

Spin Currents and Spin Caloritronics

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Barcelona, Spain

Cargèse
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Image credit: NPG

Spin currents and spin caloritronics

Outline

Spin caloritronics: *Thermoelectricity meets spintronics*

- Nonequilibrium phenomena involving charge, energy and spin transport. Usually in magnetic structures

Area of research	Topic
Electronics	Transport/manipulation of charge
Spin electronics or Spintronics	Transport/manipulation of charge and spin
Calorimetry	Study and measure the heat of chemical reactions or physical changes
Spin Caloritronics	Transport/manipulation of charge, spin and heat

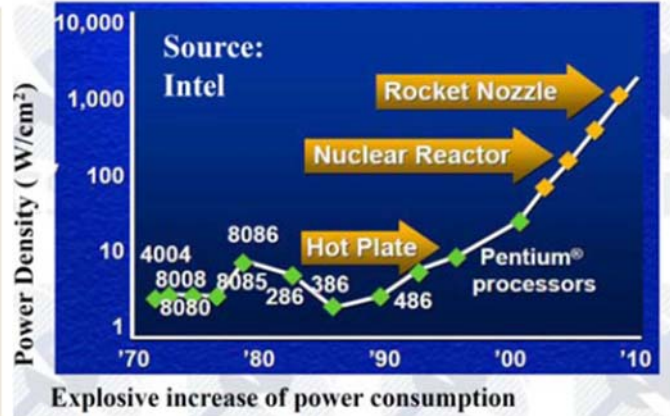
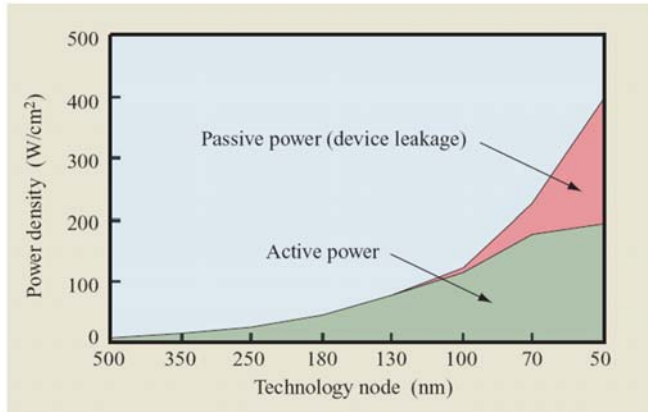
M. Johnson and R. H. Silsbee, Phys. Rev. B **35**, 4959 (1987)

Reviews: Bauer, MacDonald, Maekawa, Solid State Commun. (2010); Bauer in *Spin Current* (Oxford University Press, 2012)

Spin currents and spin caloritronics

Outline

Spin caloritronics: *Why is important?*

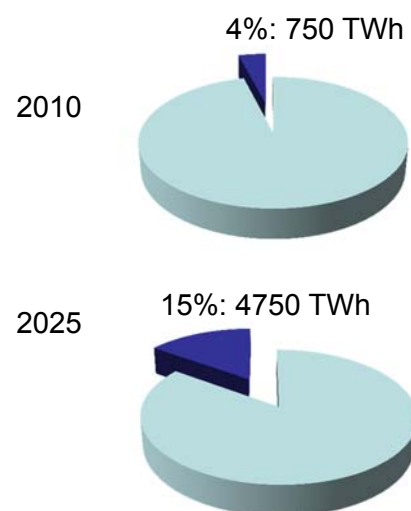
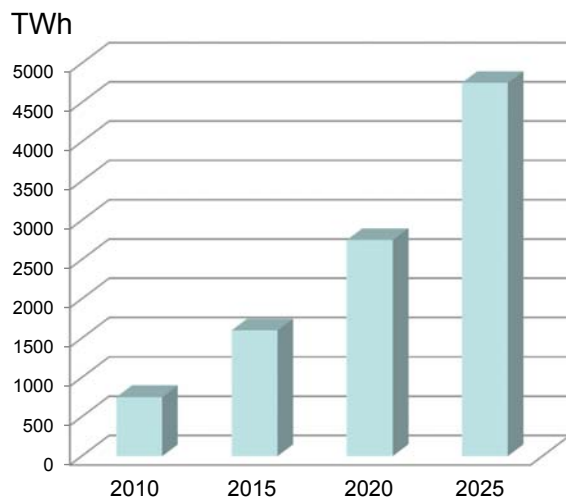


