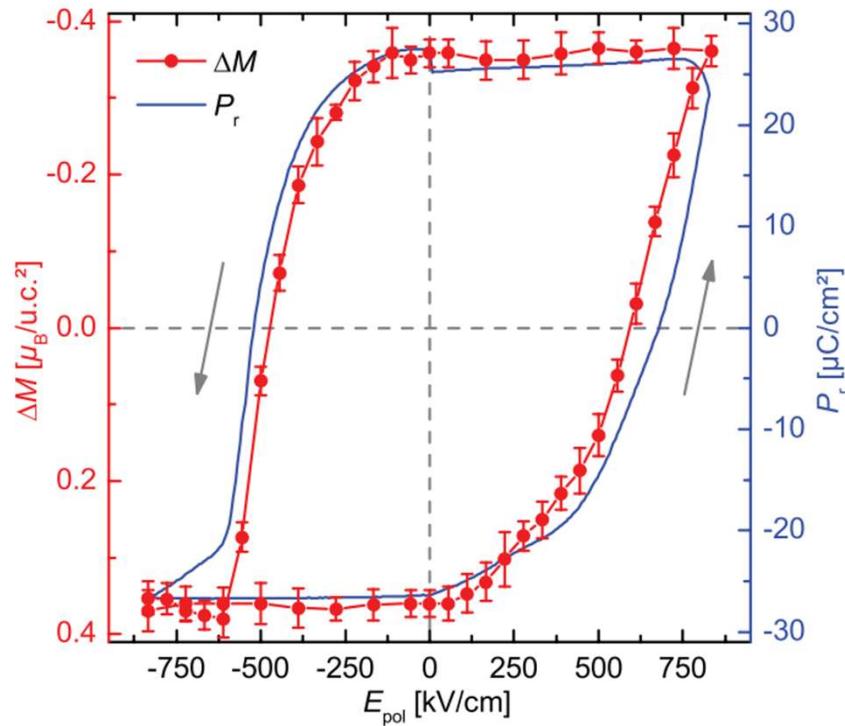
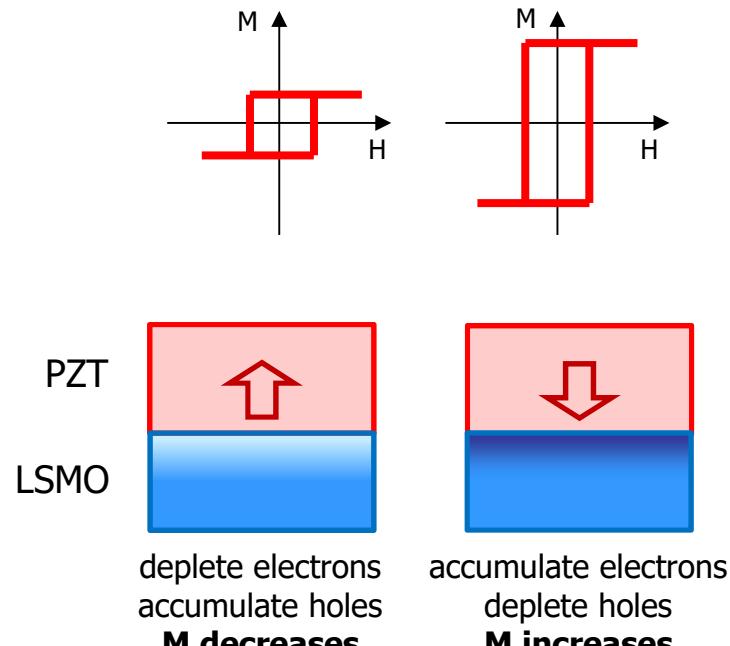
Field-effect control of Curie temperature



Vaz et al, PRL 104, 127202 (2010) & Molegraaf et al, Adv. Mater. 21, 3470 (2009)

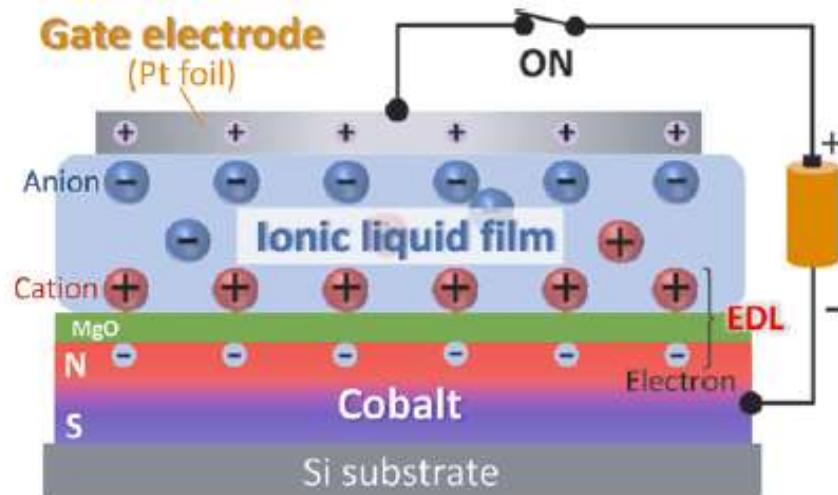
- Combination of a **ferroelectric** and a **carrier-mediated ferromagnet**
- Switching P in ferroelectric PZT produces charge accumulation/depletion in manganite
- Change T_C of manganite
- Limited to low-temperature (also with GaMnAs)

Field-effect control of magnetic moment

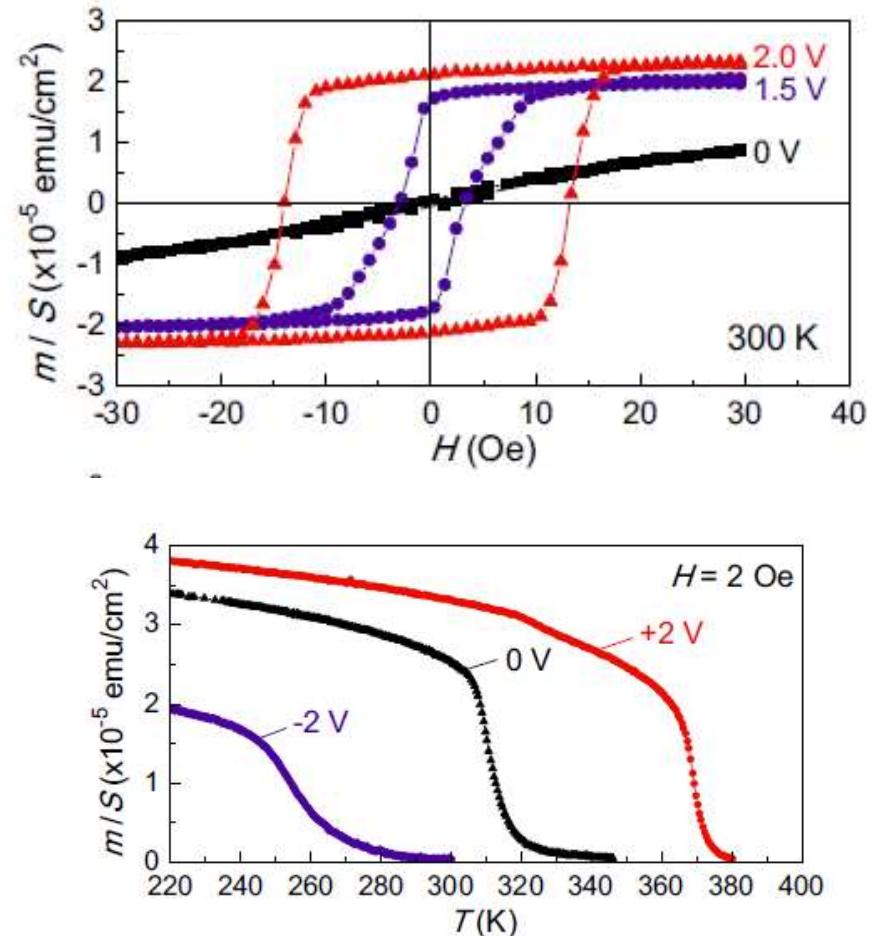


Leufke et al, Phys. Rev. B 87, 094416 (2013)

Increase accumulated charge density : ionic liquids

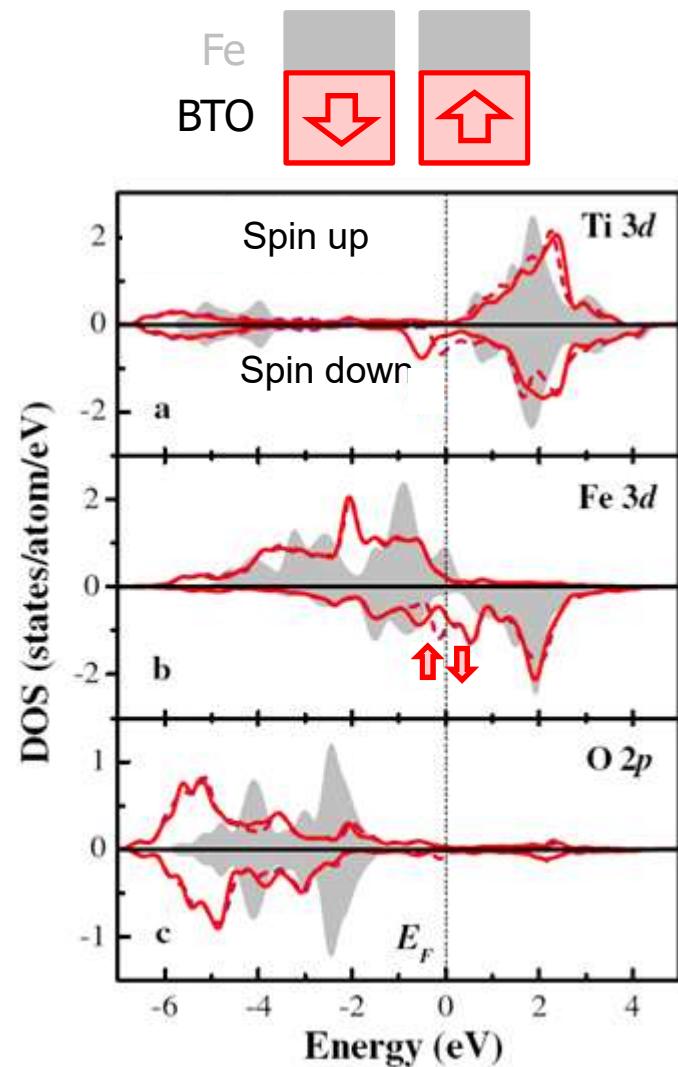


Chiba and Ono, J. Phys. D 46, 213001 (2013)



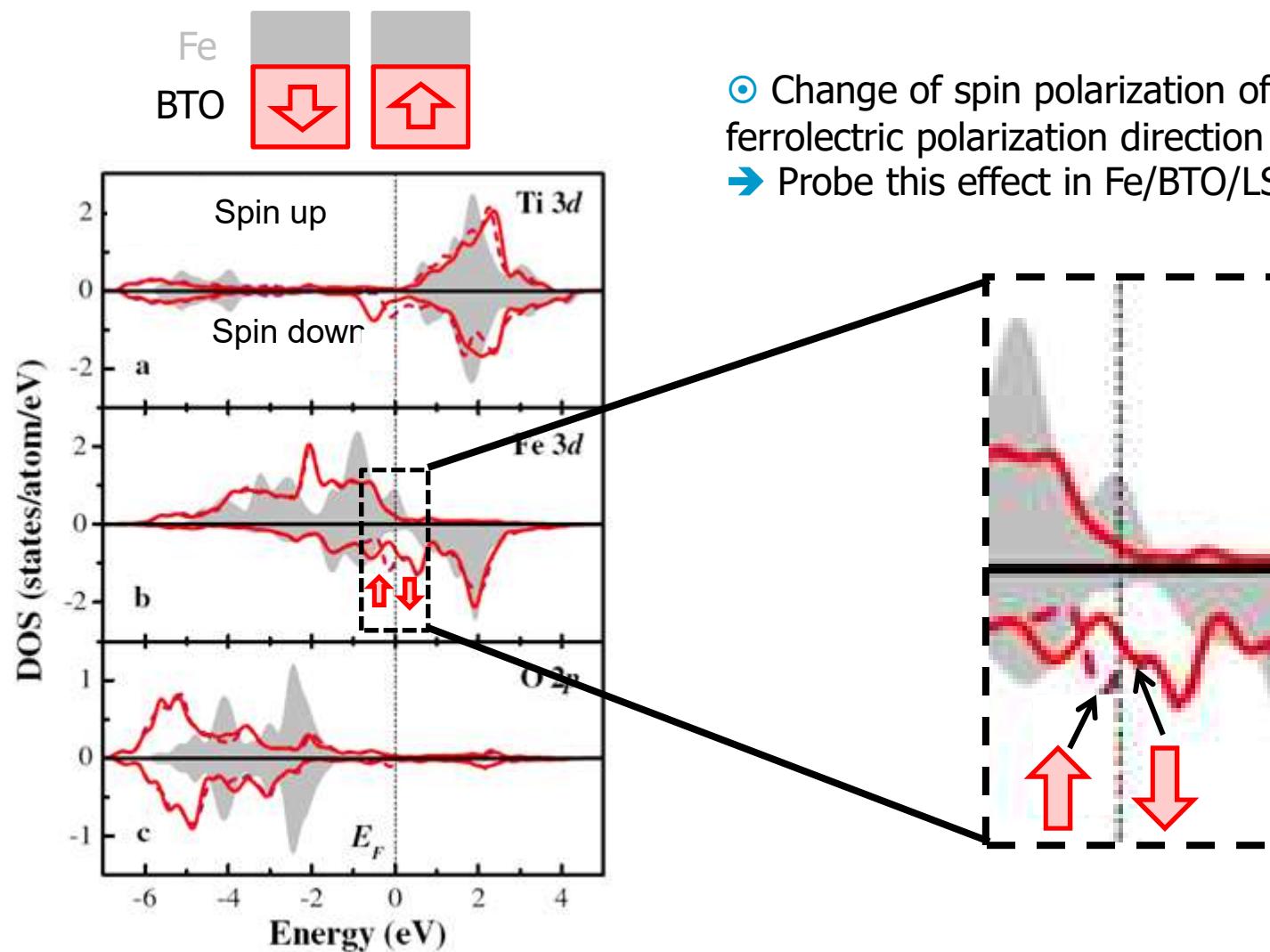
- Large field effect in 0.6 nm Co film using ionic liquid gating
- Possible with ferroelectrics (i.e. PZT/ultrathin Co) ?

Electric-field control of spin polarization



Duan et al., PRL 97, 047201 (2006)
Fechner et al., PRB 78, 212406 (2008)

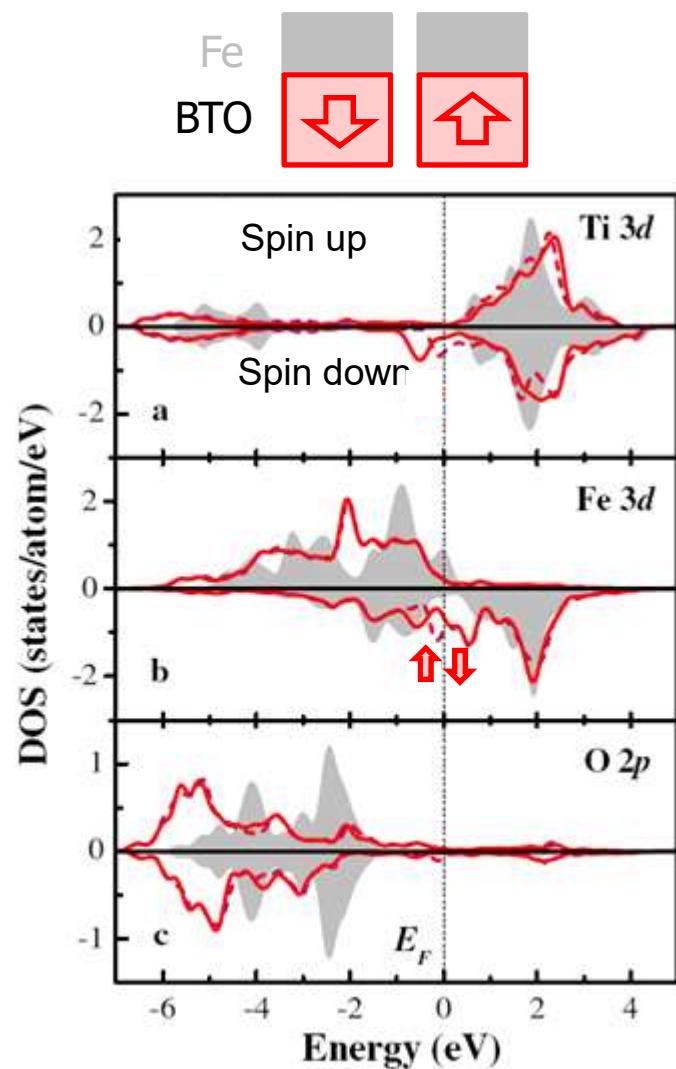
Electric-field control of spin polarization



- Change of spin polarization of Fe depending on ferroelectric polarization direction
- Probe this effect in Fe/BTO/LSMO tunnel junctions

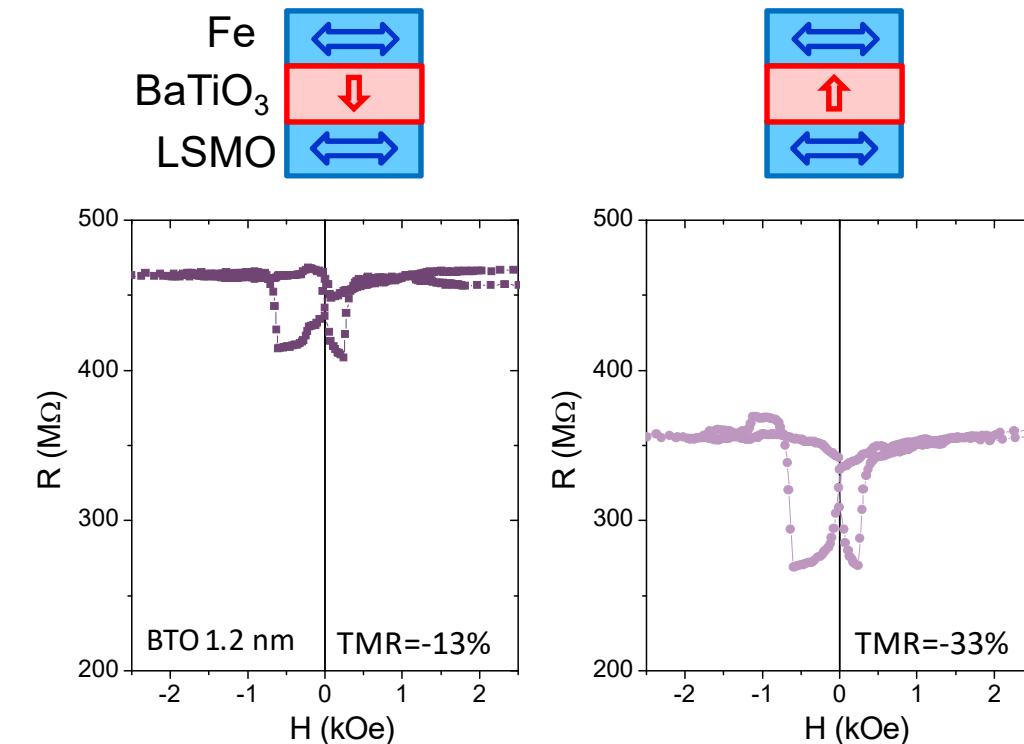
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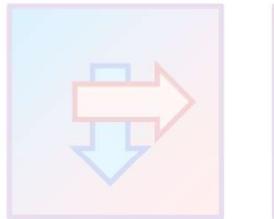
- TMR amplitude depends on direction of P
- **Ferroelectric control of spin polarization**
- Combination of field-effect and hybridization changes

V. Garcia, MB et al, Science 327, 1106 (2010)

S. Valencia, MB et al, Nature Mater. 10, 753 (2011)

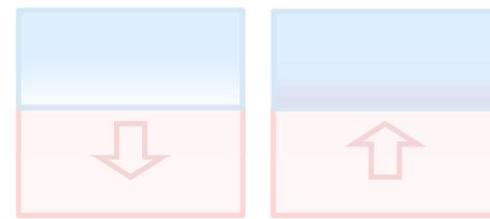
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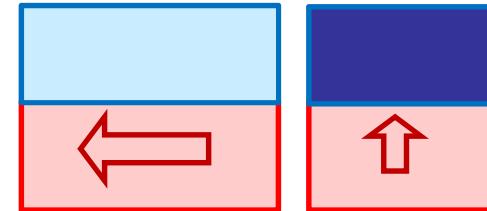
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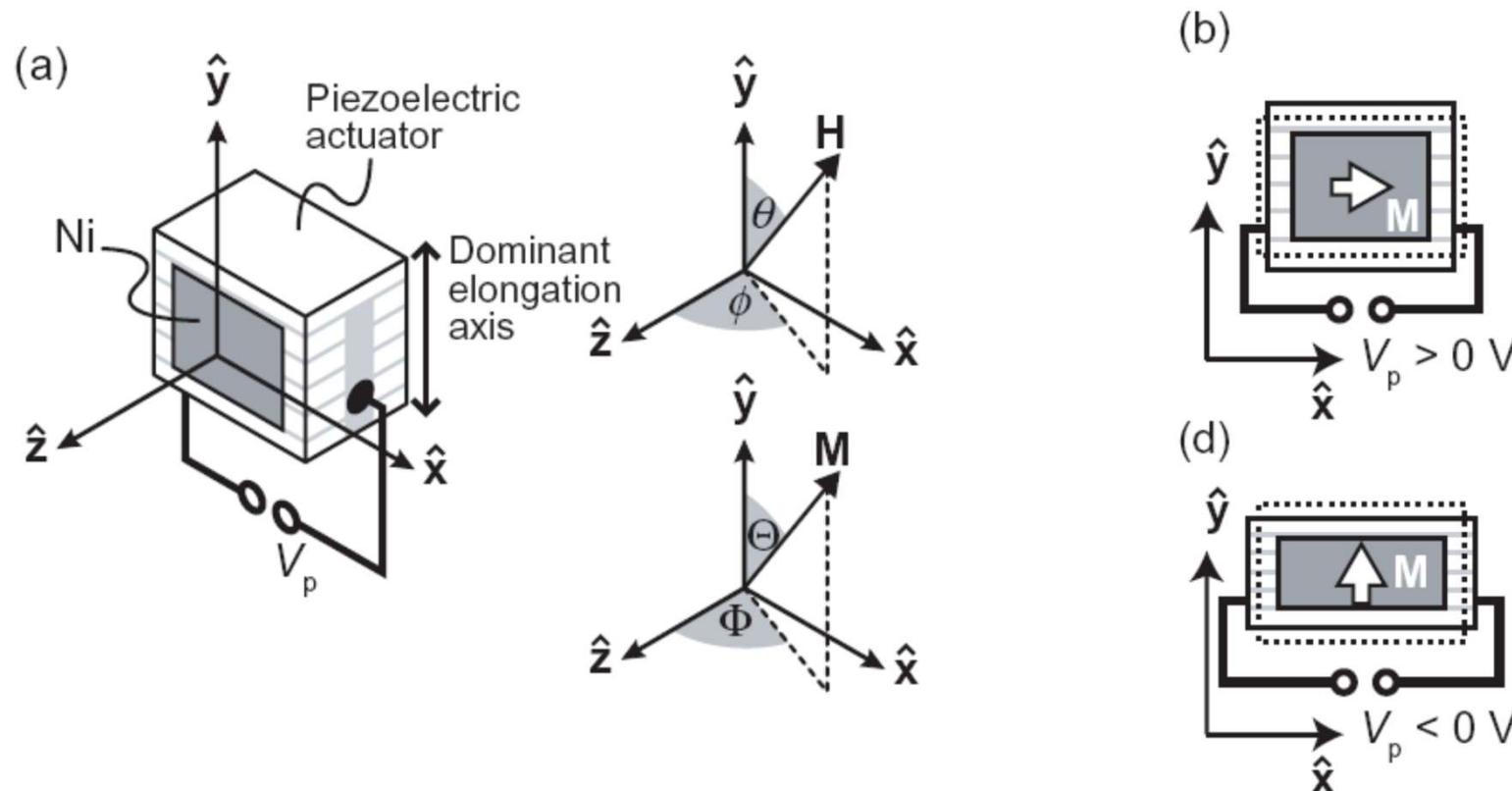
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Strain-induced control of magnetic anisotropy

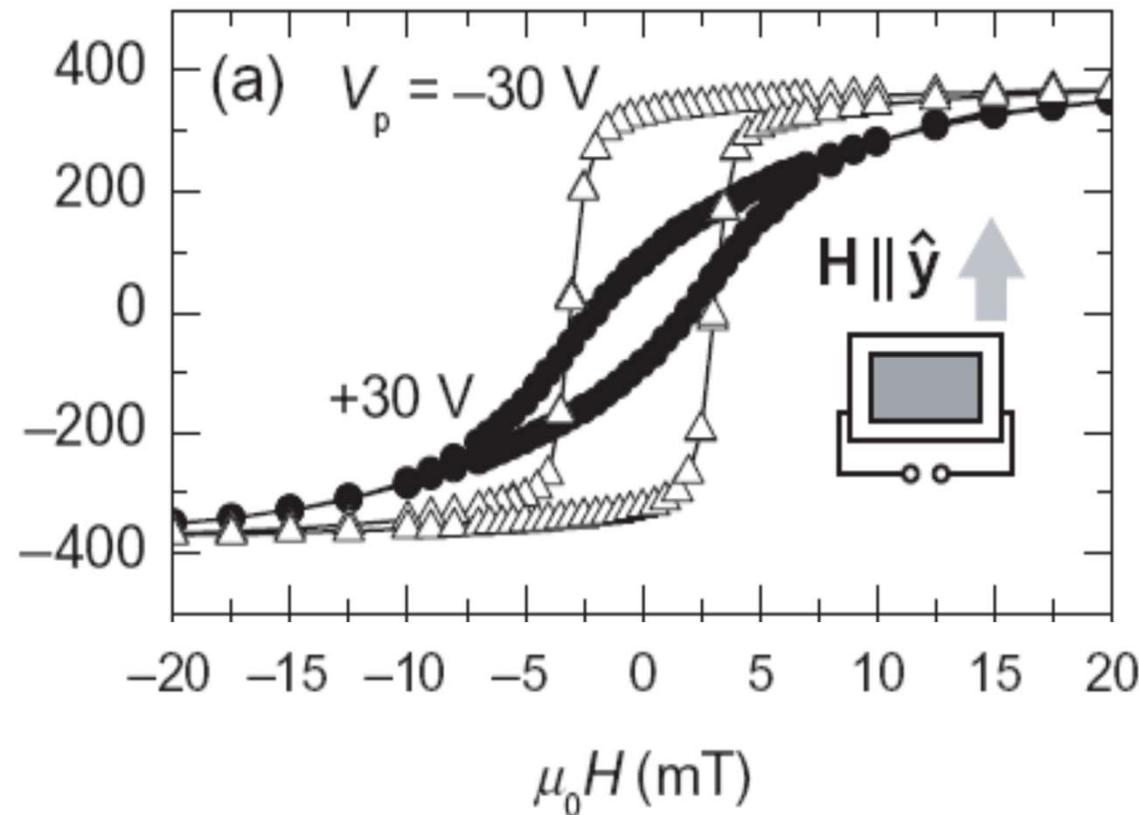
Example : experiments on PZT/Ni Weiler et al, New J. Phys 2009