Reviews: Bauer, MacDonald, Maekawa, Solid State Commun. (2010); Bauer in *Spin Current* (Oxford University Press, 2012)

Spin currents and spin caloritronics

Outline

ICT Consumption Forecast

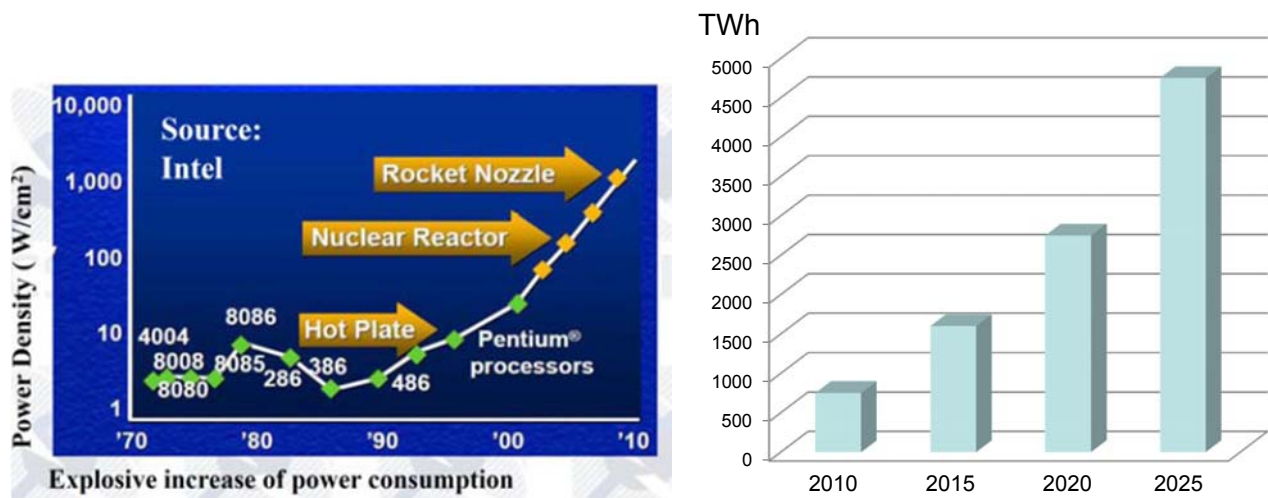


Goal: To find creative ways of using wasted heat

Spin currents and spin caloritronics

Outline

Spin caloritronics: *Why is important?*

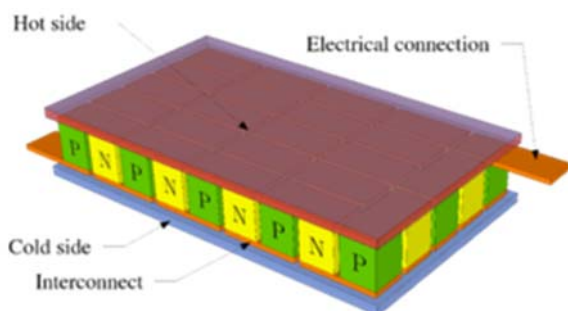


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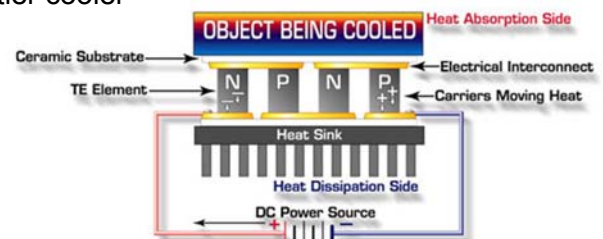
Spin currents and spin caloritronics