Principle :

- E-field applied to PZT : change in PZT dimensions due to **converse piezoelectric effect**
- Change in dimensions induced in Ni : strain effect
 - Due to **magnetostriction** in Ni, strain modifies the magnetic properties



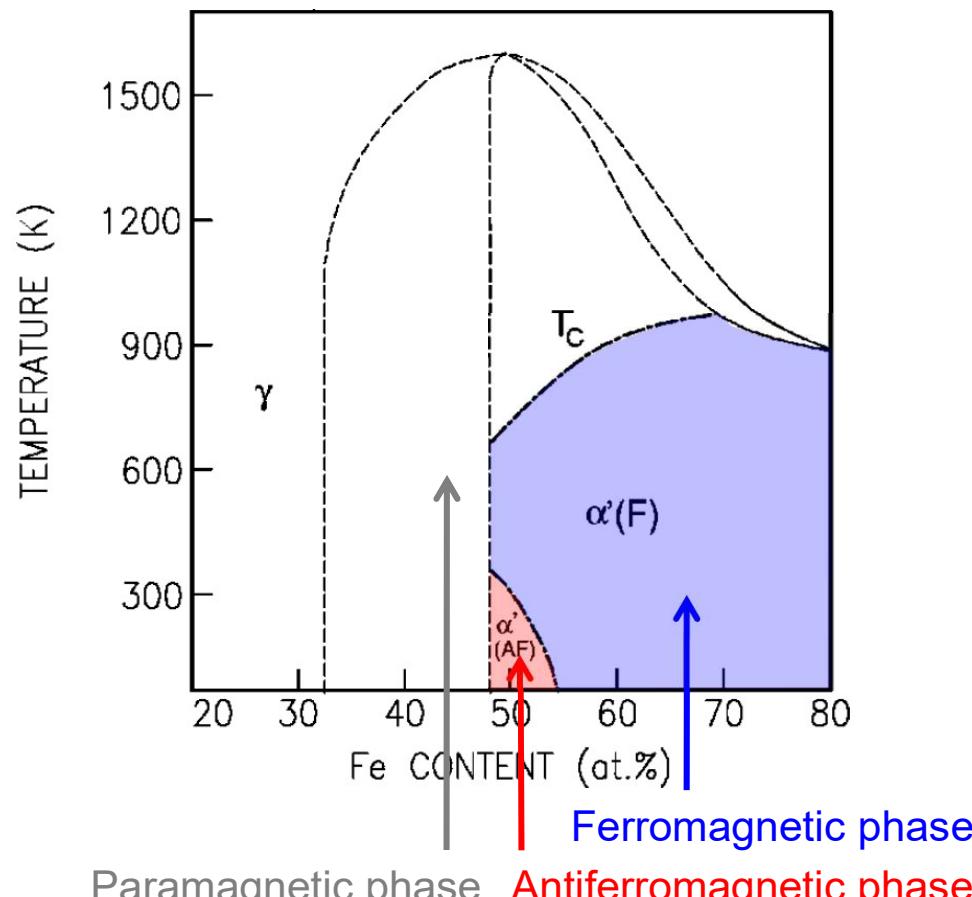
Strain-induced control of magnetic anisotropy



Weiler et al, New J. Phys 2009

Electric-field induced control of magnetization easy axis

Strain-induced control of magnetic order

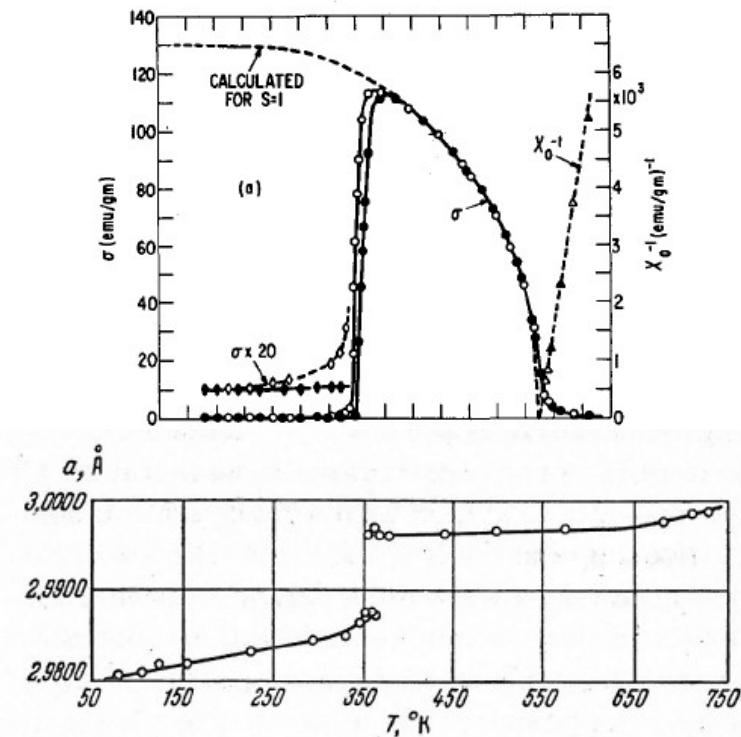
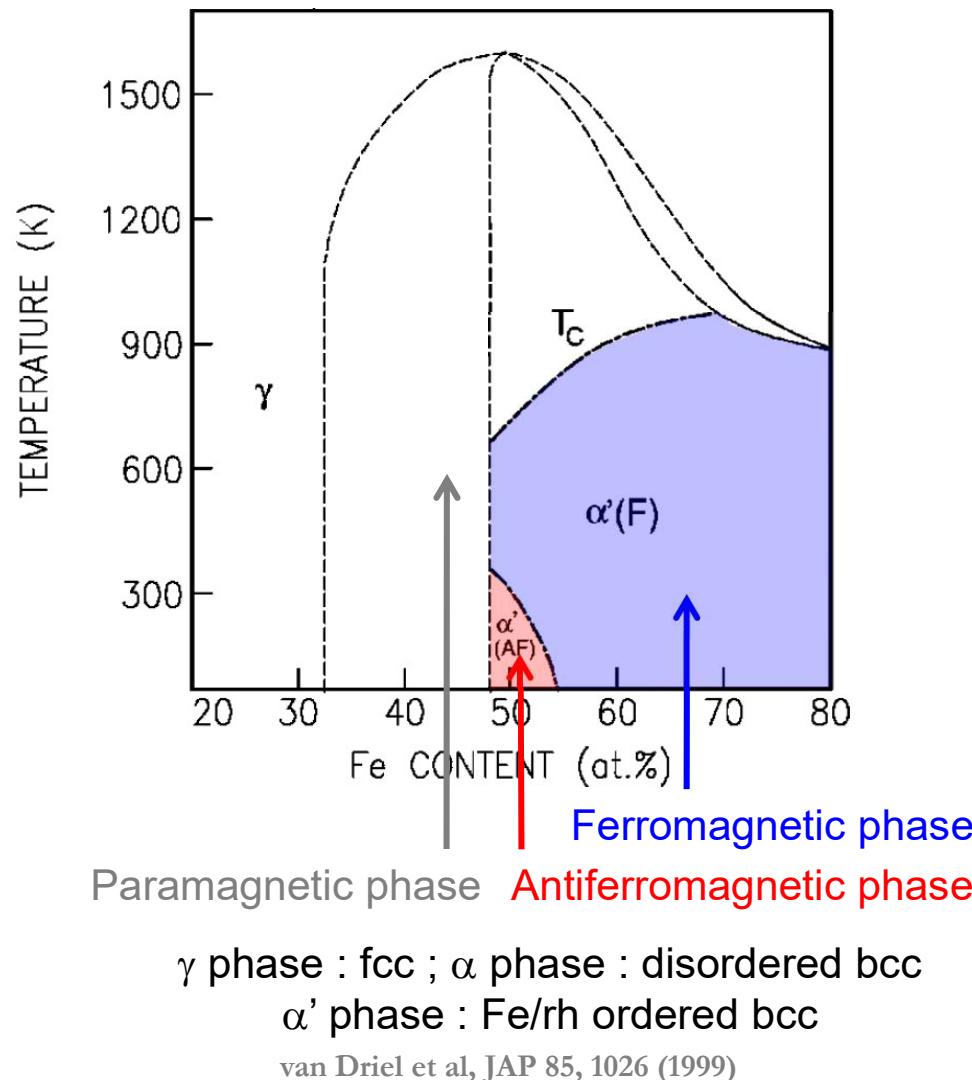


γ phase : fcc ; α phase : disordered bcc

α' phase : Fe/rh ordered bcc

van Driel et al, JAP 85, 1026 (1999)

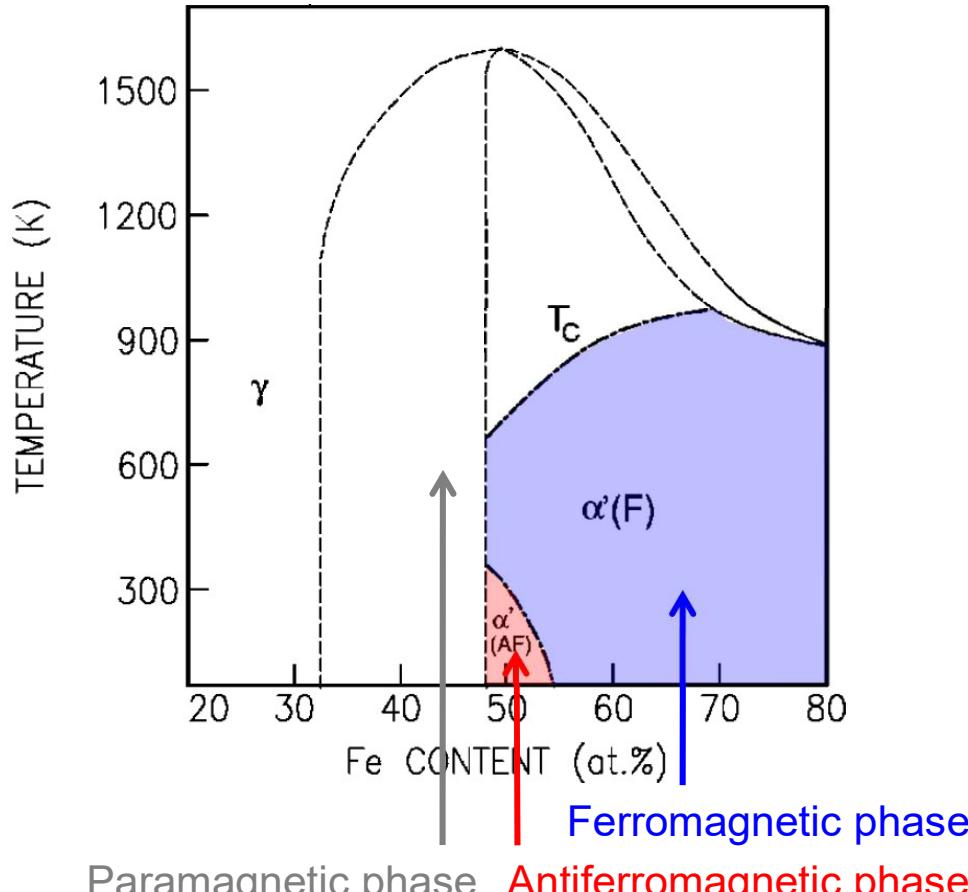
Strain-induced control of magnetic order



- ➊ Near $\text{Fe}_{50}\text{Rh}_{50}$, transition from AFM to FM at about 370K
- ➋ Transition is first order
- ➌ Associated large resistivity drop
- ➍ Jump of cell volume by ~1% at T^*

Zakharov et al, Sov. Phys. JETP 19, 1348 (1964)

Strain-induced control of magnetic order



Paramagnetic phase **Antiferromagnetic phase**

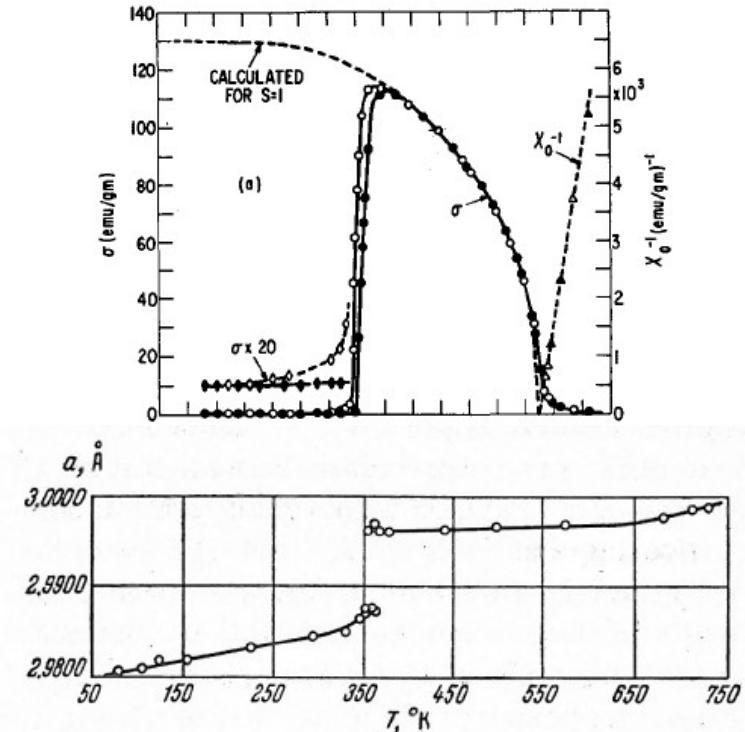
γ phase : fcc ; α phase : disordered bcc

α' phase : Fe/rh ordered bcc

van Driel et al, JAP 85, 1026 (1999)

○ Magnetic state of FeRh is sensitive to pressure

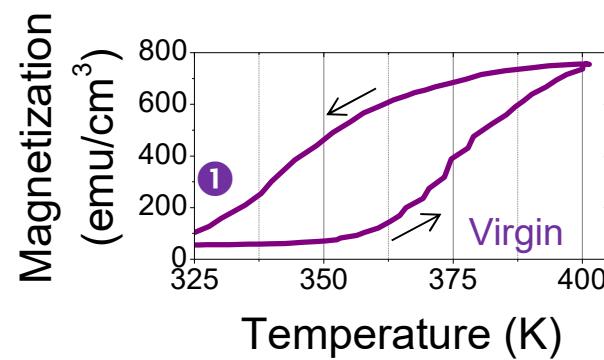
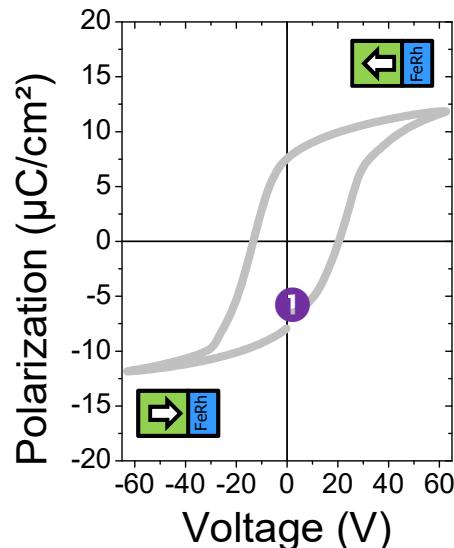
→ Grow on **ferroelectric/ferroelastic BaTiO₃** substrate to achieve E-field control



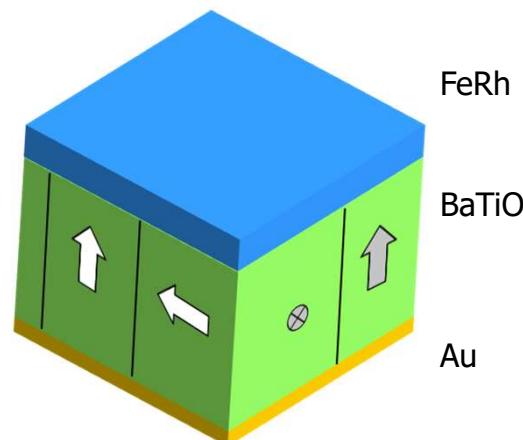
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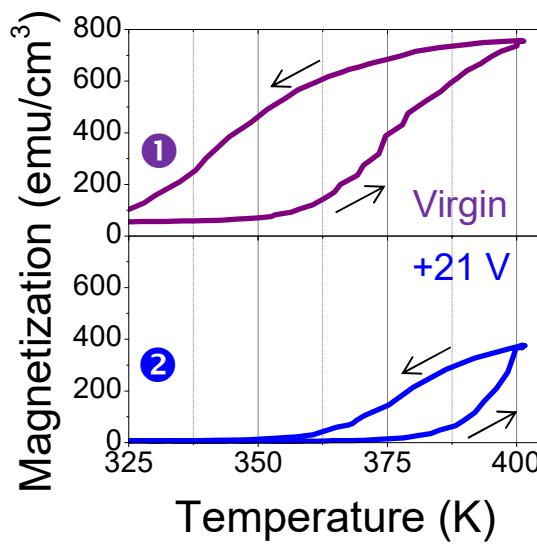
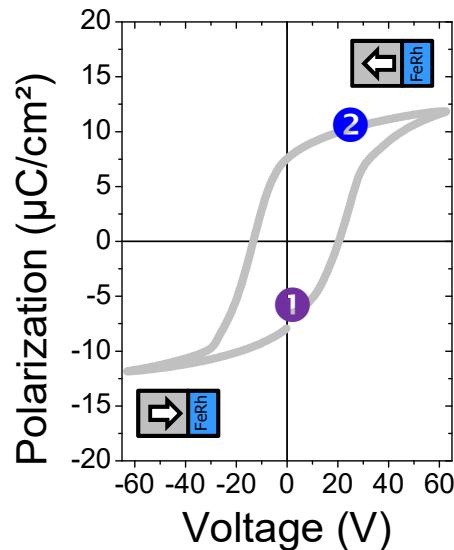
Strain-induced control of magnetic order



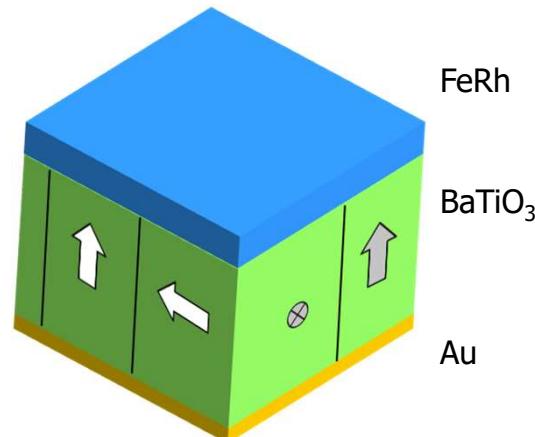
○ At 0V at 20 kOe, $T^*\approx 360$ K



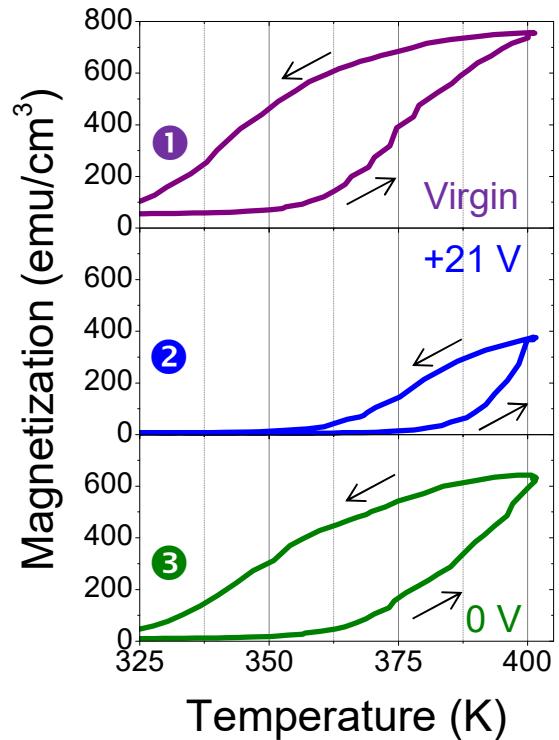
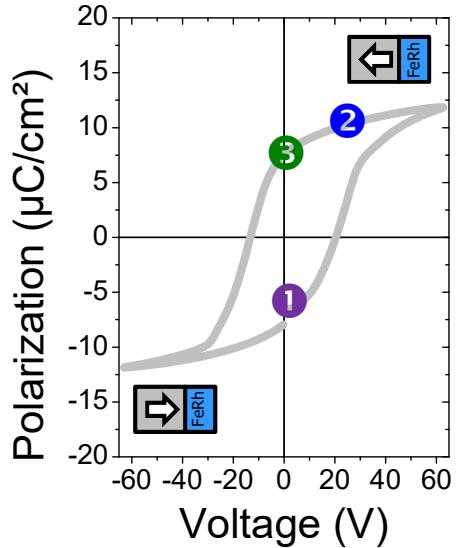
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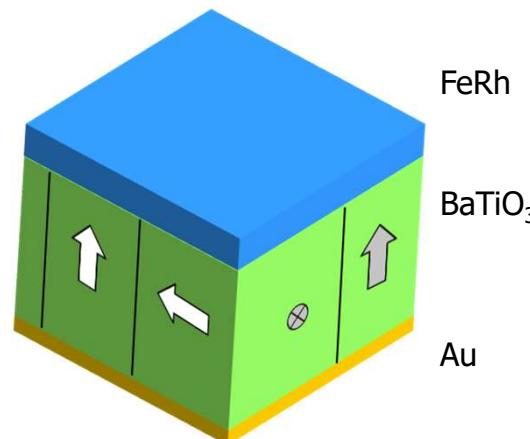
- At 0V at 20 kOe, $T^* \approx 360$ K
- Voltage shifts T^* by ~ 20 K



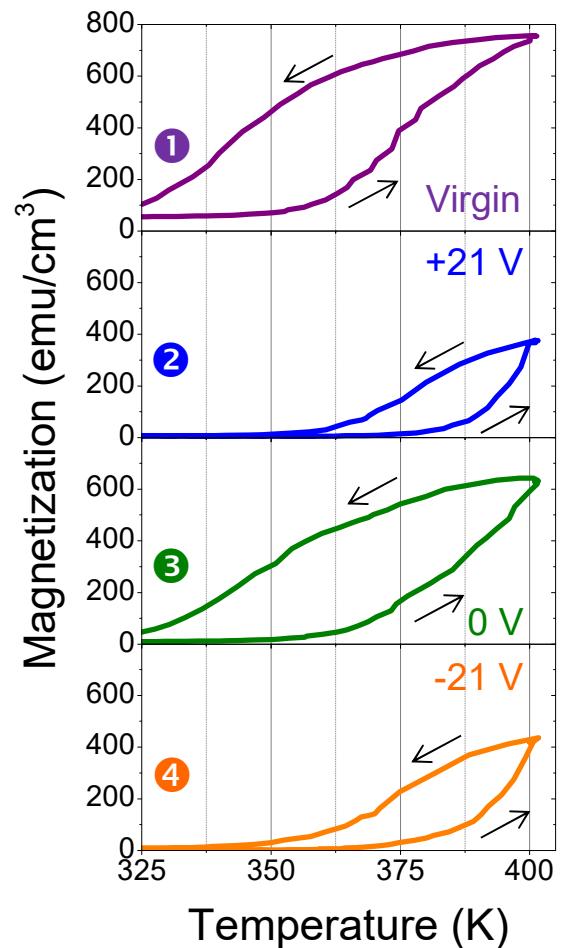
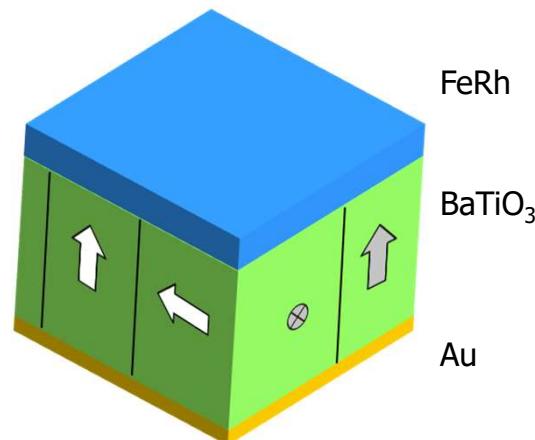
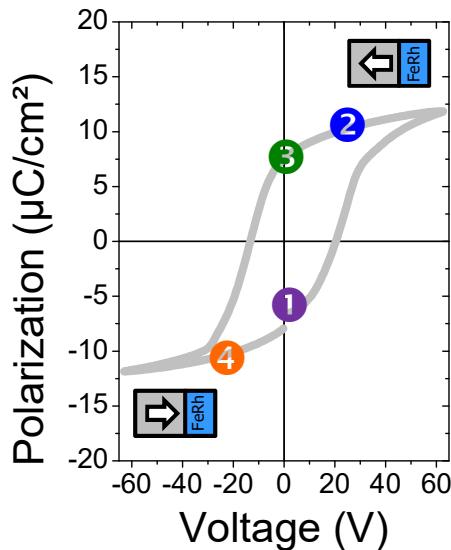
Strain-induced control of magnetic order



- At 0V at 20 kOe, $T^* \approx 360$ K
- Voltage shifts T^* by ~ 20 K
- Effect is reversible