Outline

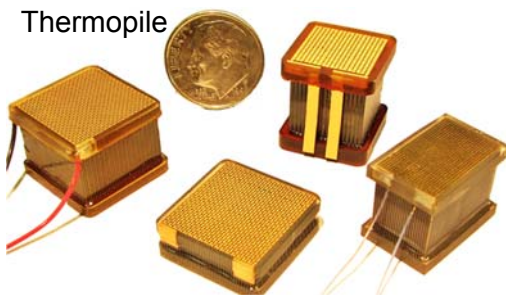
Thermoelectricity, energy harvesting, cooling



Peltier cooler



Thermopile



Marlow Industries



Spin currents and spin caloritronics

Outline

Spin caloritronics

- Thermoelectricity
 - Seebeck effect
 - Peltier effect
 - Onsager reciprocity relation
- Spin Currents: independent electron, collective
- Thermoelectricity + Spins
- Observed spin caloritronic effects: spin-dependent Seebeck effect, spin Peltier effect, thermal GMR, thermal GMR, thermal torques, magnon drag, spin Seebeck effect.

Spin currents and spin caloritronics

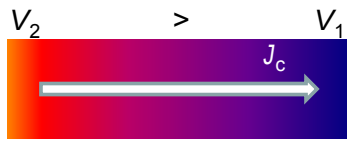
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Thermoelectricity

Electrical vs thermal transport



$$\sigma = \left(\frac{J_c}{\Delta V} \right)_{\Delta T=0} \quad \text{Electrical Conductivity}$$



$$\kappa = \left(\frac{J_Q}{\Delta T} \right)_{J_c=0} \quad \text{Thermal Conductivity}$$

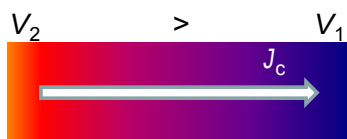
Wiedemann-Franz Law

$$\frac{\kappa}{\sigma} = LT$$

$$L = \frac{\pi^2}{3} \left(\frac{k_B}{e} \right)^2 = 2.4 \times 10^{-8} \text{ W}\Omega\text{K}^{-2}$$

Thermoelectricity

Electrical vs thermal transport



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Wiedemann-Franz Law

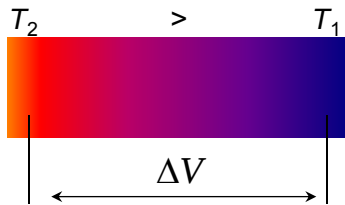
$$\frac{\kappa}{\sigma} = LT$$

$$L = \frac{\pi^2}{3} \left(\frac{k_B}{e} \right)^2 = 2.4 \times 10^{-8} \text{ W}\Omega\text{K}^{-2}$$

Thermoelectricity

Seebeck effect

Conversion of temperature differences directly into electricity



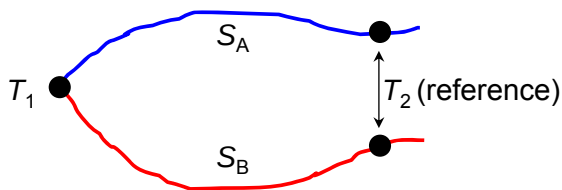
$$S = \left(\frac{\Delta V}{\Delta T} \right)_{J_c=0}$$

Seebeck coefficient



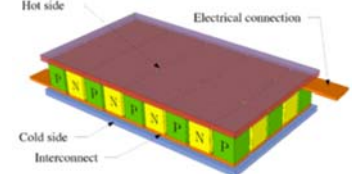
Thomas J. Seebeck
(1770-1831)