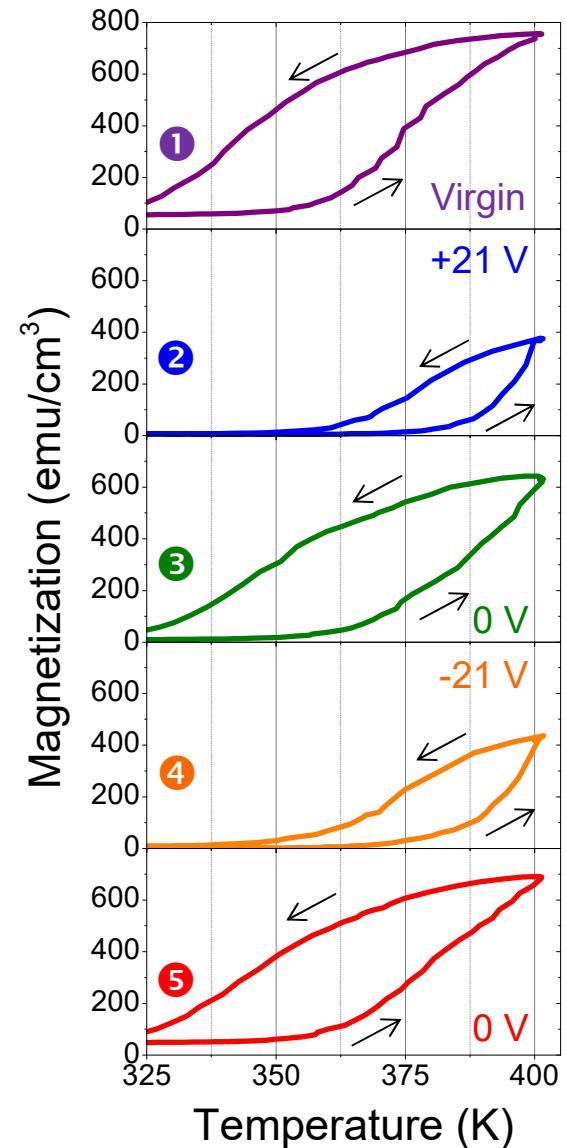
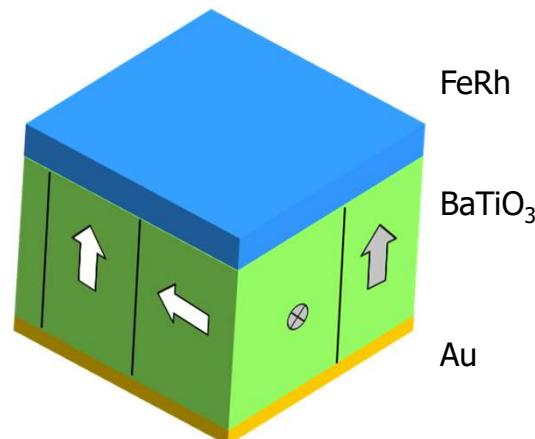
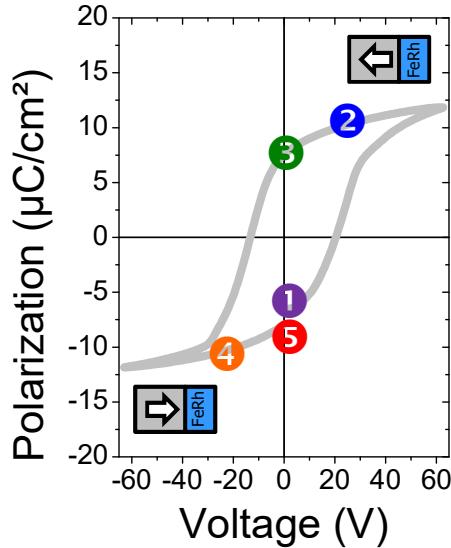


Strain-induced control of magnetic order



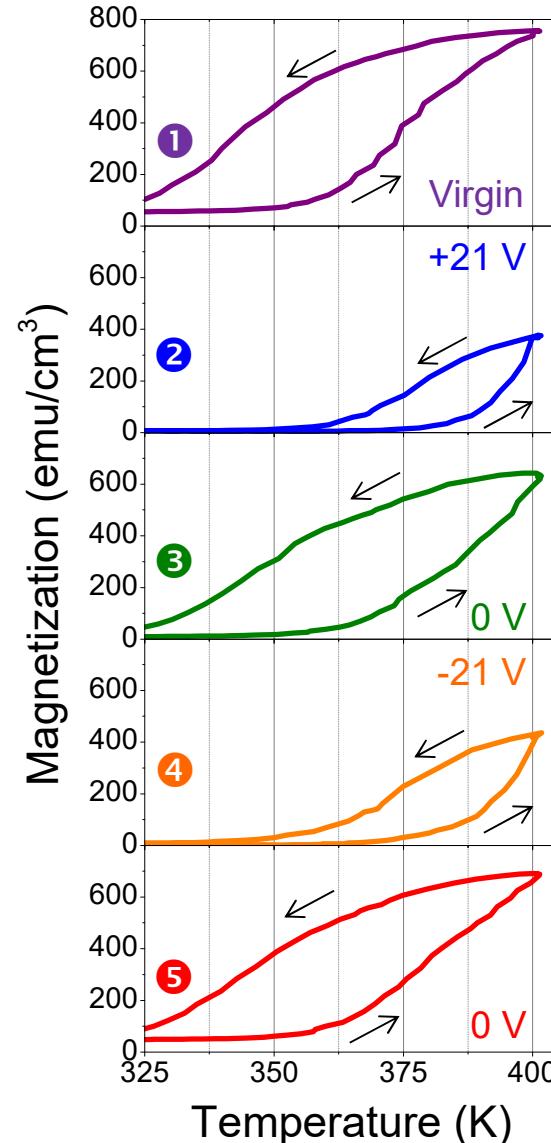
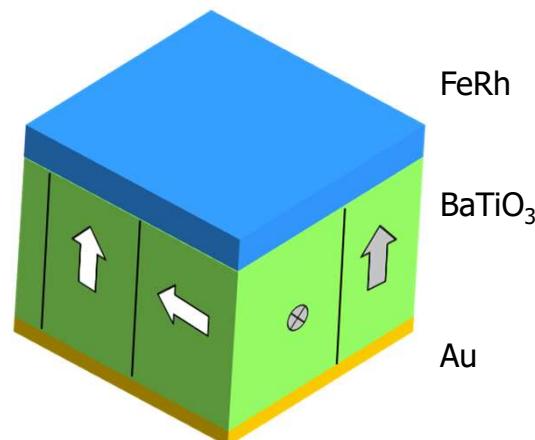
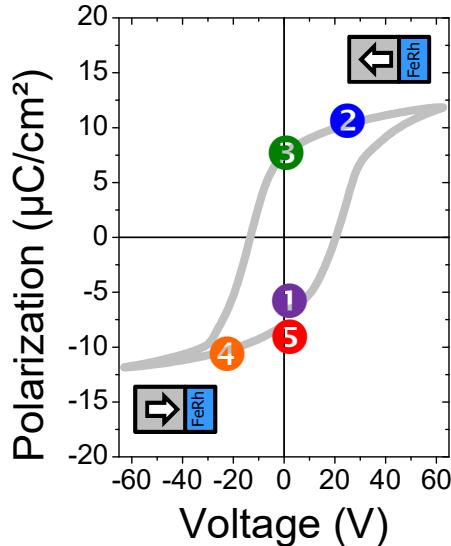
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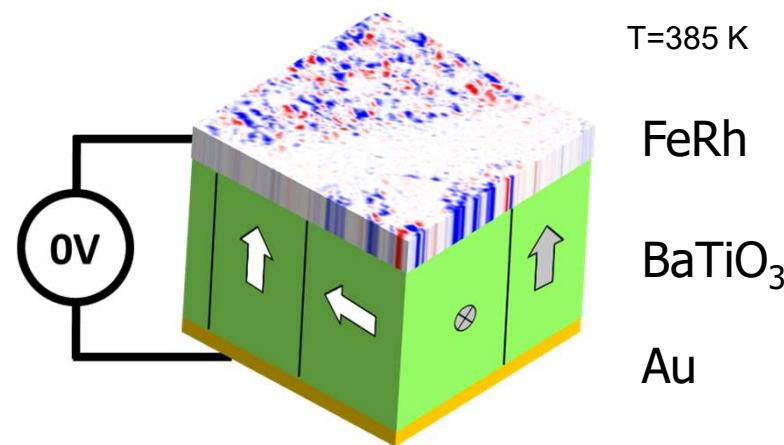
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- Max magnetization change $\sim 600 \text{ emu}/\text{cm}^3$
- ME coupling $\alpha = 1.6 \cdot 10^{-5} \text{ s/m}$
- Larger than in any single phase material by 5 orders
- Larger than in any artificial multiferroic by factor > 10

Cherifi, MB et al, Nature Mater. 13, 345 (2014)

Strain-induced control of magnetic order

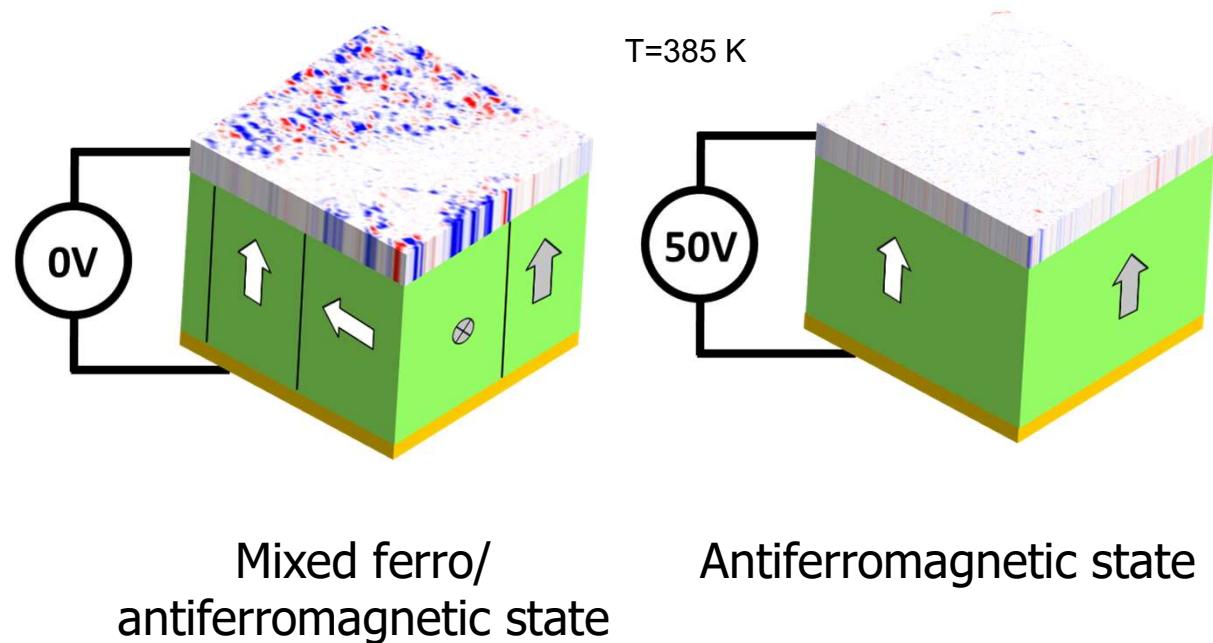
Direct imaging of magnetic state using XCMD-PEEM



Mixed ferro/
antiferromagnetic state

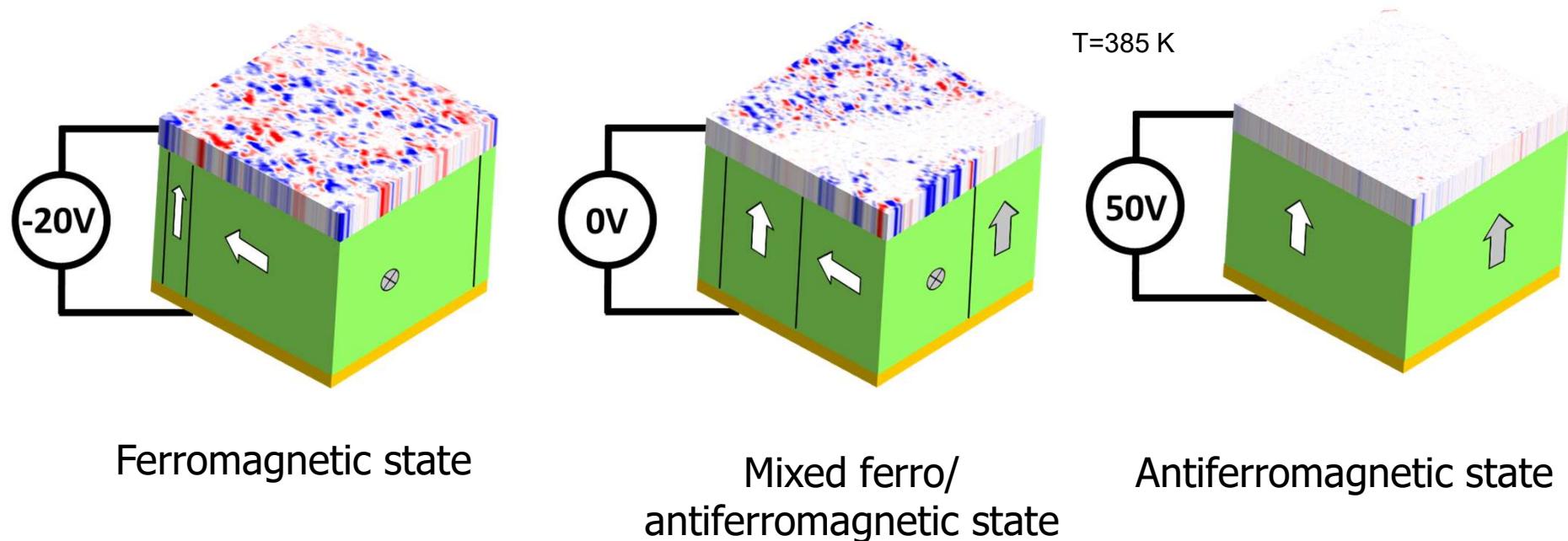
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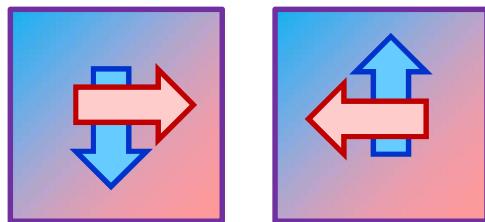
Direct imaging of magnetic state using XCMD-PEEM



- Switch ferromagnetism **OFF and ON** by electric field, just above room temperature

Phillips, MB et al, Sci. Rep. 5, 10026 (2014)

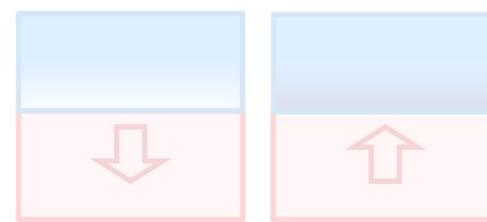
Intrinsic magnetolectric



Use single-phase multiferroic material

- ✓ Simple approach, just one material
- ✓ Beautiful physics, potential for new science
- ✗ BFO only RT multiferroic
- ✗ Can be leaky, hard to switch

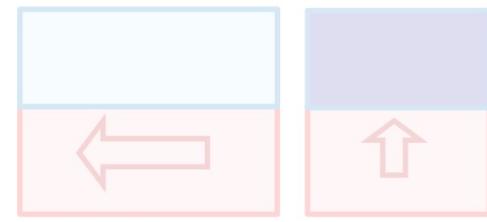
Field-effect



Combine strong ferroelectric with carrier-mediated ferromagnet

- ✓ Broader choice of materials
- ✓ Well-suited for perpendicular transport
- ✗ Few ferromagnetic oxides with high T_c ; need simple metals
- ✗ Effect occurs over very small thickness (few nm max)
- ✗ Needs very large fields

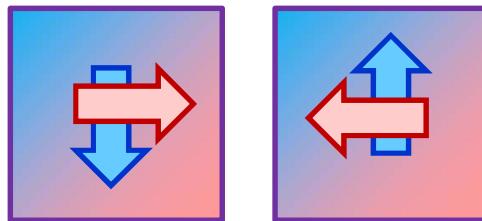
Strain-driven



Combine piezoelectric or ferroelectric/ferroelastic with magnetostrictive ferromagnet

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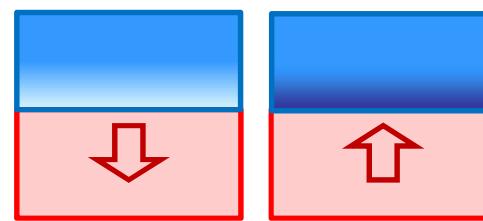
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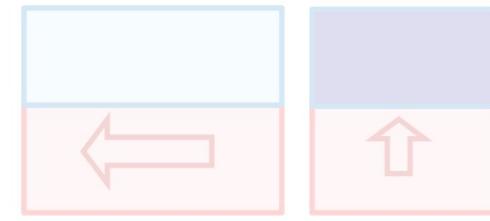
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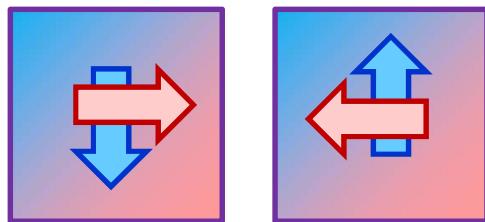
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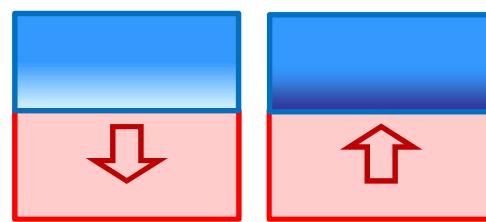
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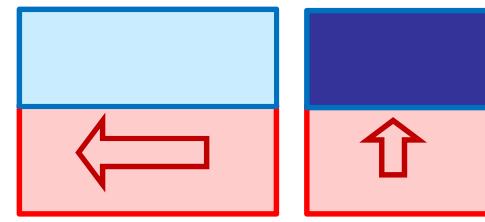
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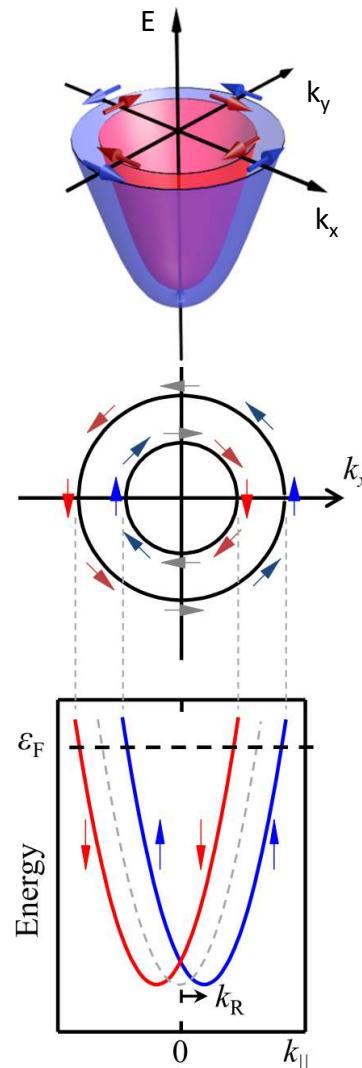


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- Ferroic orders, multiferroics and magnetoelectric coupling
- Approaches for the electric-field control of magnetism
- ➔ ○ Electric-field control of spin-charge interconversion
- Low-power spin-based devices for logic-in-memory

Rashba and Edelstein effects



The **Rashba effect** – manifestation of spin-orbit interaction (SOI) in solids, more particularly in two-dimensional electron systems, where spin degeneracy is lifted due to a symmetry-breaking electric field normal to an heterointerface.

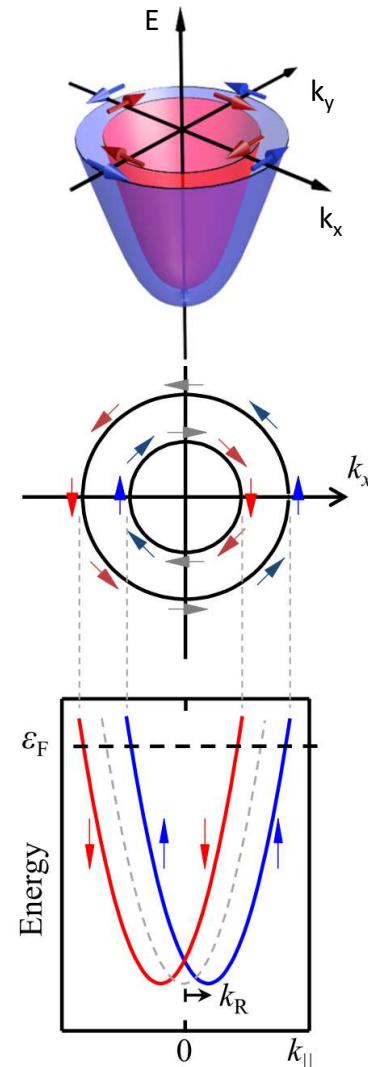
$$H_R = \frac{\alpha_R}{\hbar} (\vec{\sigma} \wedge \vec{p}) \bullet \vec{z}$$

Bychkov & Rashba, Sov. Phys. JETP Lett. 39, 78 (1984)

Manchon et al., Nature Mater. 14, 871 (2015)

$$\varepsilon_k = \frac{\hbar^2 k^2}{2m^*} \pm \alpha_R k$$

Rashba and Edelstein effects

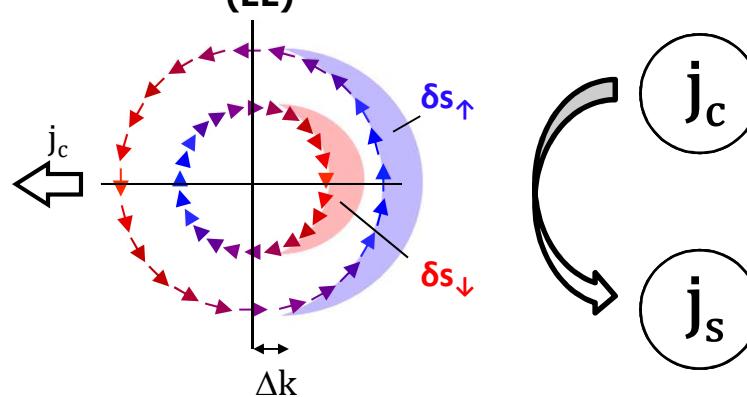


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Direct Edelstein effect (EE)

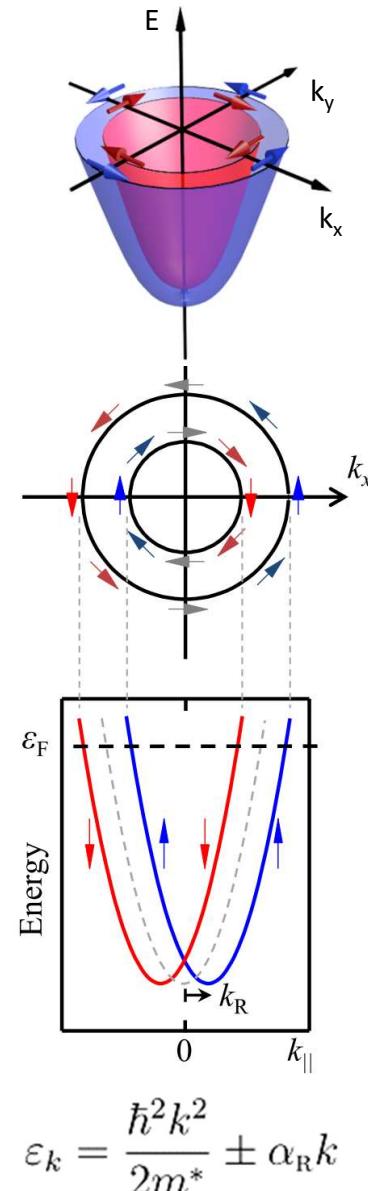


Injection of a **charge current**
Generation of **spin accumulation**

Edelstein, Solid State Commun. 73, 233 (1990)
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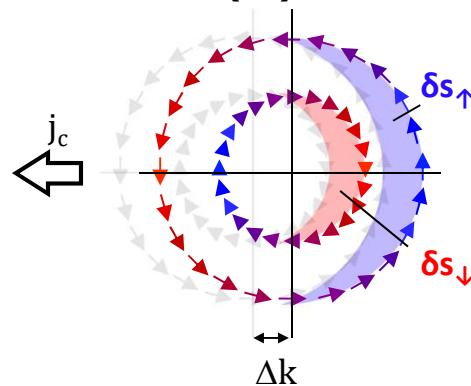


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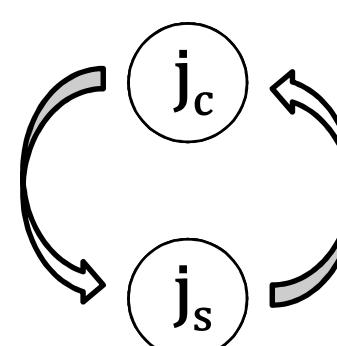
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Edelstein, Solid State Commun. 73, 233 (1990)
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Inverse Edelstein effect (IEE)

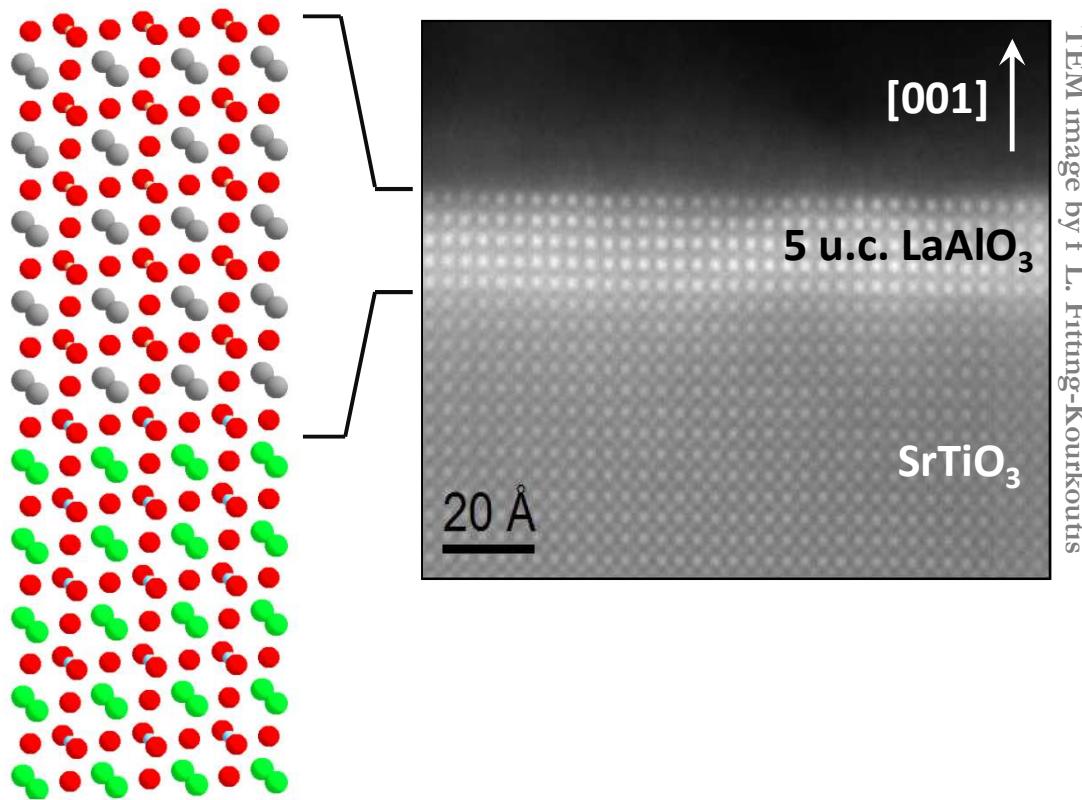


Injection of a **spin current**
Generation of a **charge current**

Rojas-Sánchez et al., Nat. Commun. 4, 2944 (2013)
Shen et al., Phys. Rev. Lett. 112, 096601 (2014)

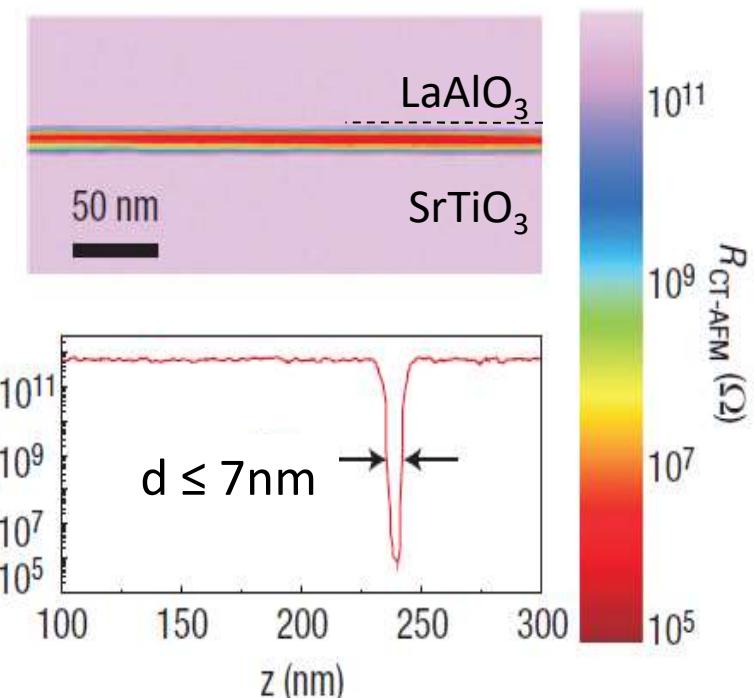
Figure of merit : $\lambda_{IEE} = \frac{j_c^{2D}}{j_s^{3D}} = \frac{\alpha_R \tau}{\hbar}$

The LaAlO₃/SrTiO₃ system



TEM image by f L. Fitting-Kourkoutis

Resistance map in cross-section (CT-AFM)

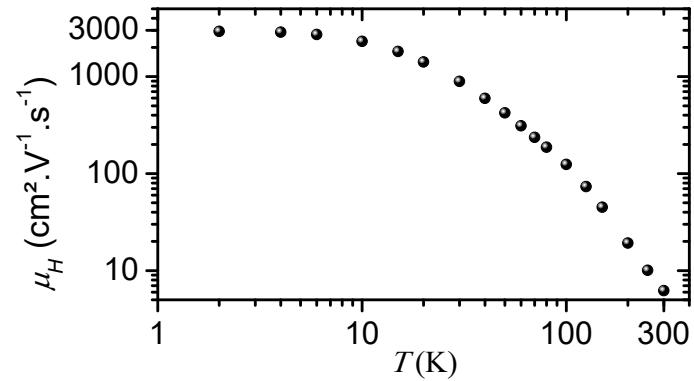


➔ Quasi-two-dimensional
conduction

M. Basletic, MB et al., Nat. Mater. 7, 621 (2008)
O. Copie, MB et al., Phys. Rev. Lett. 102, 216804 (2009)

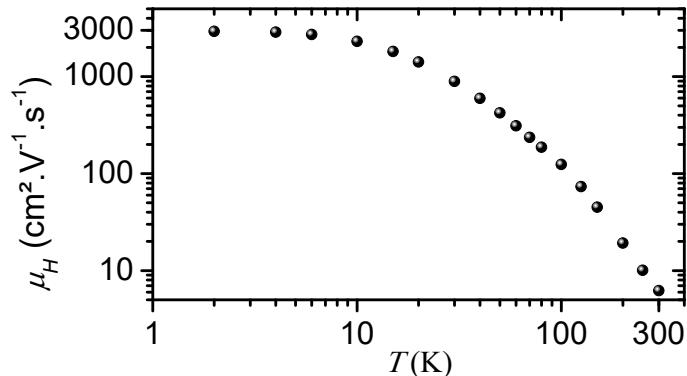
The LaAlO₃/SrTiO₃ system

⦿ High Electron Mobility

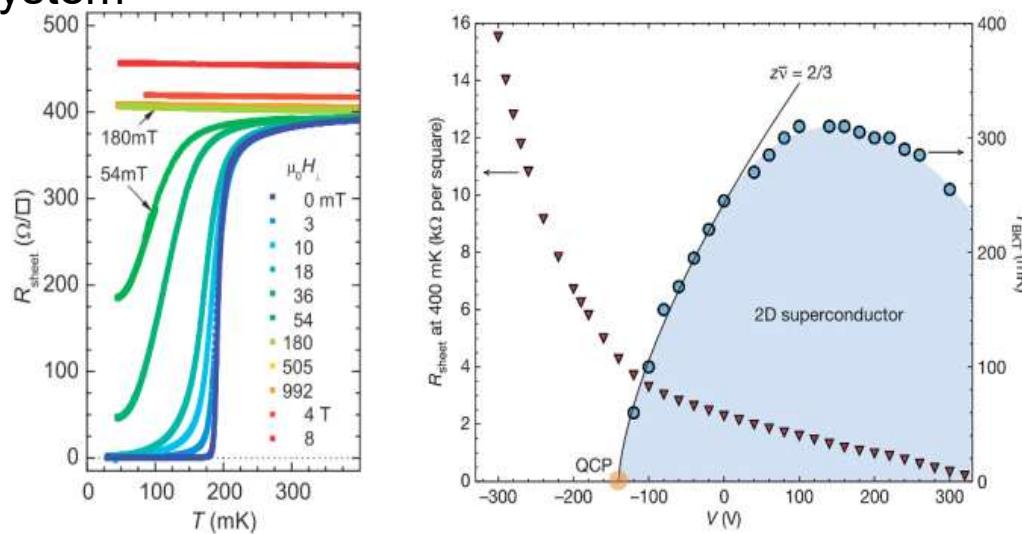


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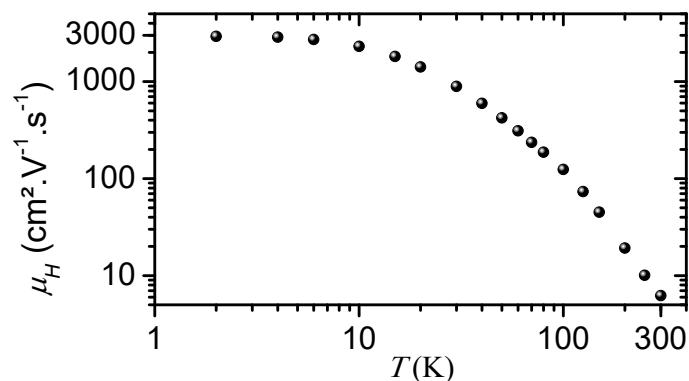
⦿ Two-dimensional Superconductivity + gateable system



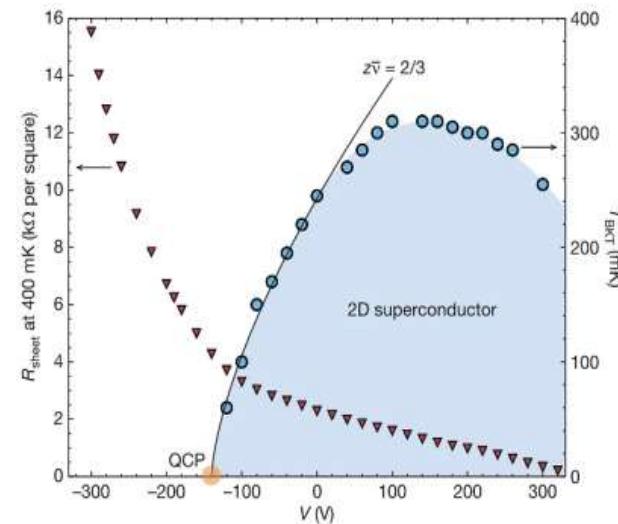
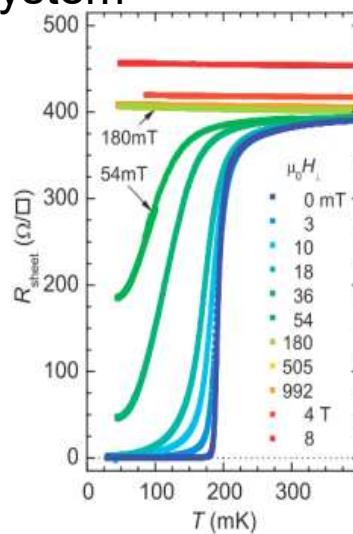
N. Reyren et al., Science 317, 1196 (2007) ; A. D. Caviglia et al., Nature 456, 624 (2008)

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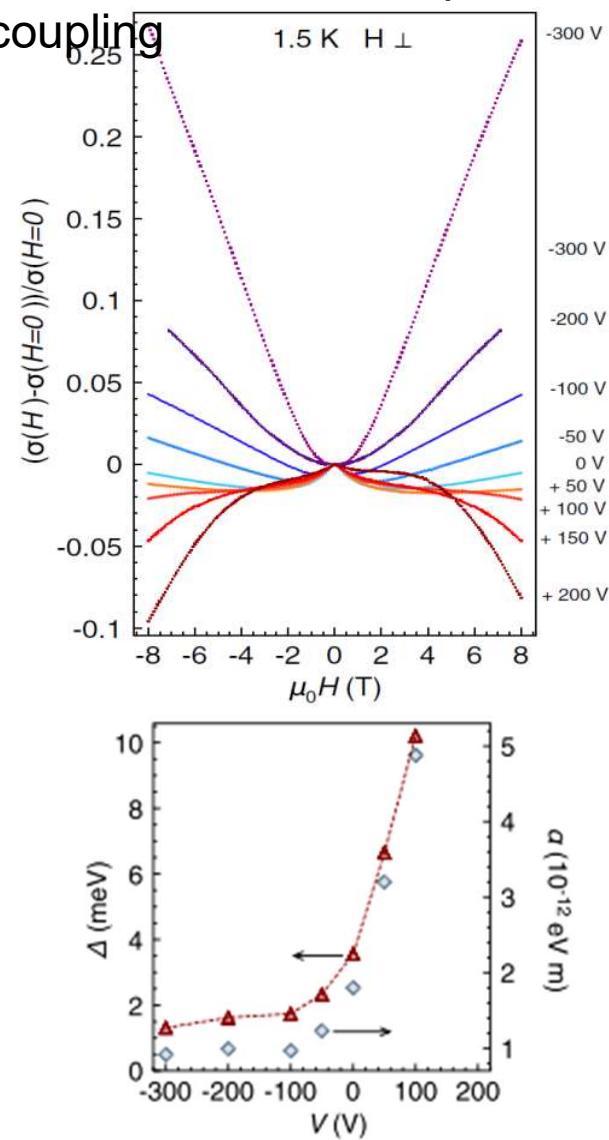


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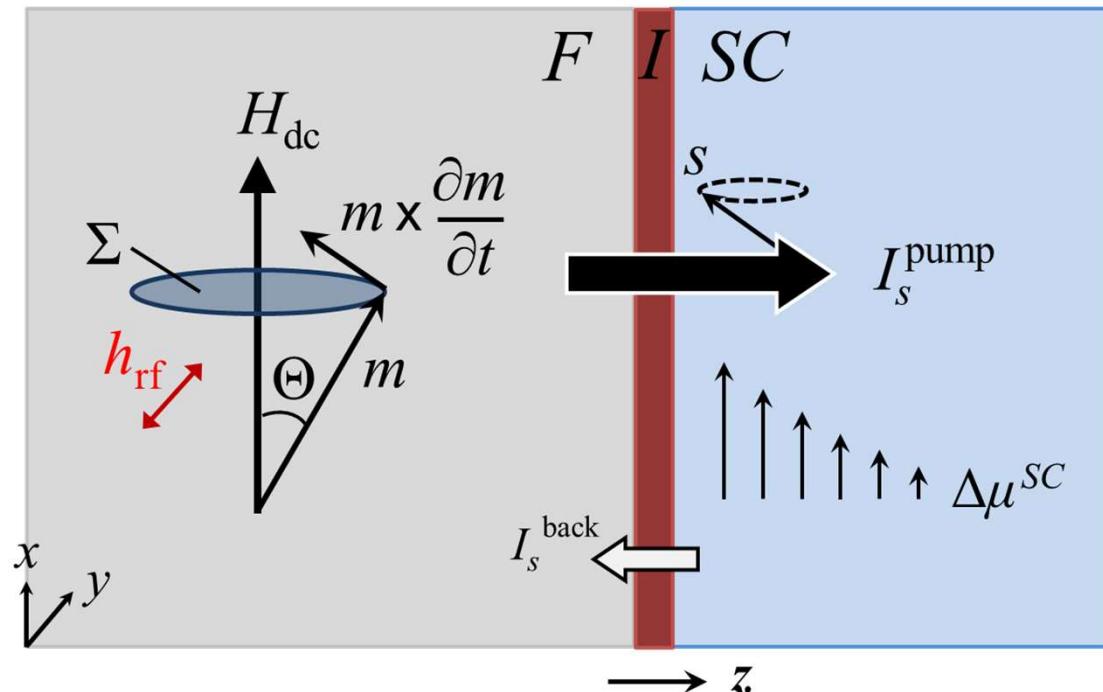
○ Sizeable Rashba spin-orbit coupling



A.D. Caviglia et al., PRL. 104, 126803 (2010)

Spin pumping

- An microwave excitation field drives magnetization dynamics
- Ejects transverse spin angular momentum
- Creates an imbalance of spin up and spin down populations
- ➔ Generation of a pure spin current I_s .

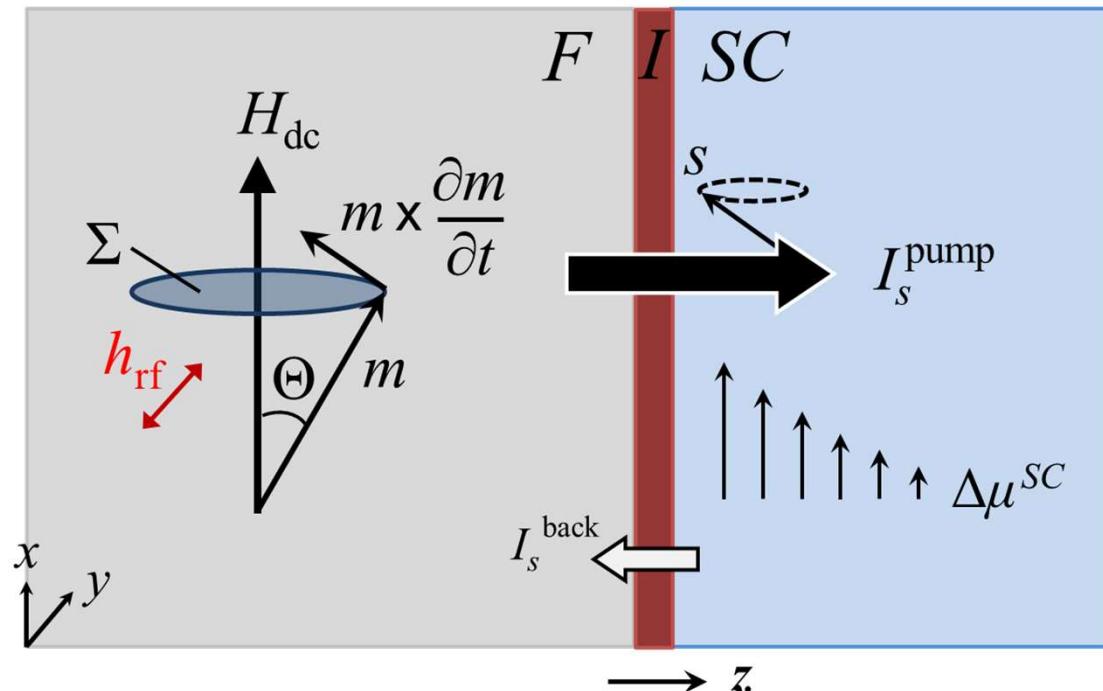


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A. Brataas et al., Phys. Rev. B 66, 060404(R) (2002)
Y. Tserkovnyak, et al., Rev. Mod. Phys. 77, 1375 (2005)

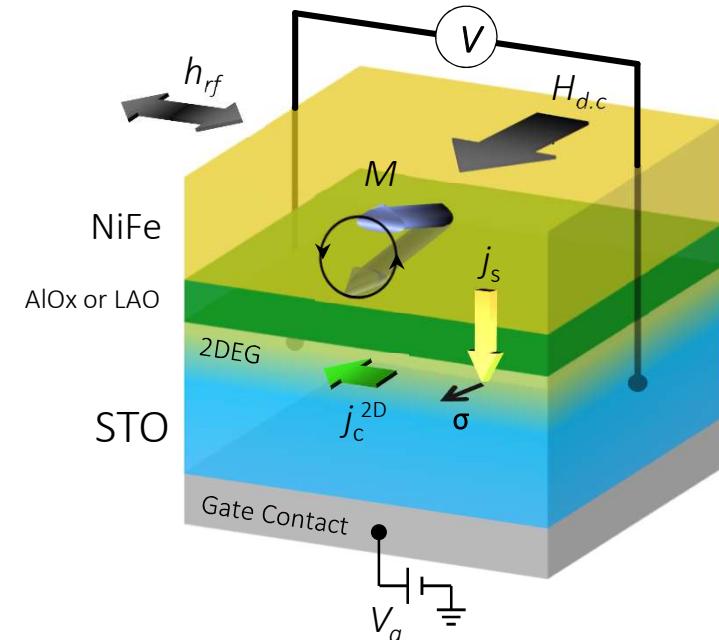
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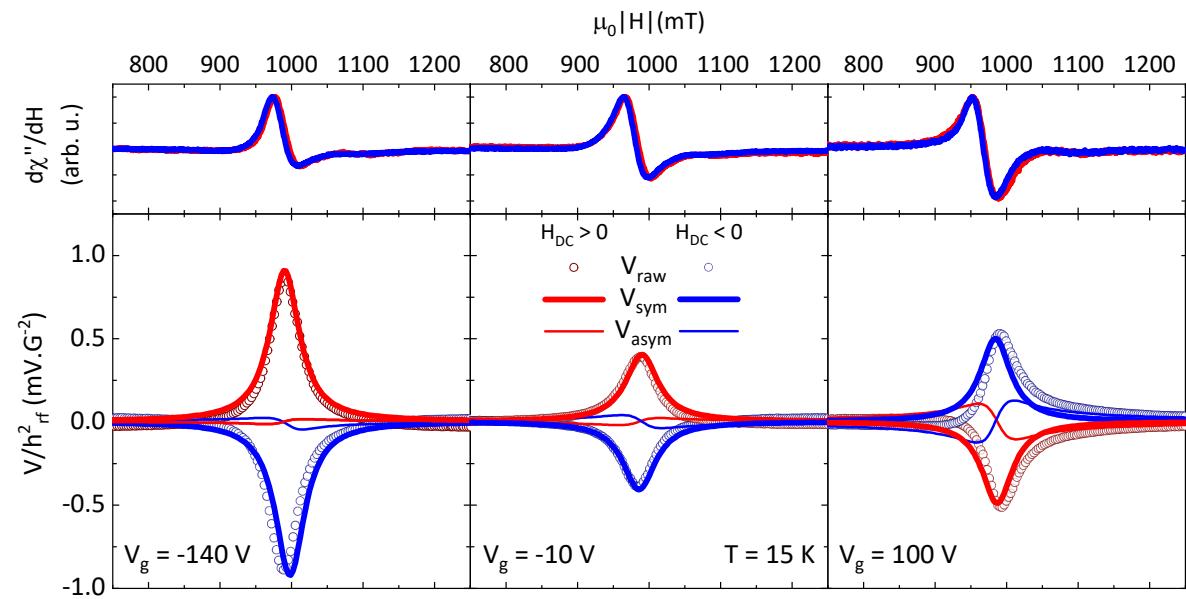
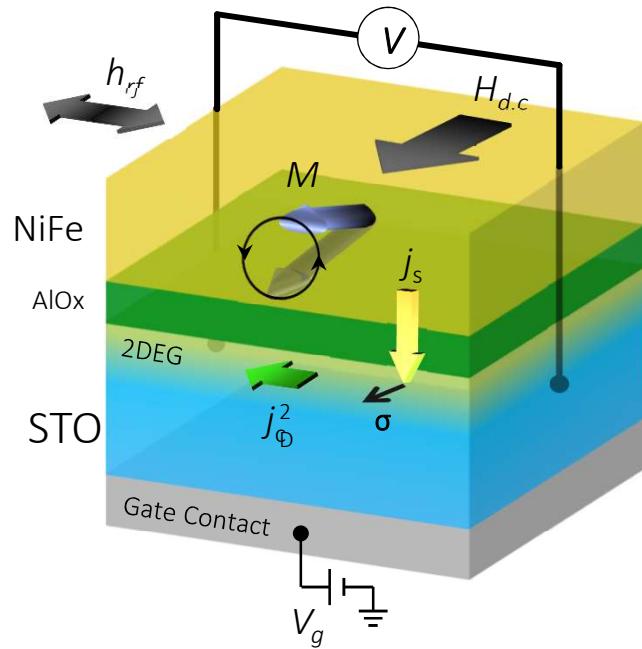


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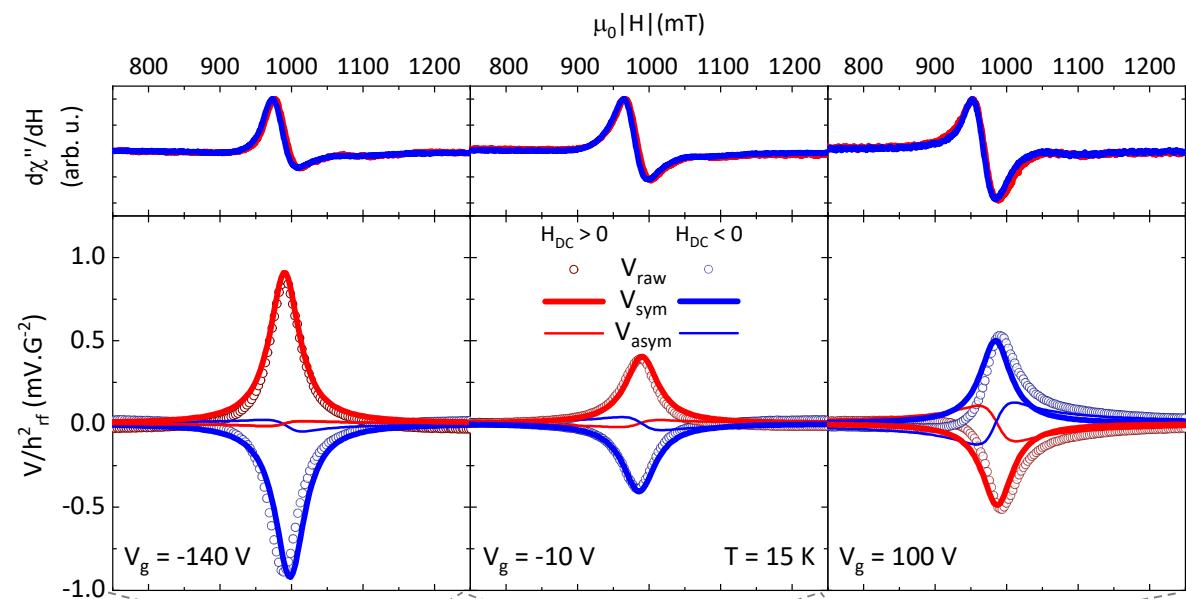
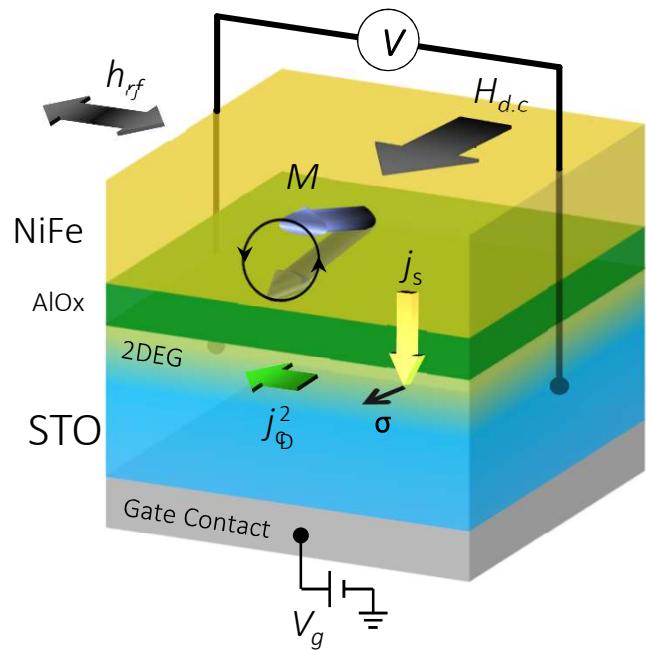


Spin-charge conversion in Al/STO 2DEGs



- ➊ Voltage peak almost entirely symmetric
- ➔ **Inverse Edelstein effect**

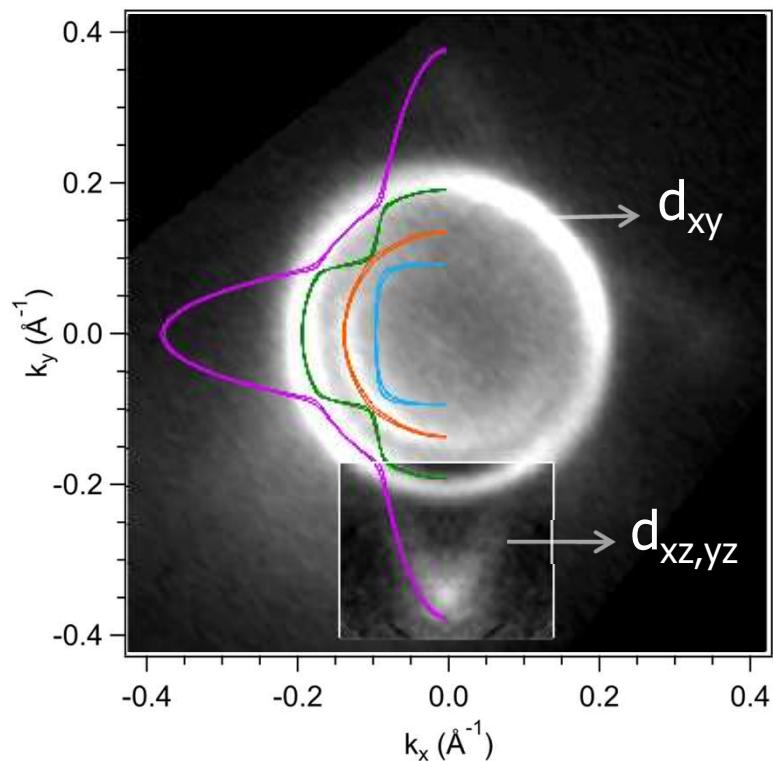
Spin-charge conversion in Al/STO 2DEGs



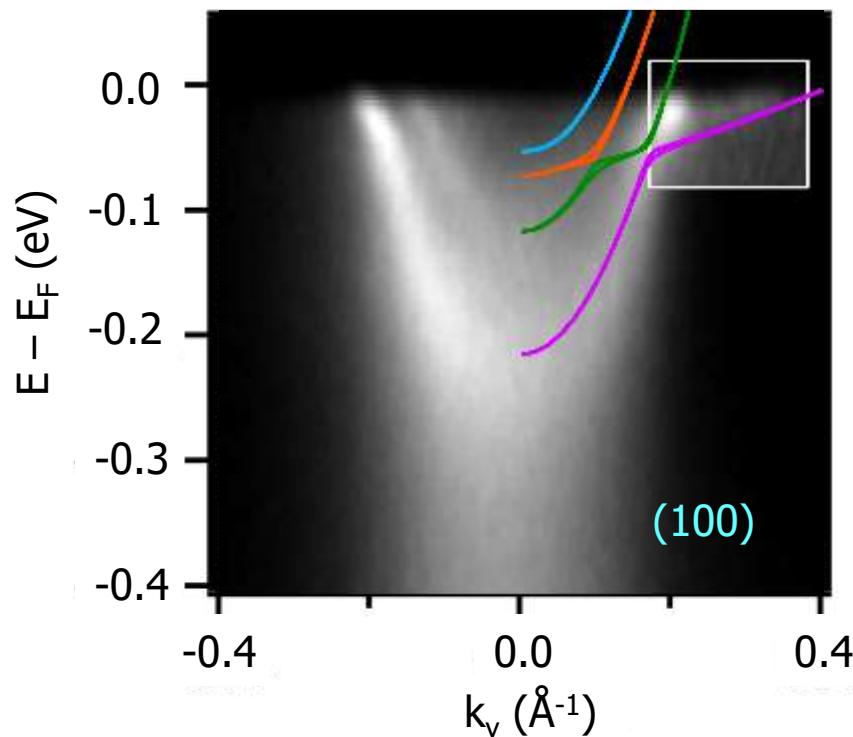
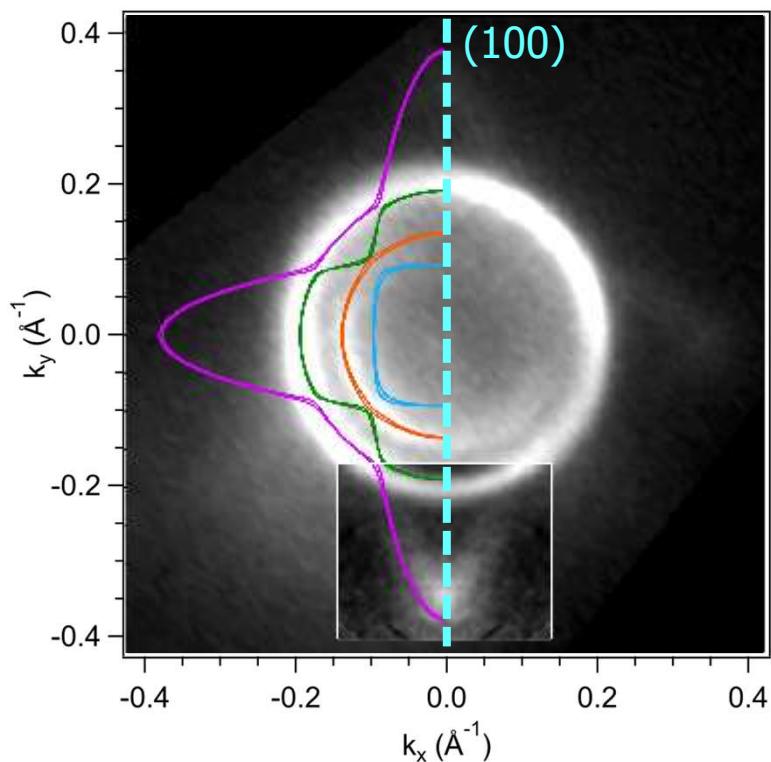
- ➊ Voltage peak almost entirely symmetric
- ➔ **Inverse Edelstein effect**
- ➋ Very large conversion effect with $\lambda_{\text{IEE}} > 25 \text{ nm}$
- ➋ Strong gate dependence of amplitude & sign
- ➋ Efficiency larger topological insulators and LAO/STO

Rojas-Sánchez et al., Nature Comm. 4, 2944 (2013)
 Rojas-Sánchez et al., Phys. Rev. Lett. 116, 096602 (2016)
 E. Lesne, MB et al., Nature Mater. 15, 1261 (2016)

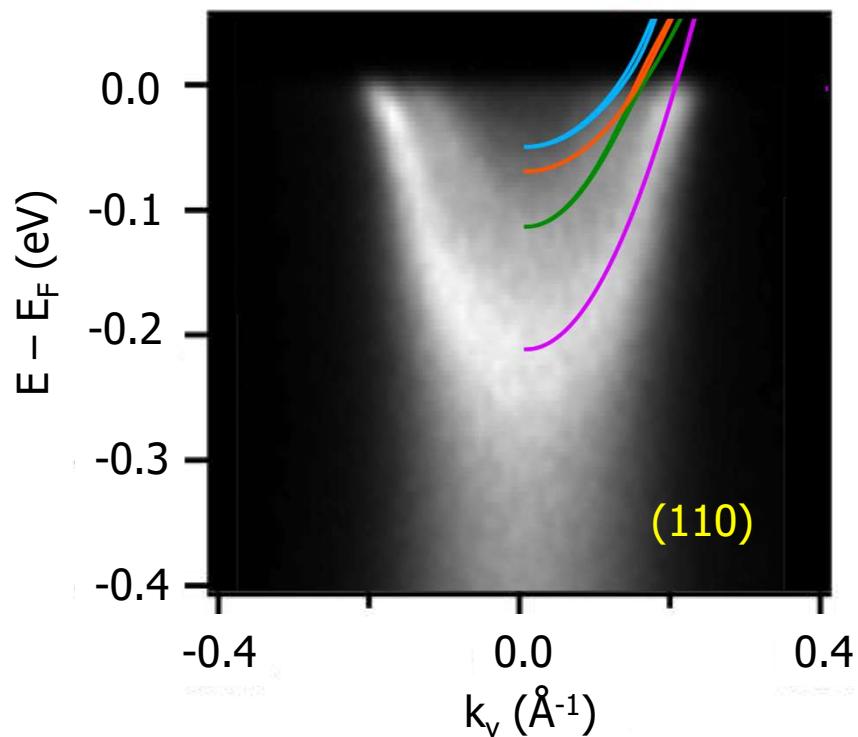
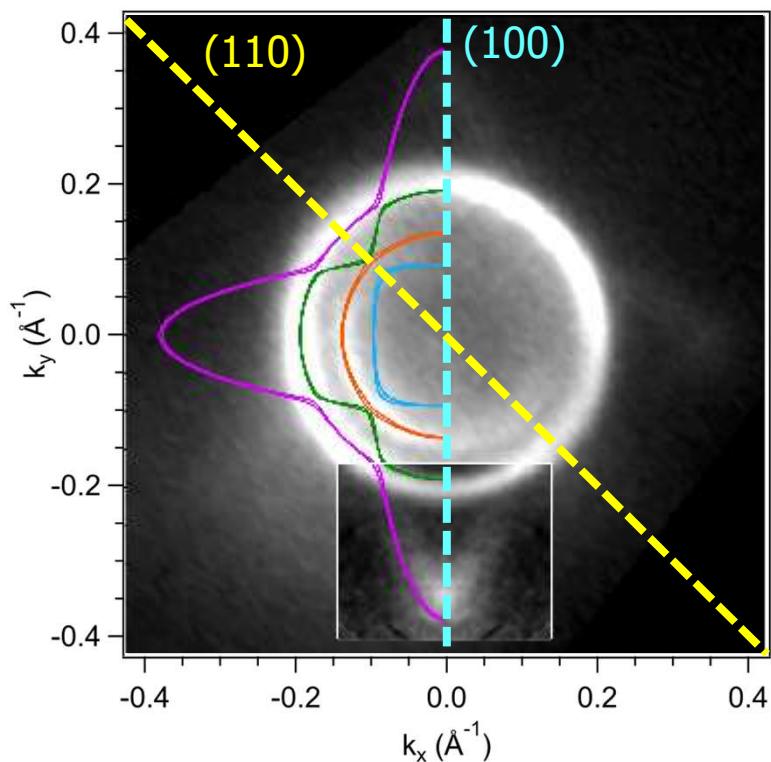
Electronic structure of STO 2DEGs



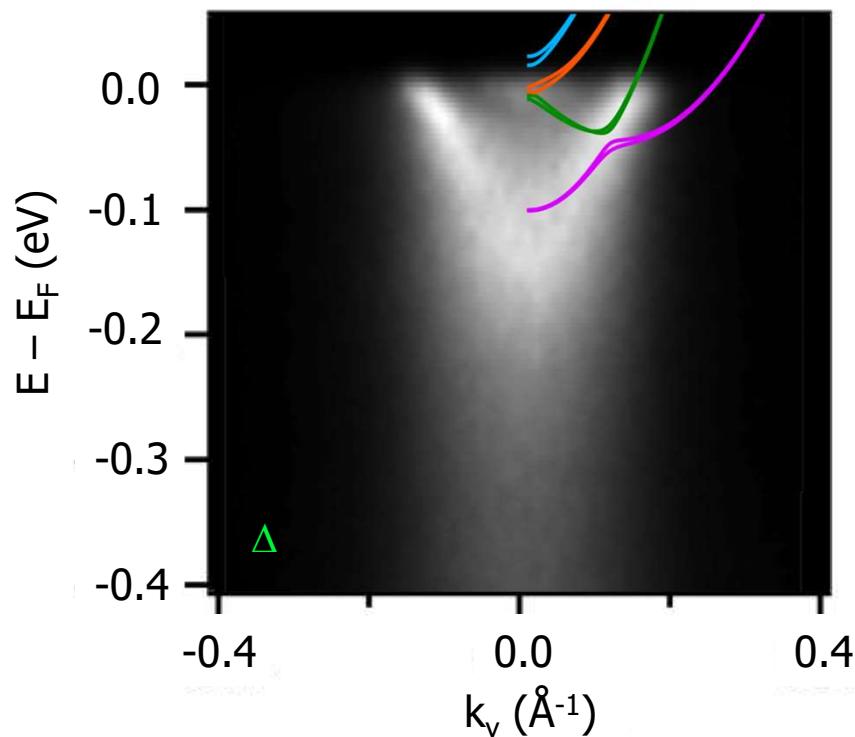
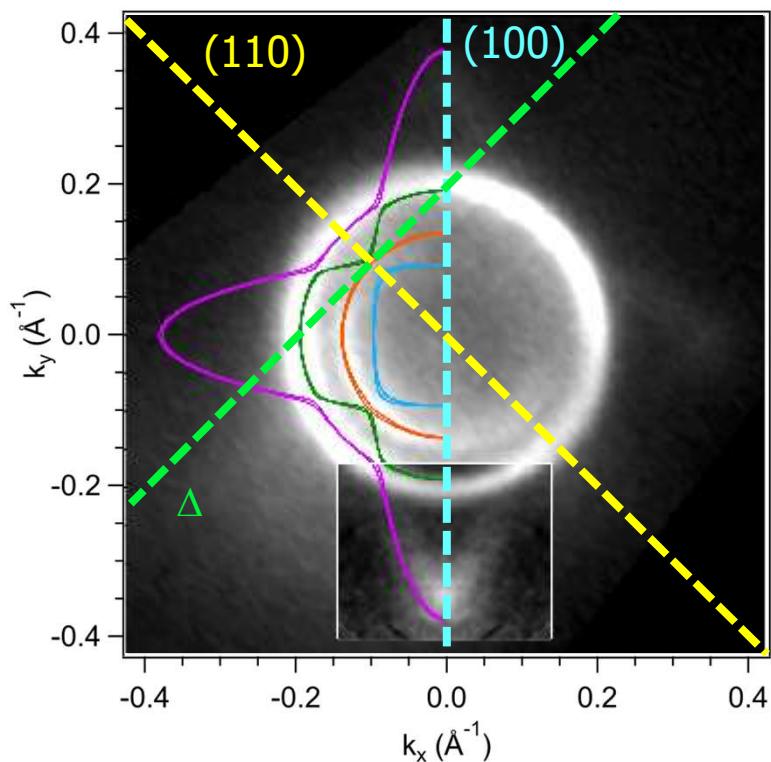
Electronic structure of STO 2DEGs



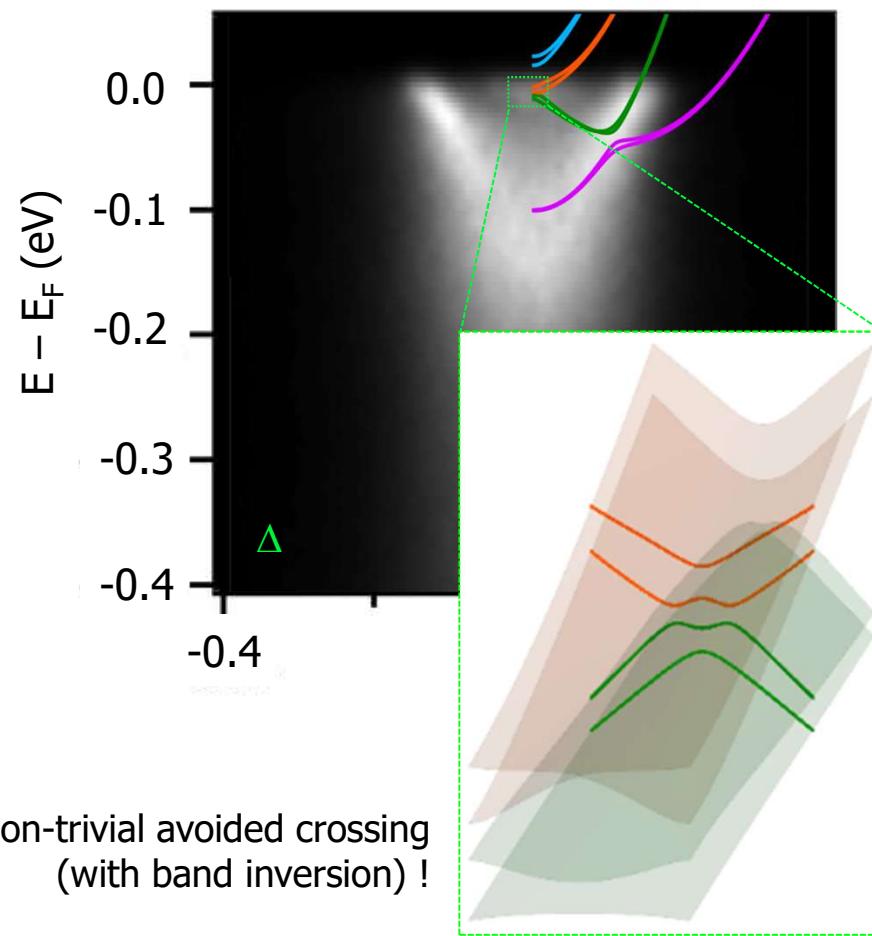
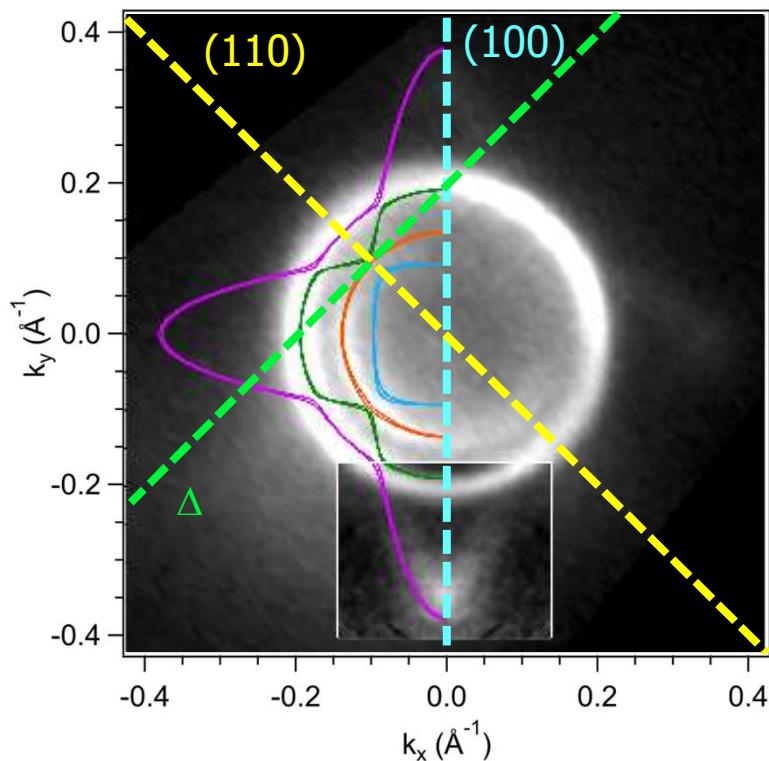
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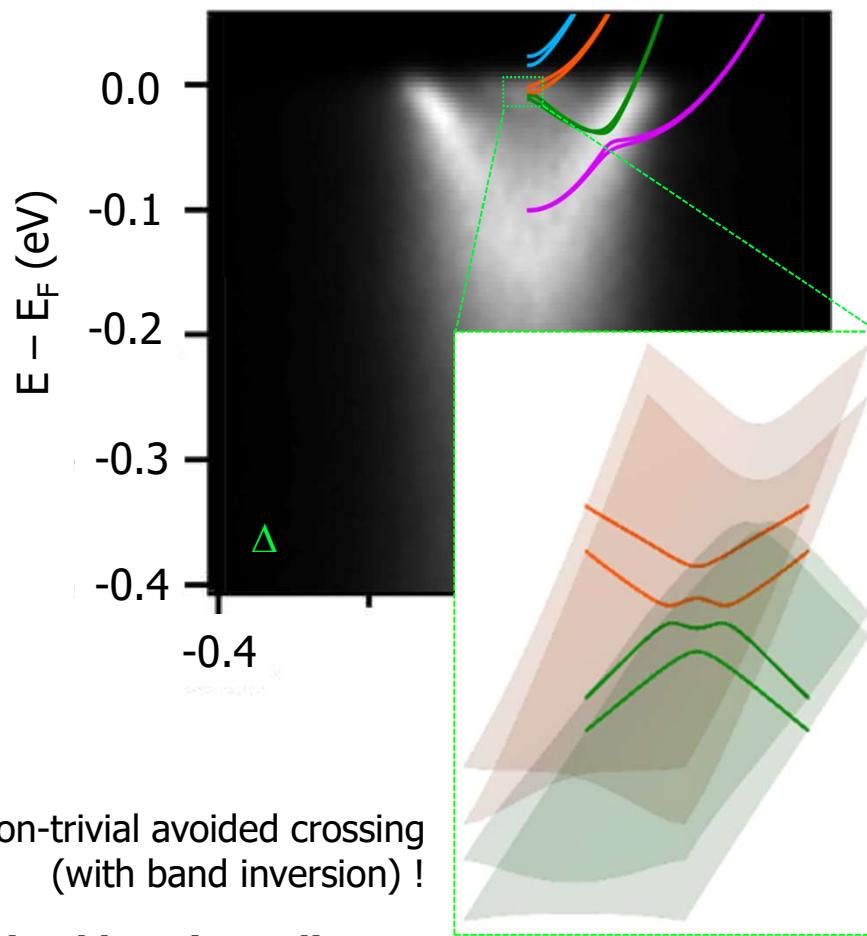
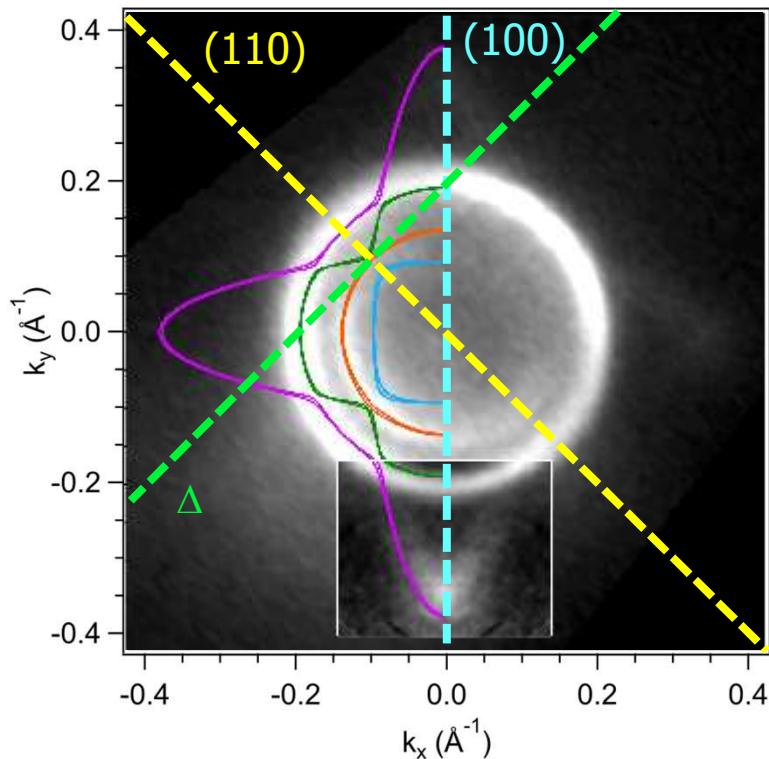
Electronic structure of STO 2DEGs



Topologically non-trivial avoided crossing
(with band inversion) !

M. Vivek et al, PRB 95, 165117 (2017)

Electronic structure of STO 2DEGs



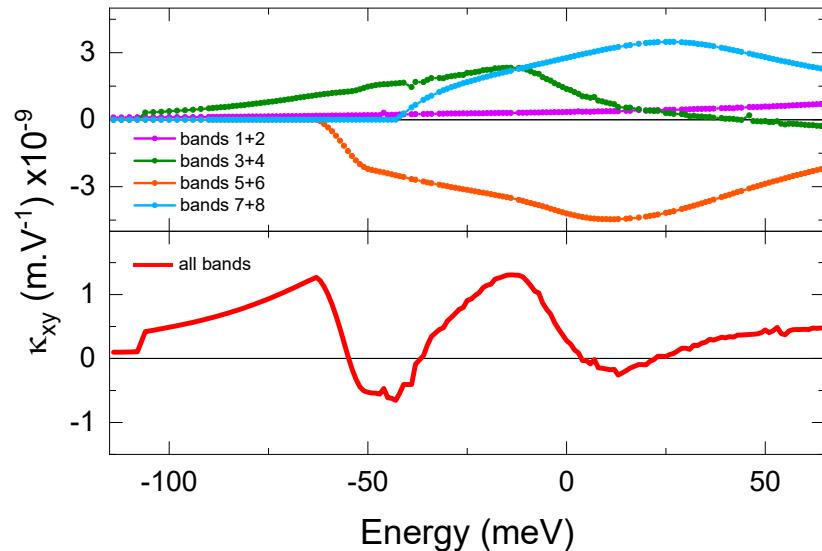
Topologically non-trivial avoided crossing
(with band inversion) !

Electronic structure of STO 2DEGs determined by 4 ingredients

1. Splitting of d_{xy} and $d_{xz/yz}$ bands by confinement potential
2. Sub-bands due to quantum confinement
3. Spin-orbit coupling
4. Orbital mixing

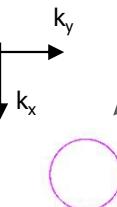
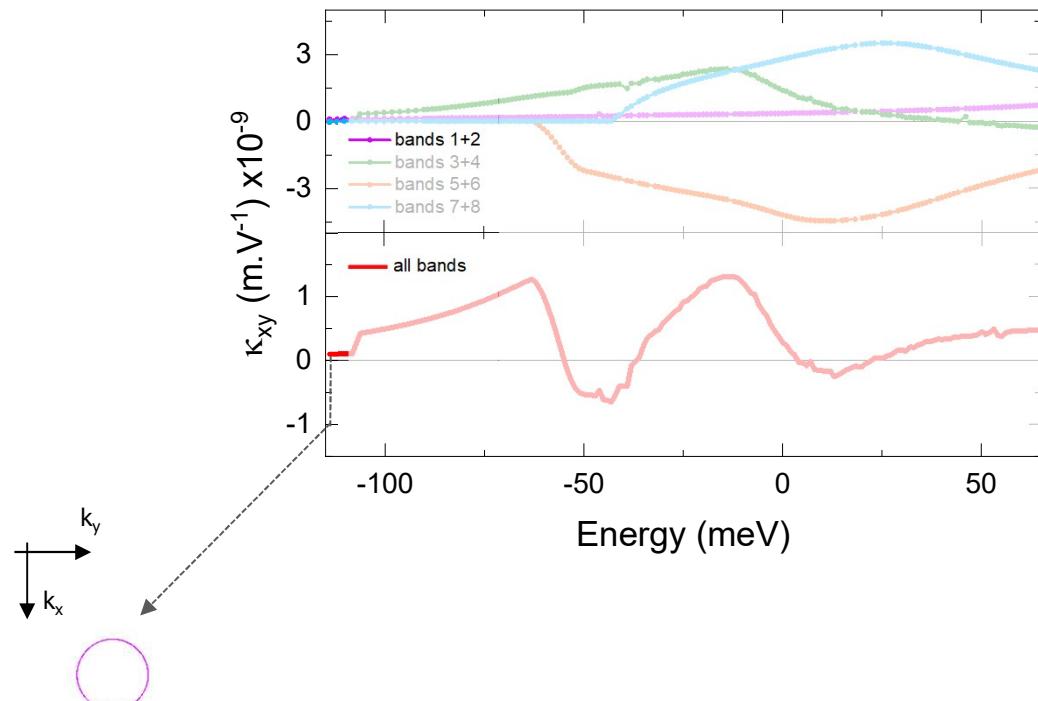
Vaz, MB et al, Nature Mater. 18, 1187 (2019)

Energy dependence of the Edelstein effect



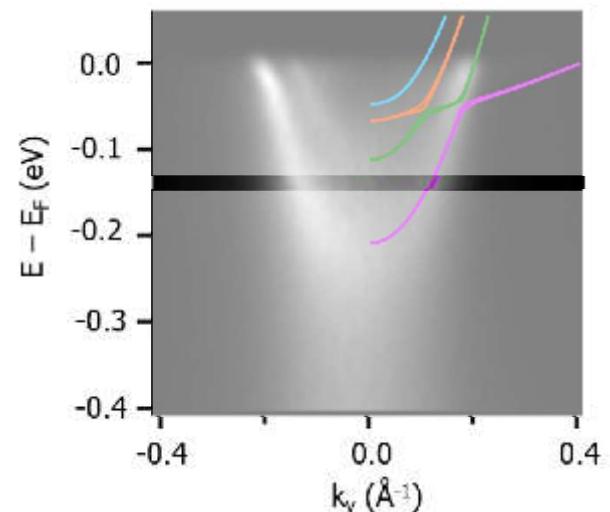
D.C. Vaz, MB et al, Nature Mater. 18, 1187 (2019)

Energy dependence of the Edelstein effect



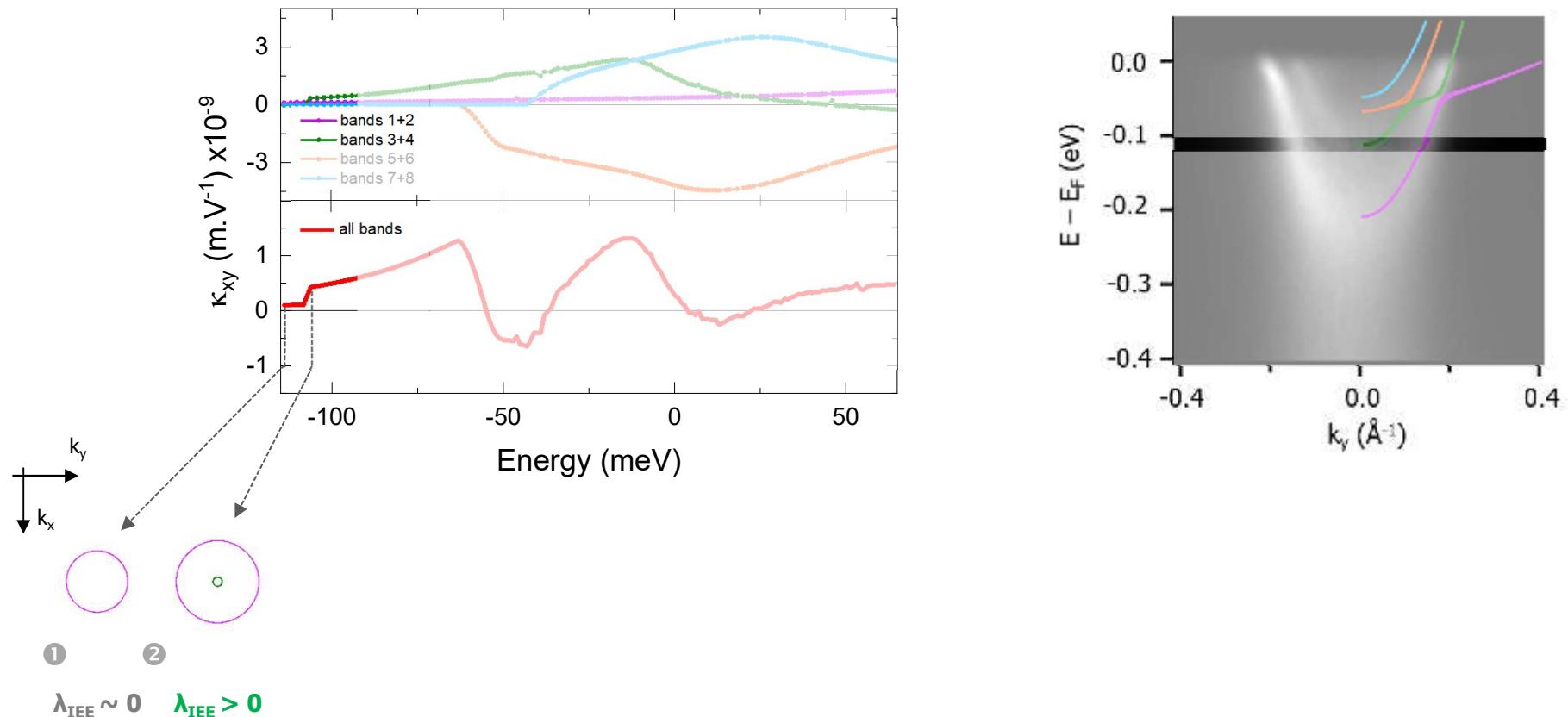
①

$\lambda_{\text{IEE}} \sim 0$



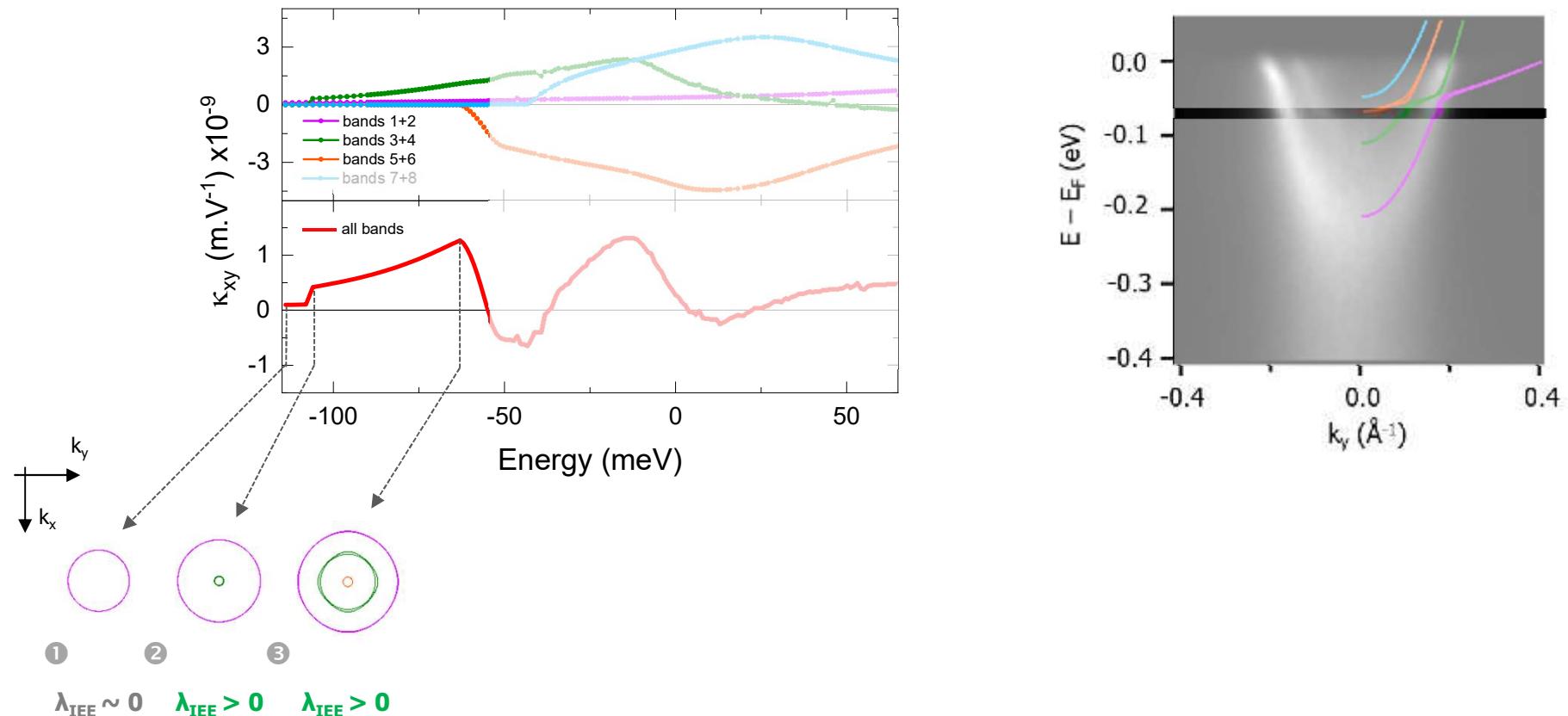
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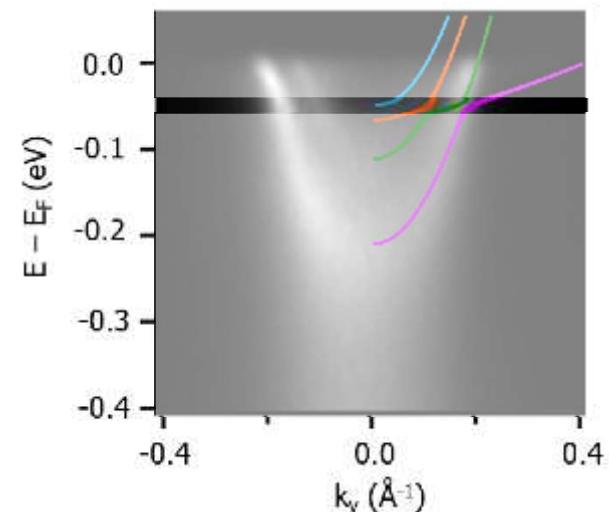
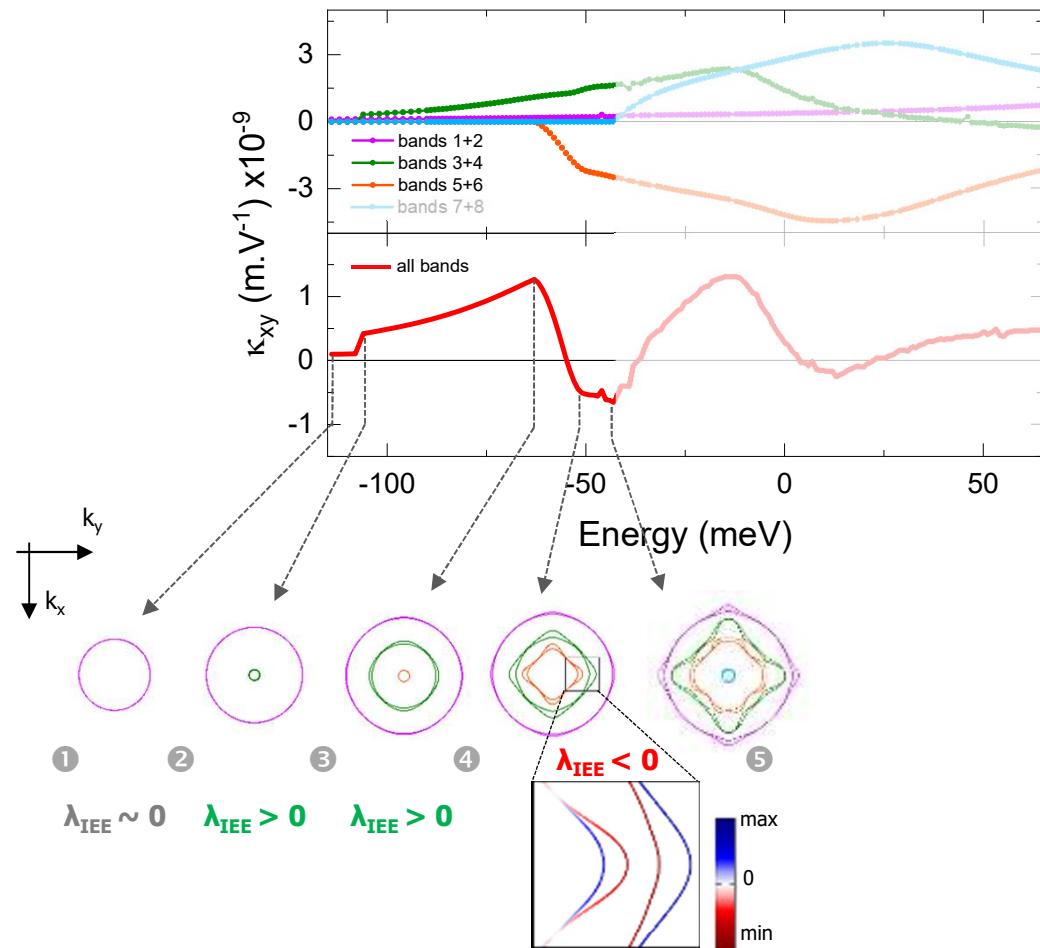
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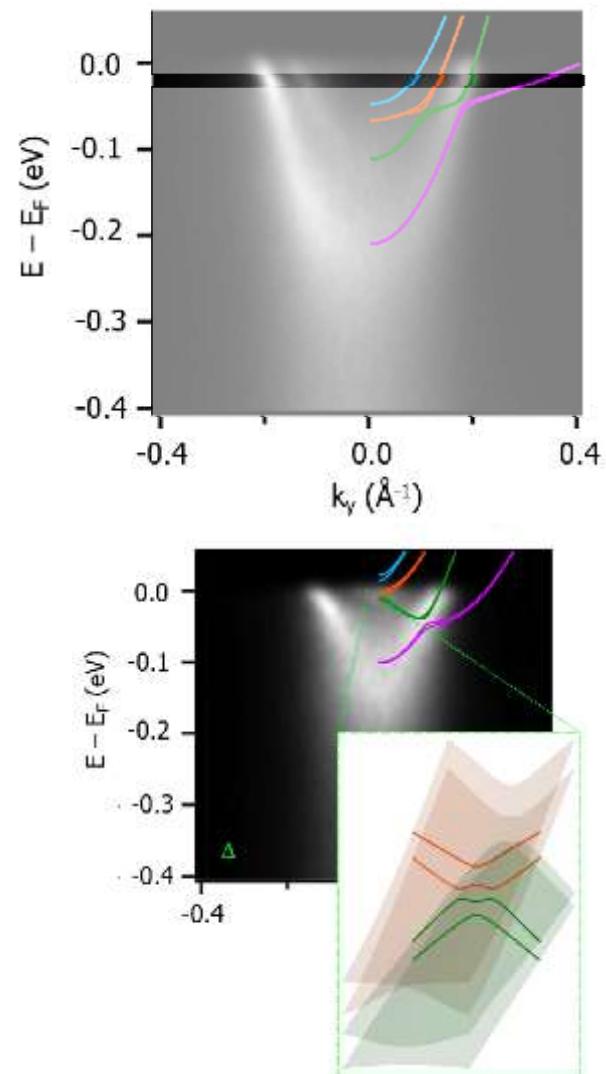
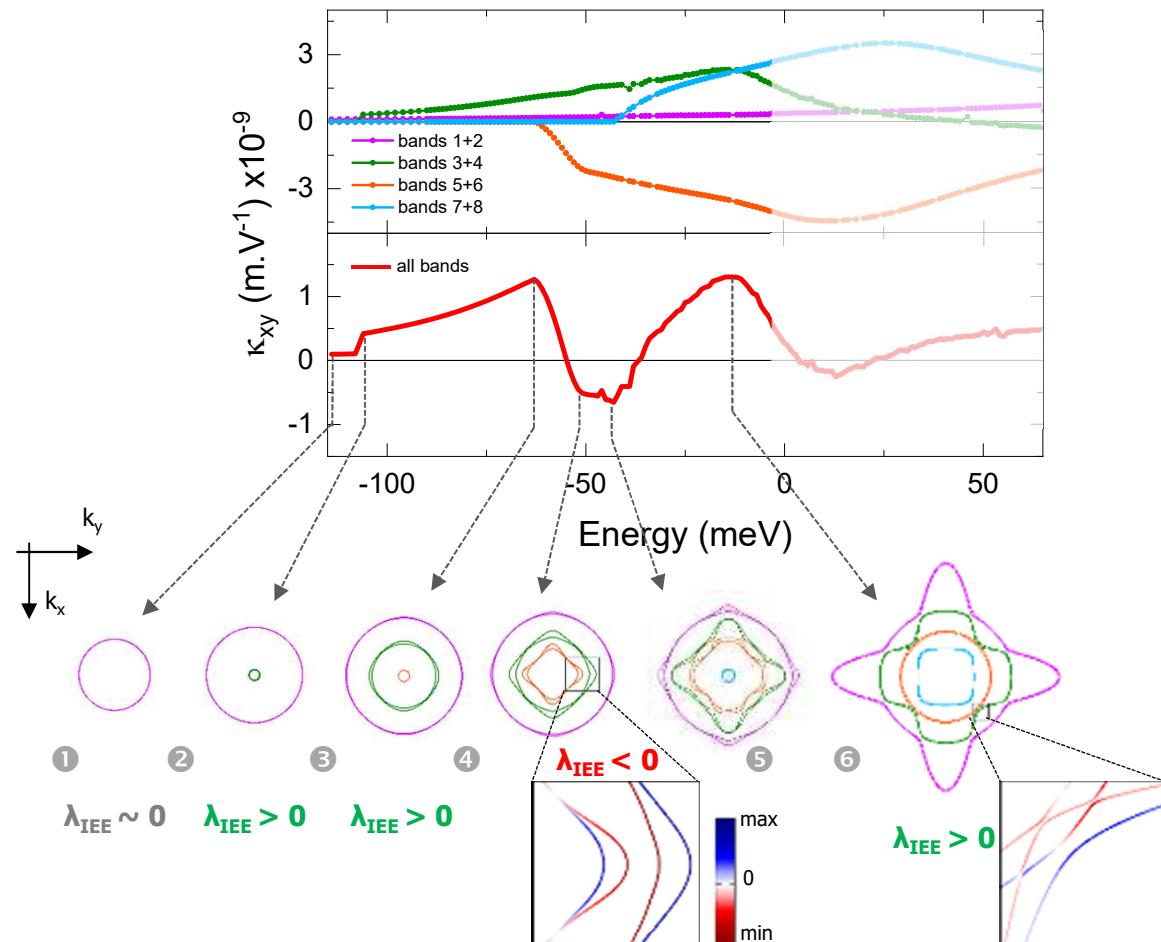
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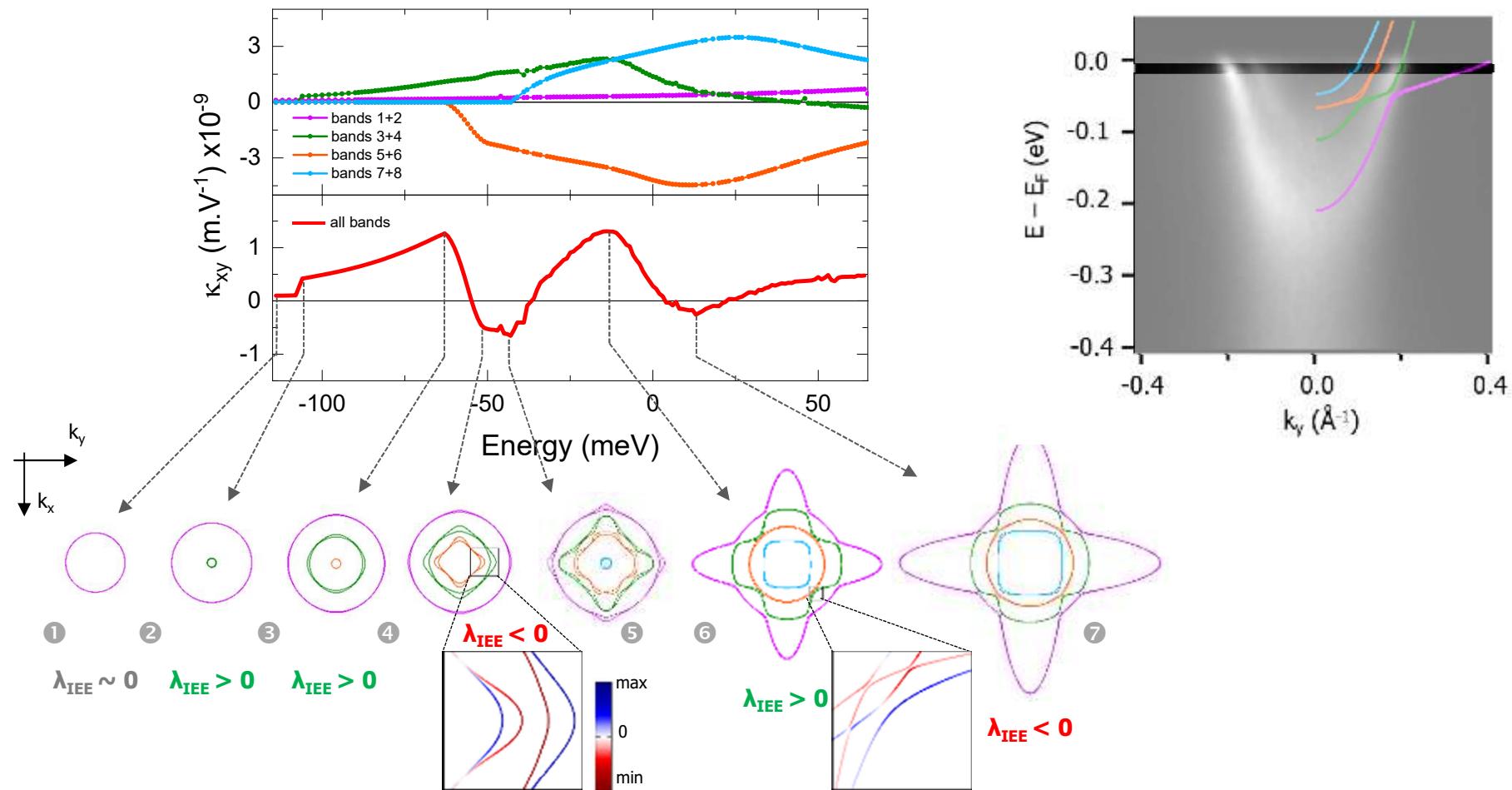
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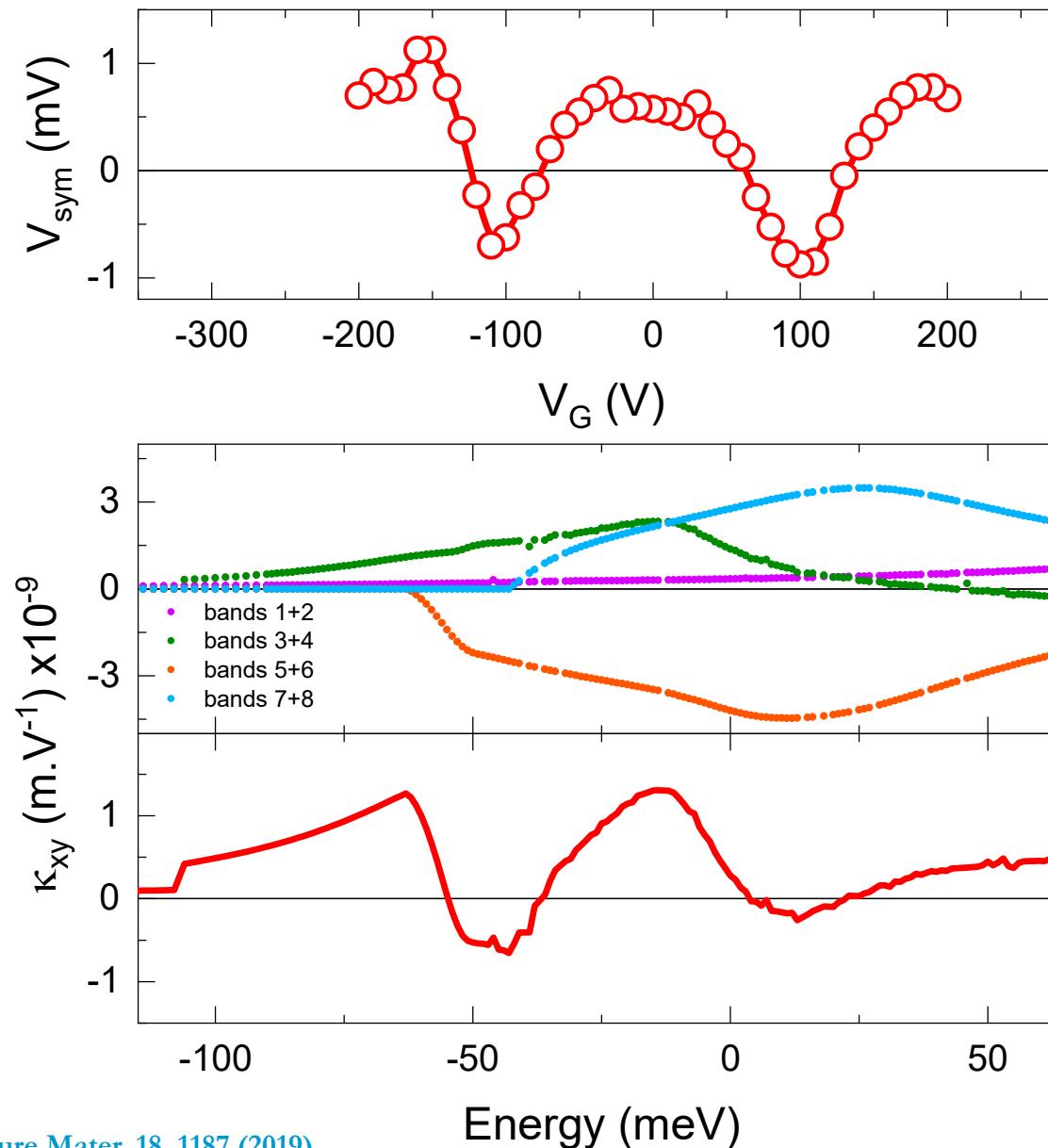
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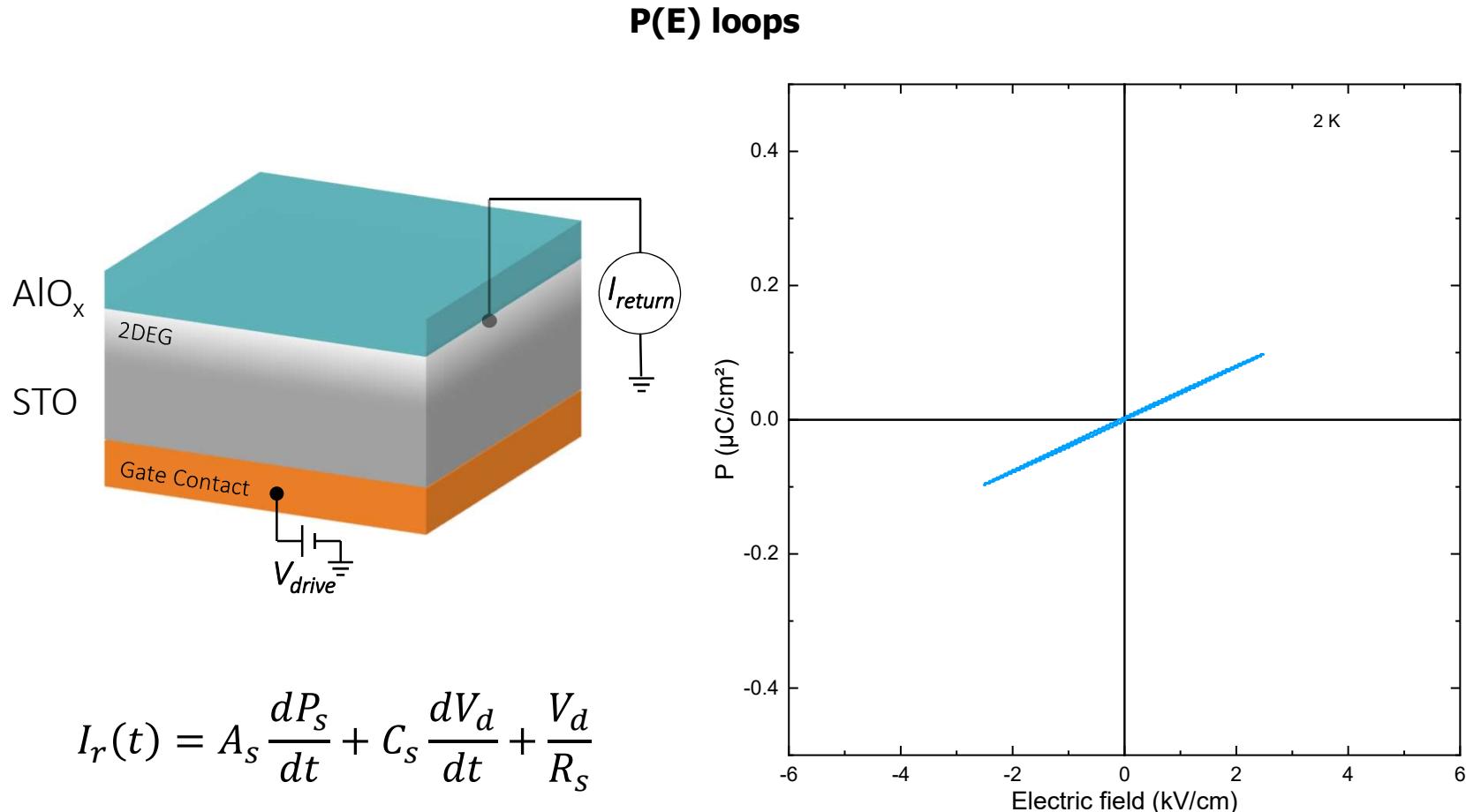
D.C. Vaz, MB et al, Nature Mater. 18, 1187 (2019)

Comparison with experiments



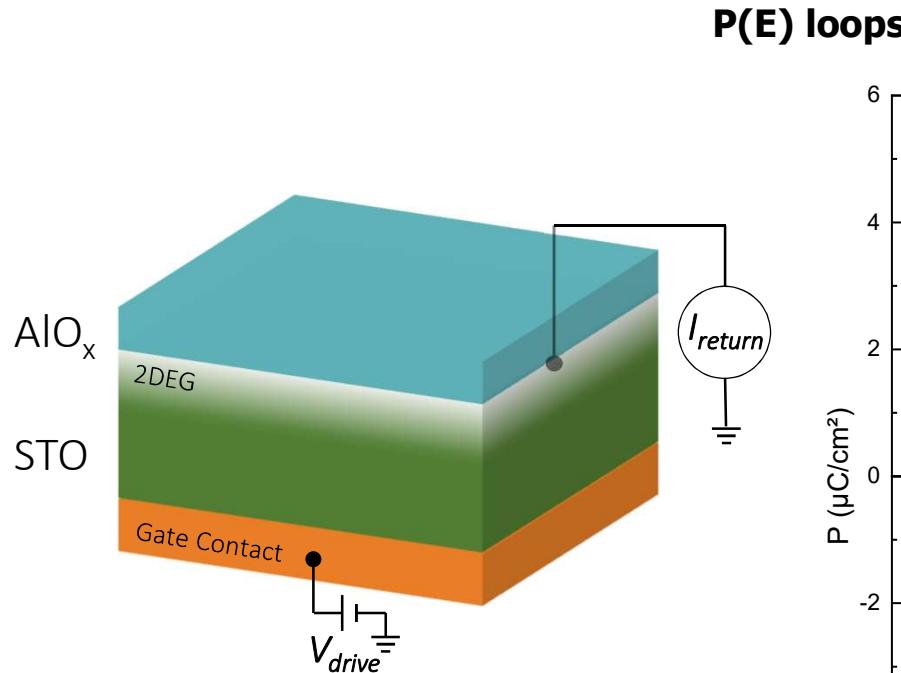
D.C. Vaz, MB et al, Nature Mater. 18, 1187 (2019)

E-field induced ferroelectricity in SrTiO₃

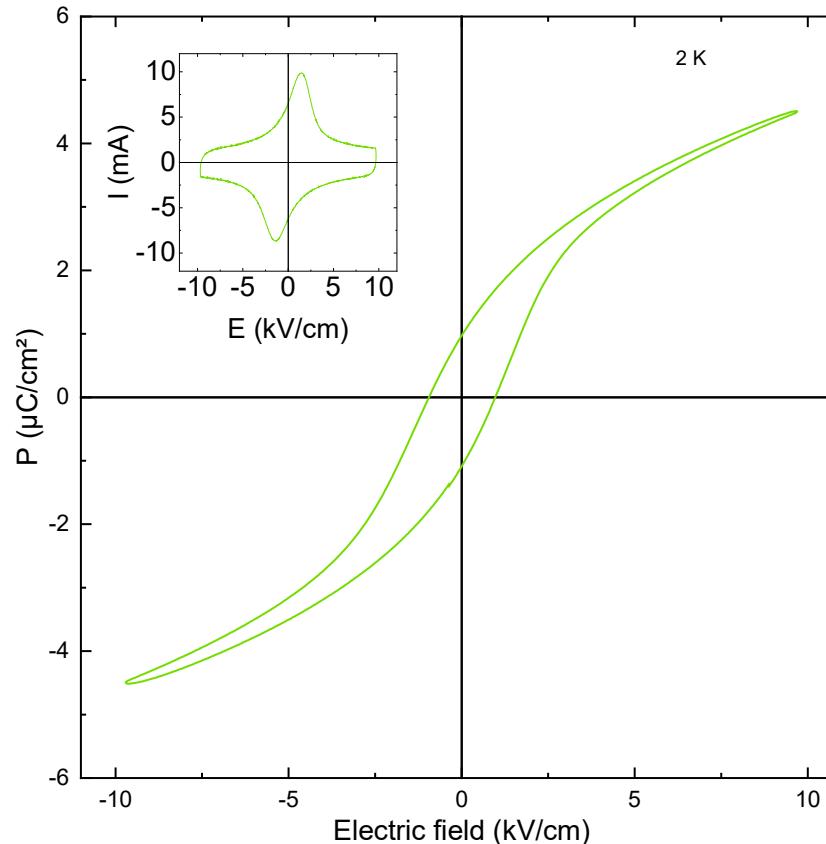


- After 2DEG formation by Al deposition, STO is thinned down to 200 μm , allowing to apply large electric fields

E-field induced ferroelectricity in SrTiO₃

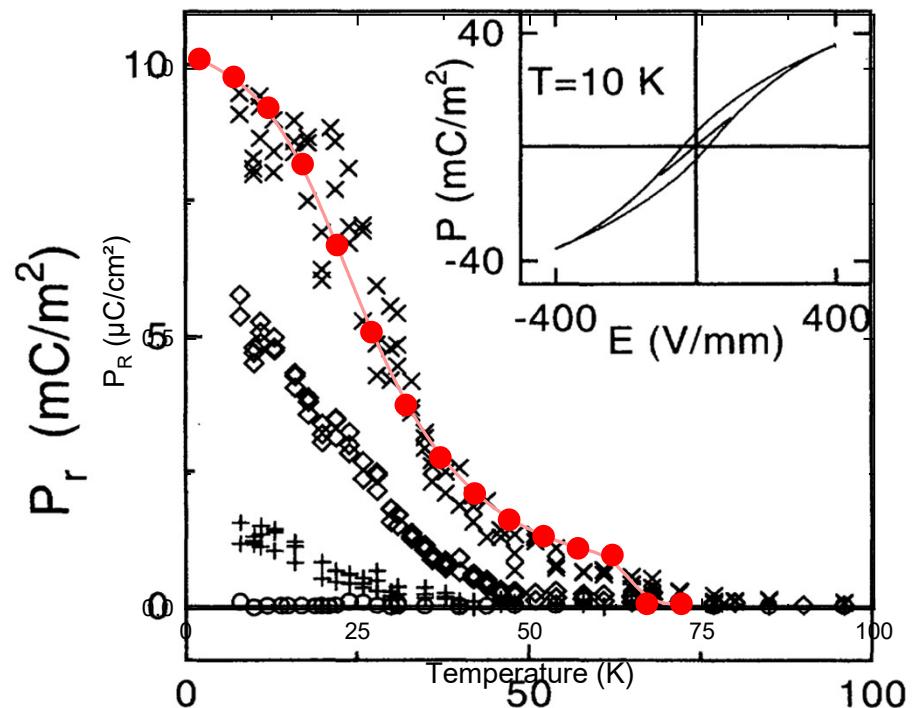
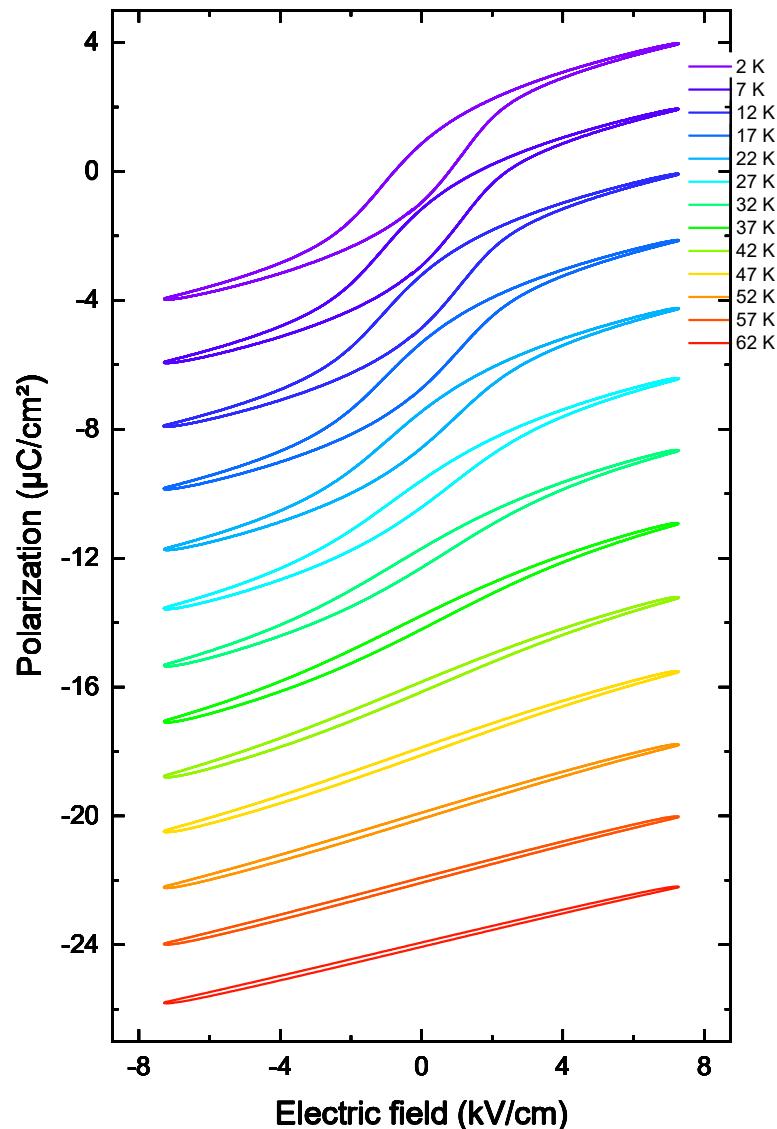


$$I_r(t) = A_s \frac{dP_s}{dt} + C_s \frac{dV_d}{dt} + \frac{V_d}{R_s}$$



- After 2DEG formation by Al deposition, STO is thinned down to 200 μm , allowing to apply large electric fields
- Open hysteresis loops appear beyond a threshold E field, indicative of a **ferroelectric-like state**

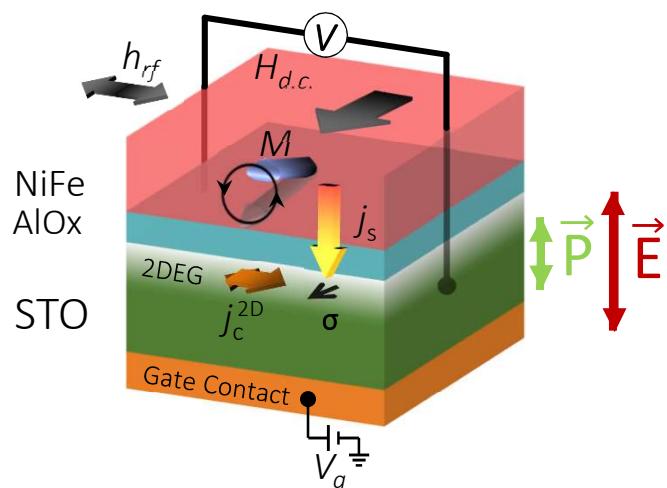
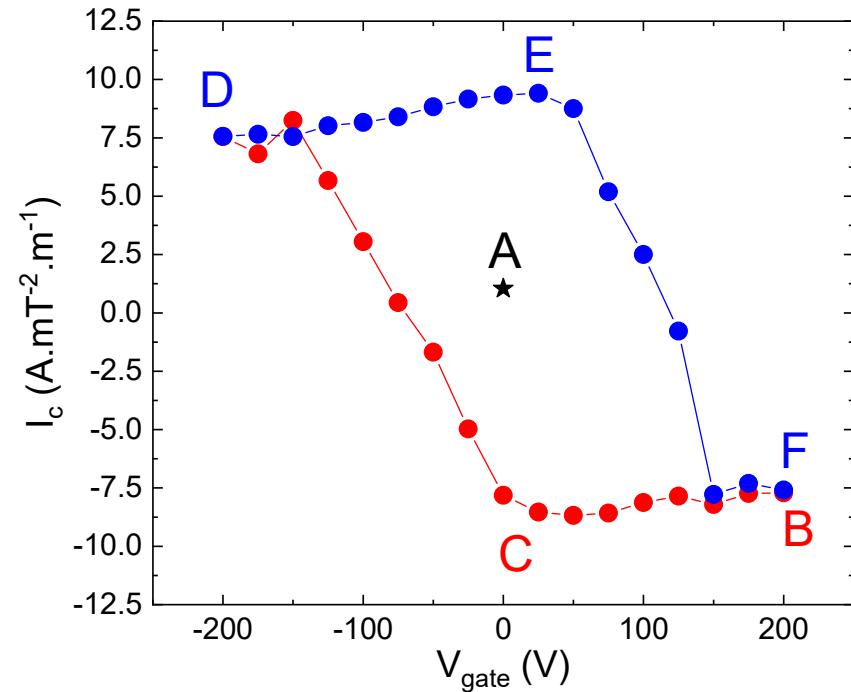
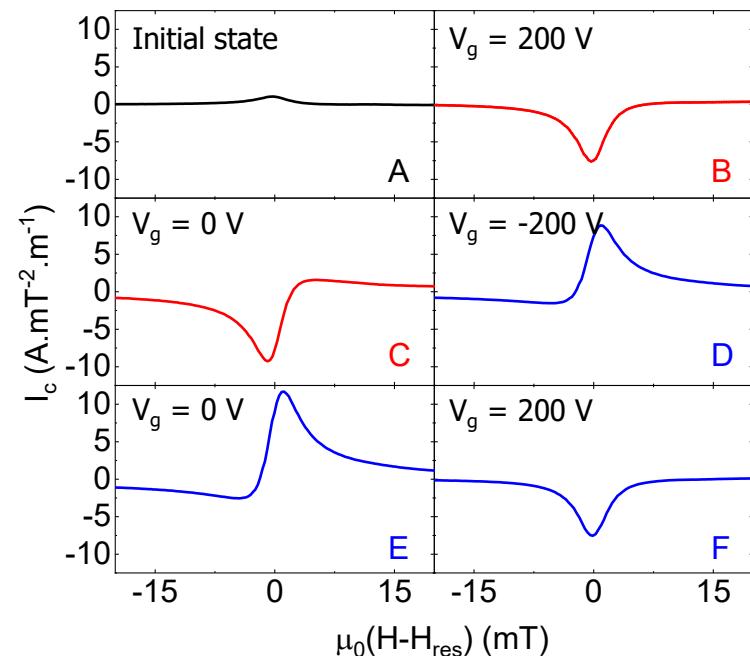
E-field induced ferroelectricity in SrTiO₃



⌚ Ferroelectric state disappears around 50 K

H. Gränicher, Helvetica Phys. Acta 29, 210 (1956) ;
J. Hemberger et al PRB 52, 13159 (1995) ;
J. Sidoruk, J. et al. JPCM 22, 235903 (2010) ;
H. Manaka, JPSJ. 86, 114702 (2017).

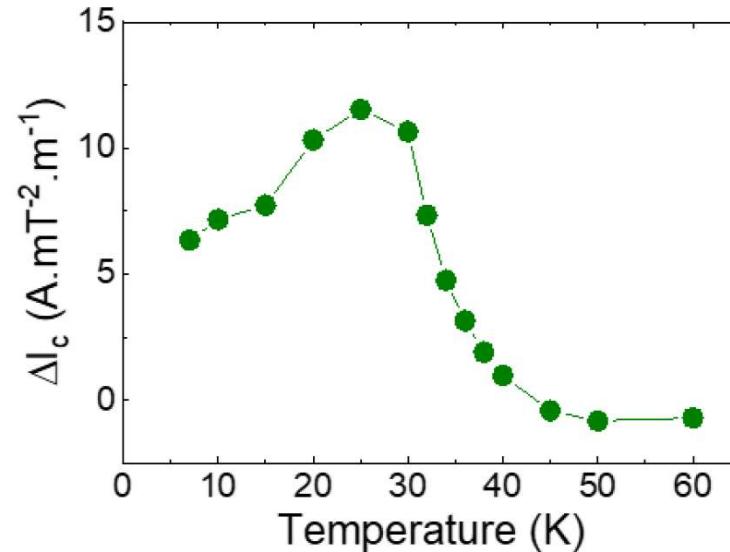
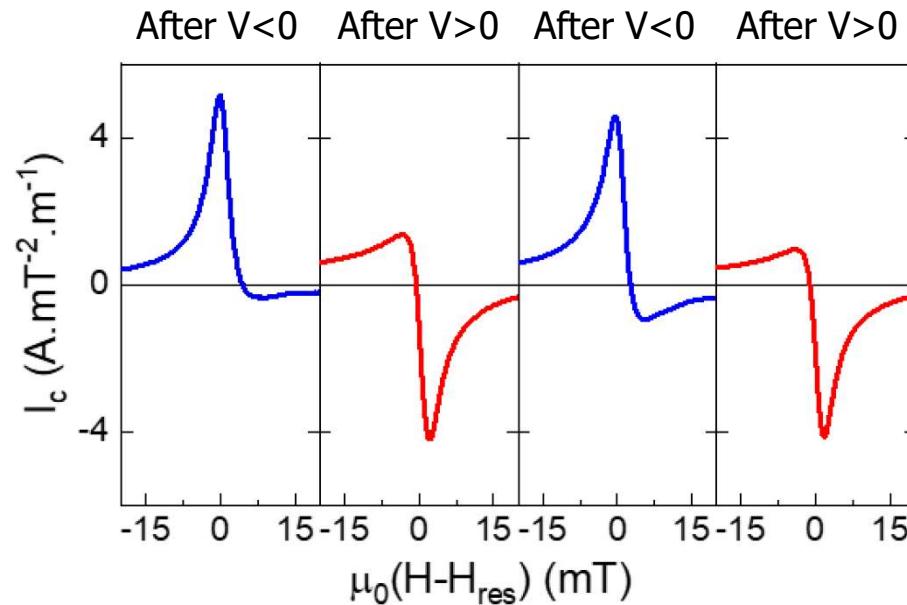
Spin-charge conversion in ferroelectric Al/STO 2DEGs



- Charge current production by the inverse Edelstein effect shows a hysteretic dependence with the electric field, with two remanent states corresponding to positive and negative charge currents

P. Noël, MB et al., Nature 580, 483 (2020)

Spin-charge conversion in ferroelectric Al/STO 2DEGs

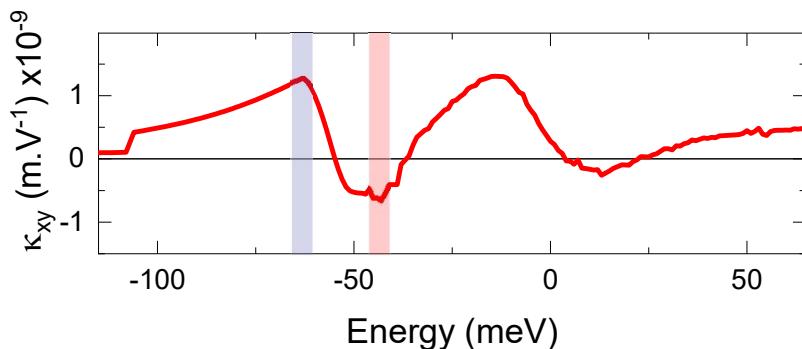


P. Noël, MB et al., Nature 580, 483 (2020)

- Effect is reversible and disappears around 40-50 K just as the remanent polarization in the P-E loops
- Strong indication that the electrical switching of spin-charge conversion is driven by ferroelectricity
- Control of spin-charge conversion by ferroelectricity**
- Huge potential for low-power logic-in-memory architectures beyond CMOS**

Possible mechanism ?

- Modulation of carrier density and E_F by polarization direction
- P_{up} : higher $E_F \rightarrow$ negative conversion
- P_{dn} : lower $E_F \rightarrow$ positive conversion

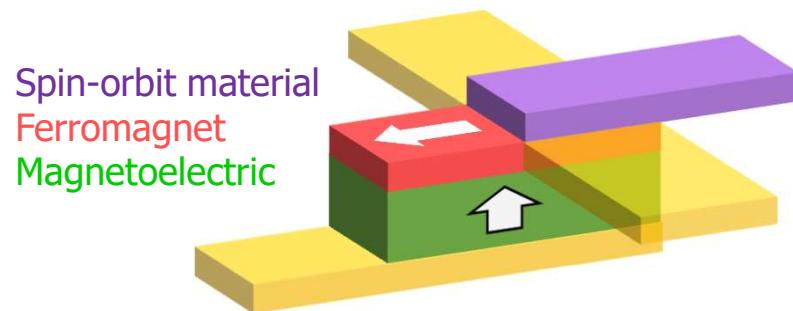


- Ferroic orders, multiferroics and magnetoelectric coupling
- Approaches for the electric-field control of magnetism
- Electric-field control of spin-charge interconversion
- ○ Low-power spin-based devices for logic-in-memory

Magneto-electric spin-orbit transistor (MESO)

Original design by Intel

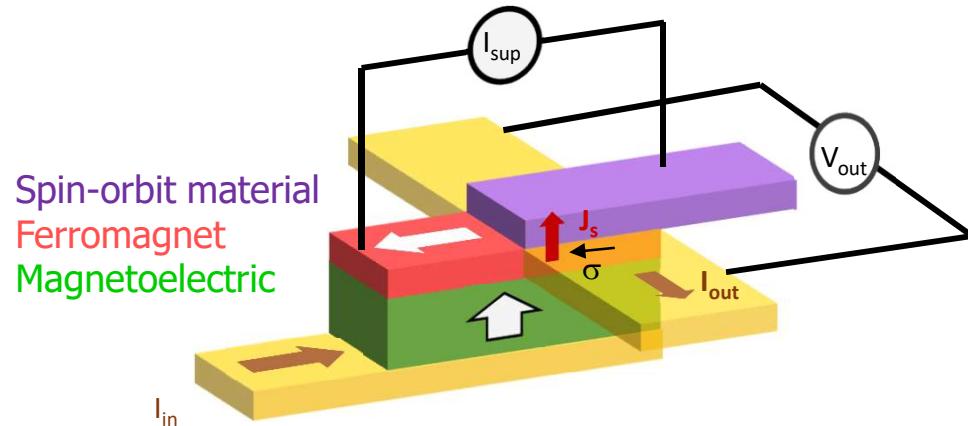
Manipatruni et al, Nature 565, 25 (2019)



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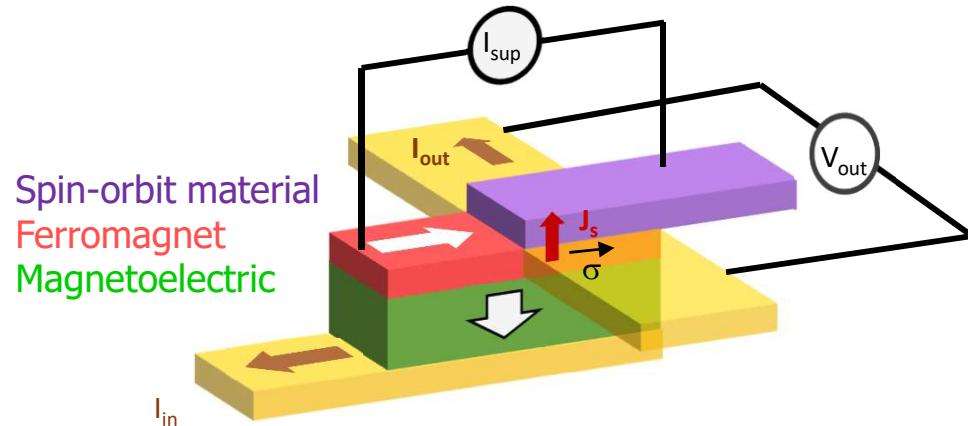
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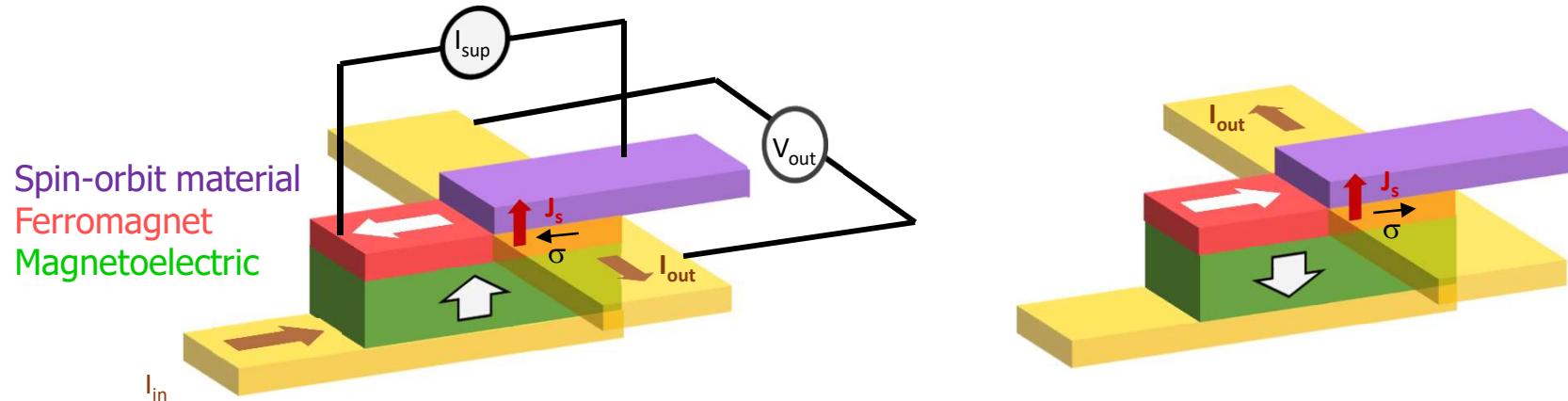
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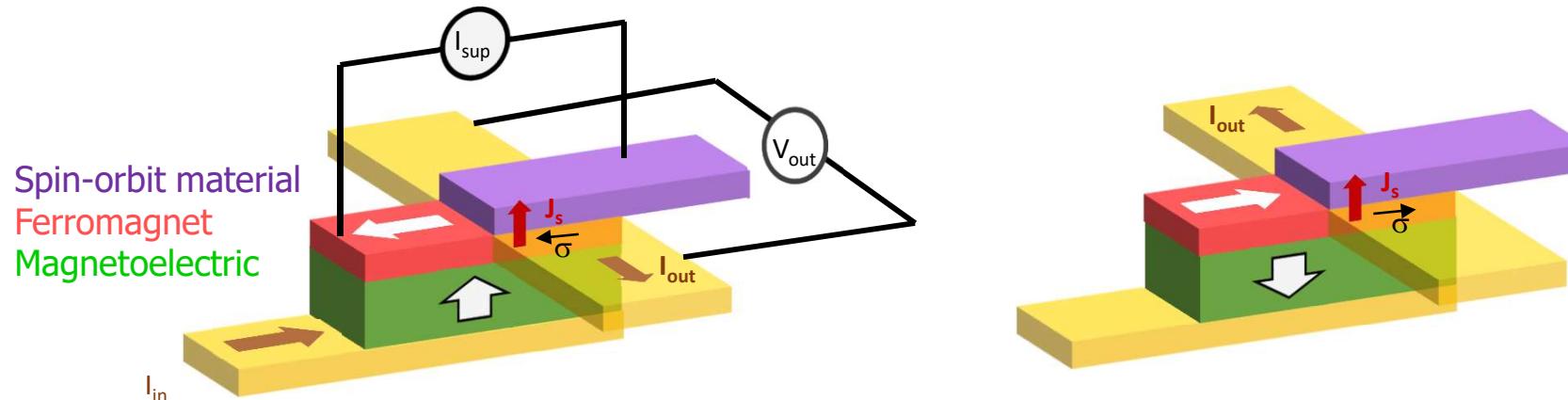


- New type of non-volatile spin-based transistor proposed by **Intel**
- Operates through **magnetoelectric coupling** (input) and **spin-orbit coupling** (output)
- Memory and logic embedded
- Scalable, concatenable, implementable as majority gate
- Low power (30 times less than CMOS for same size)
- 100 mV operating voltage (ME needs to switch with <100 mV ; SOC must generate >100 mV)

Magneto-electric spin-orbit transistor (MESO)

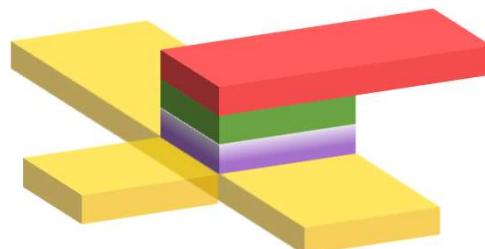
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Alternative design CNRS/Thales/Spintec

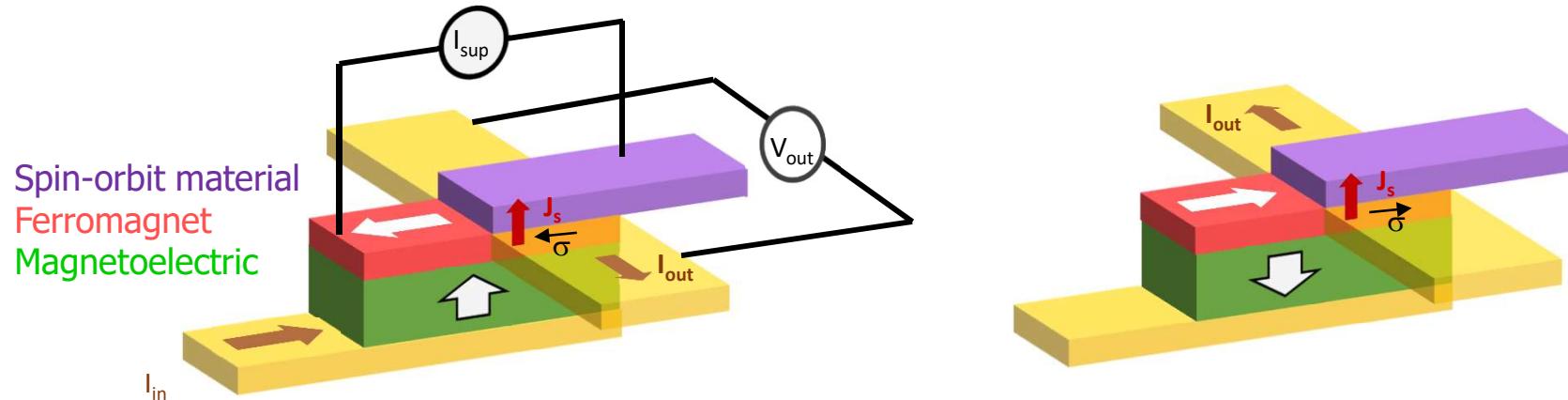
MB et al, patent FR18 74319 ; Nature 580, 483 (2020)



Magneto-electric spin-orbit transistor (MESO)

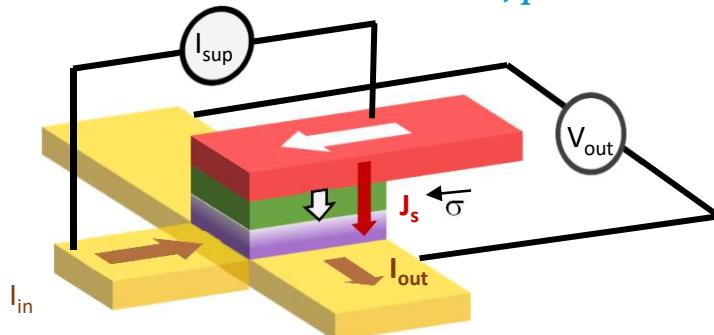
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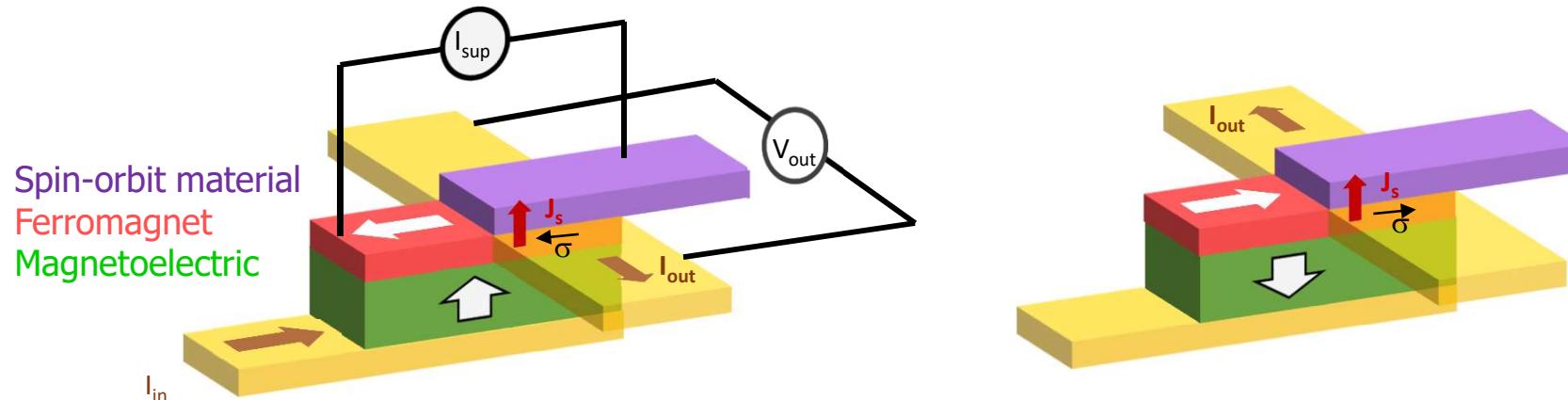
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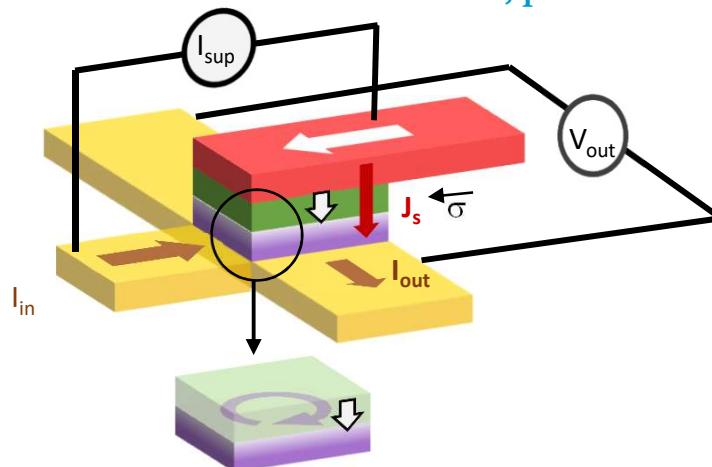
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Manipatruni et al, Nature 565, 25 (2019)



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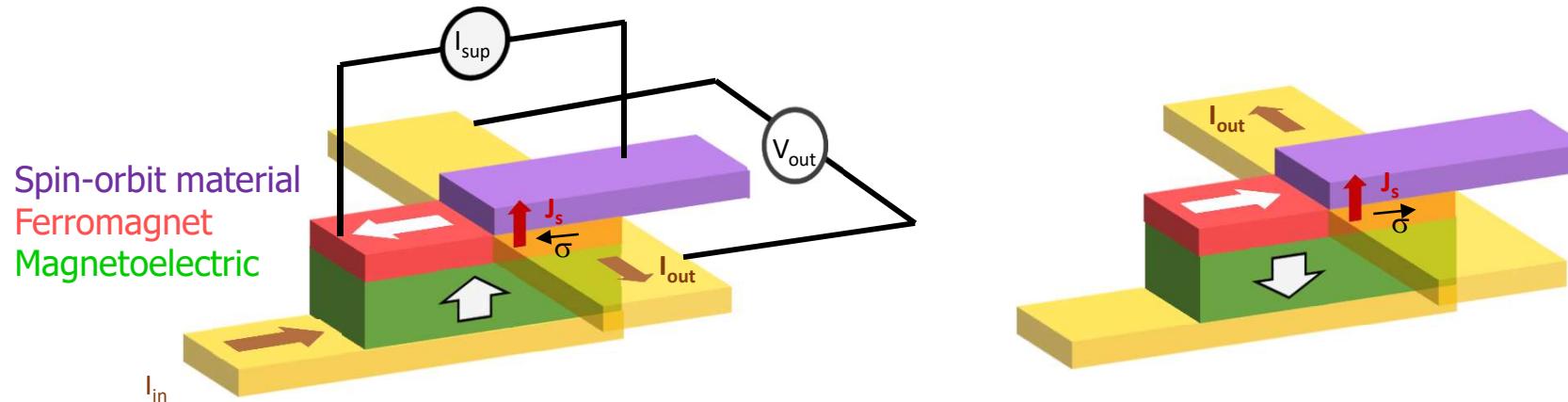
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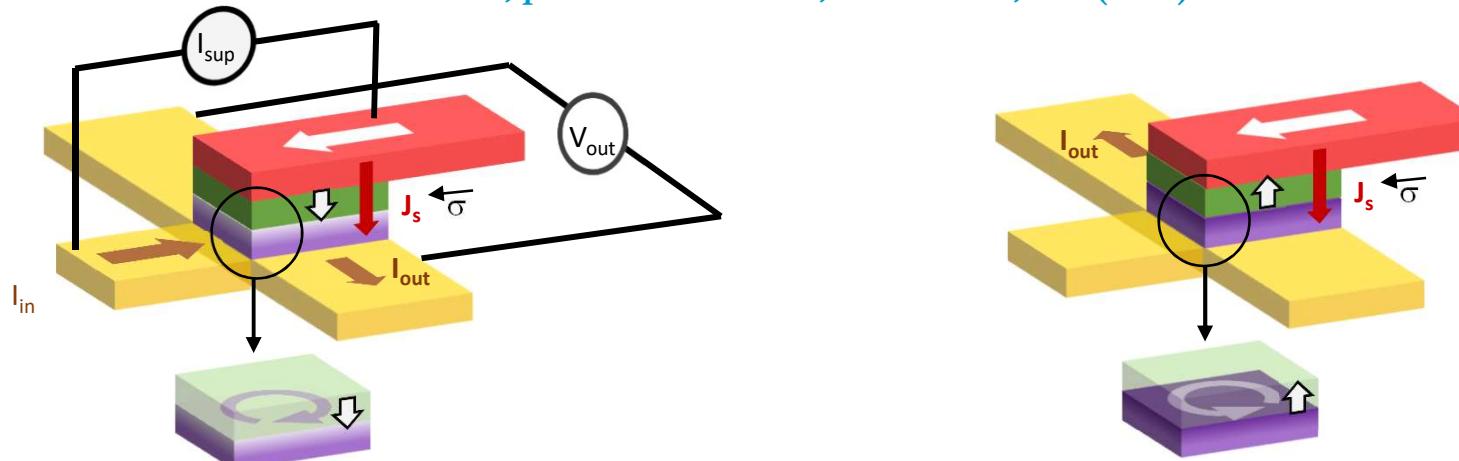
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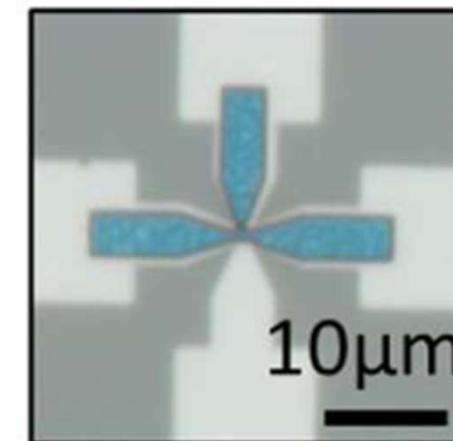
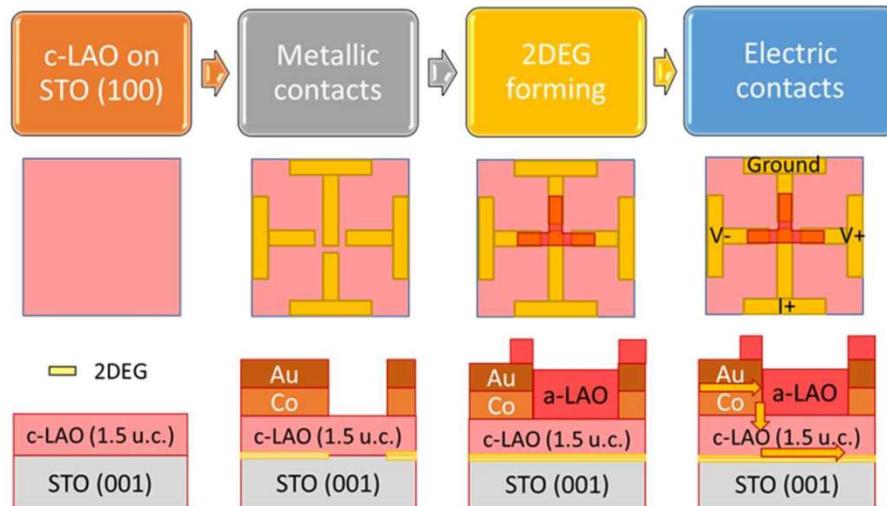
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MB et al, patent FR18 74319 ; Nature 580, 483 (2020)

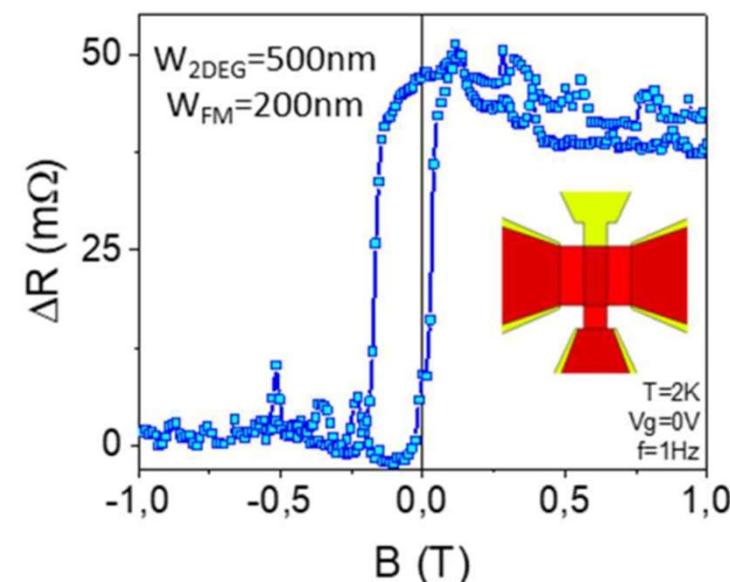
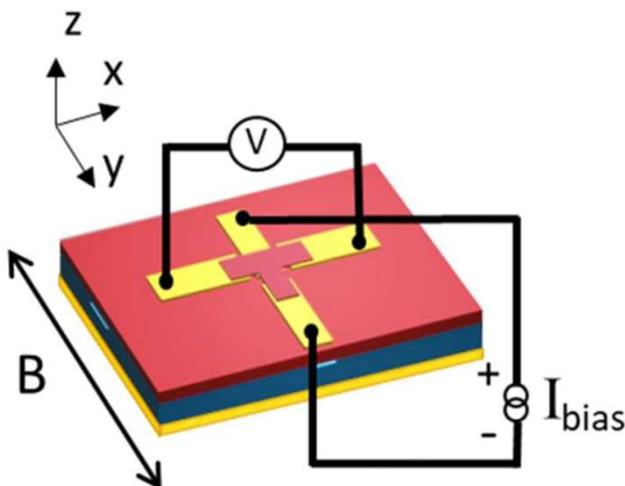


- ➊ Magnetoelectric replaced by non-magnetic ferroelectric
- ➋ No need for magnetization switching
- ➌ Ferroelectricity controls Rashba state at interface with spin-orbit material

All electrical detection of spin-charge conversion in STO nanodevices



Patterning of T-shape devices by e-beam lithography

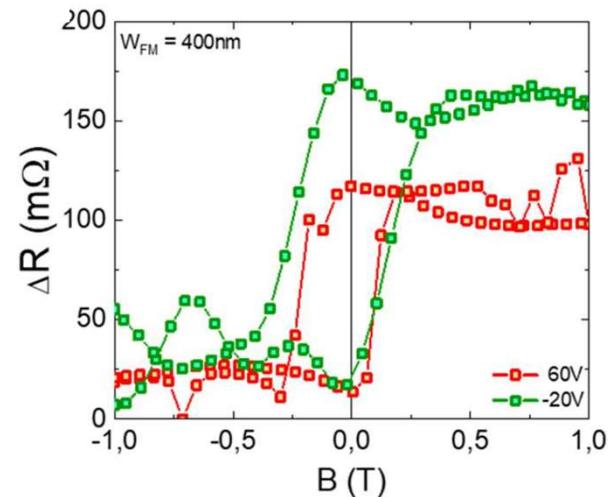


Direct evidence for spin-charge conversion from magnetotransport in STO 2DEGs

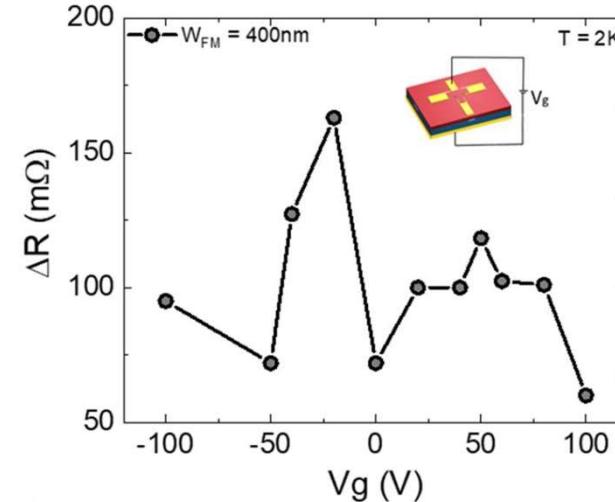
F. Gallego, MB et al, Adv. Func. Mater. 34, 2307474 (2024)

All electrical detection of spin-charge conversion in STO nanodevices

Gate dependence

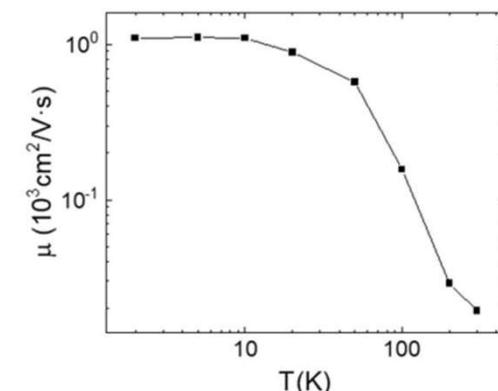
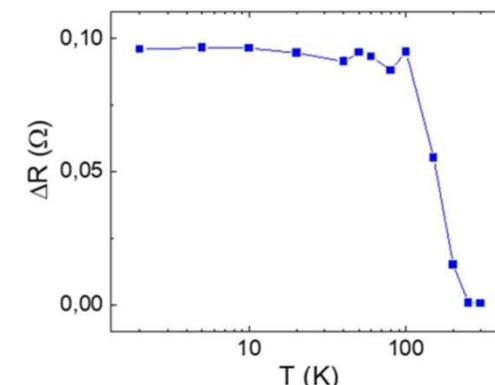
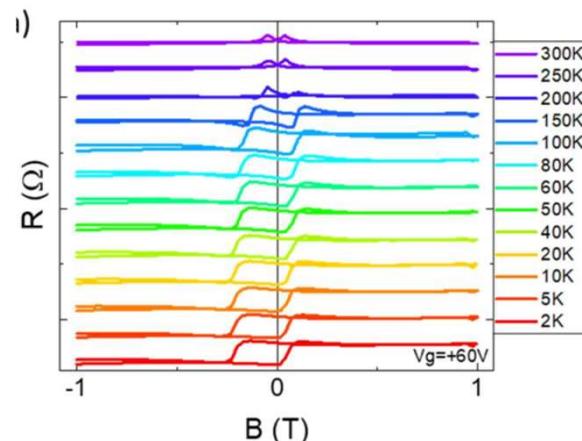


Signal amplitude depends on gate voltage



Non-monotonic gate dependence as in spin-pumping

Temperature dependence



Most devices show exchange bias, probably due to CoOx at Co/LAO interface

Signal disappears > 200 K, probably due to decrease in conversion efficiency / scattering time

Conclusions and perspectives

There are many possibilities for E-field control of magnetism

- With multiferroics (works but at room temperature limited to BiFeO₃)
- Through field effect (limited efficiency, requires ultrathin films and large fields)
 - ❖ Find better materials combinations between FE and FM → ferroelectric nitrides (e.g. AlScN)?
- Through field-controlled strain (efficient but hard to miniaturize + poor endurance)
- New possibilities for spin/charge interconversion at oxide interfaces with Rashba SO coupling
 - ❖ FE control demonstrated
 - ❖ Need reliable, low leakage demonstration at room temperature
 - ❖ Implementation in FESO-type nanoscale devices in under way

<https://nellow.eu/>



Thank you !