Thermocouple



$$\Delta V = \int_{T_1}^{T_2} [S_B(T) - S_A(T)] dT$$

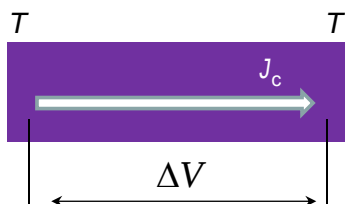
$$\Delta V \approx (S_B - S_A) \Delta T$$



Thermoelectricity

Peltier effect

Conversion of charge current into heat flow



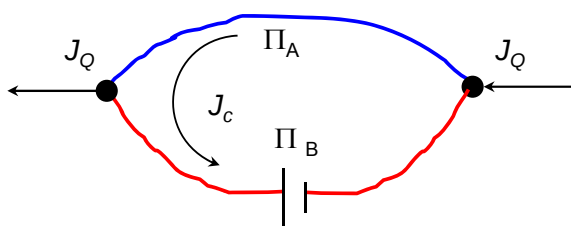
$$\Pi = \left(\frac{J_Q}{J_c} \right)_{\Delta T=0}$$

Seebeck coefficient

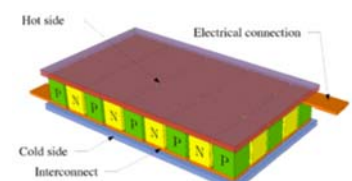


Jean C. A. Peltier
(1785-1845)

Heat pump: cooler / heater

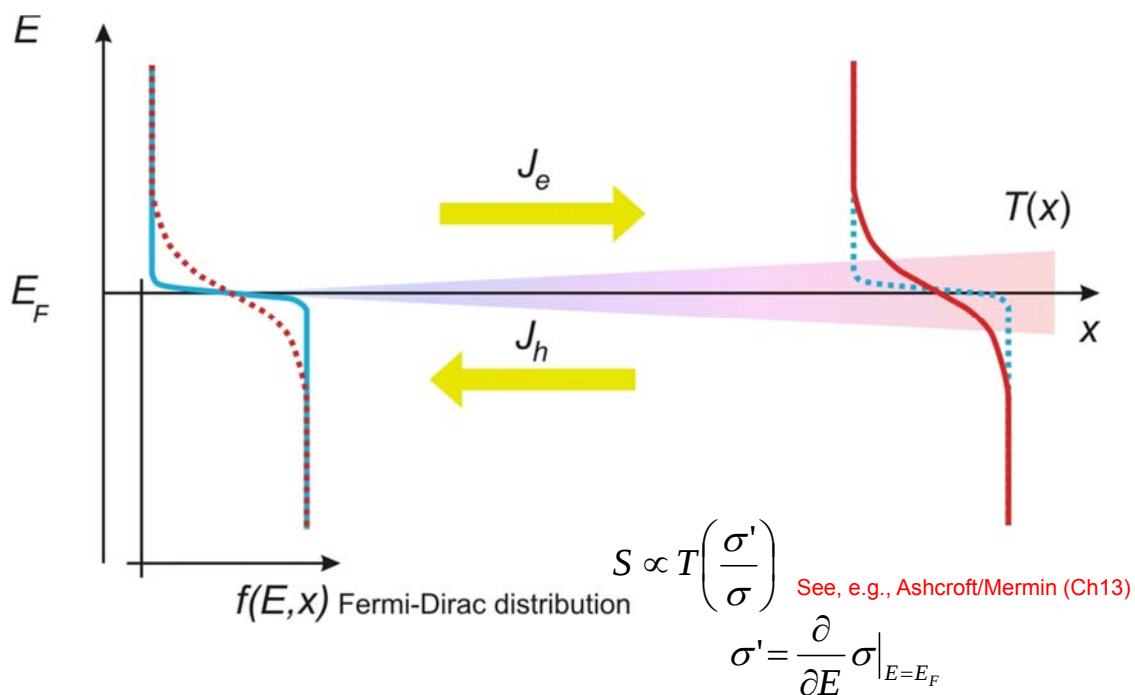


$$J_Q \approx (\Pi_B - \Pi_A) J_c$$



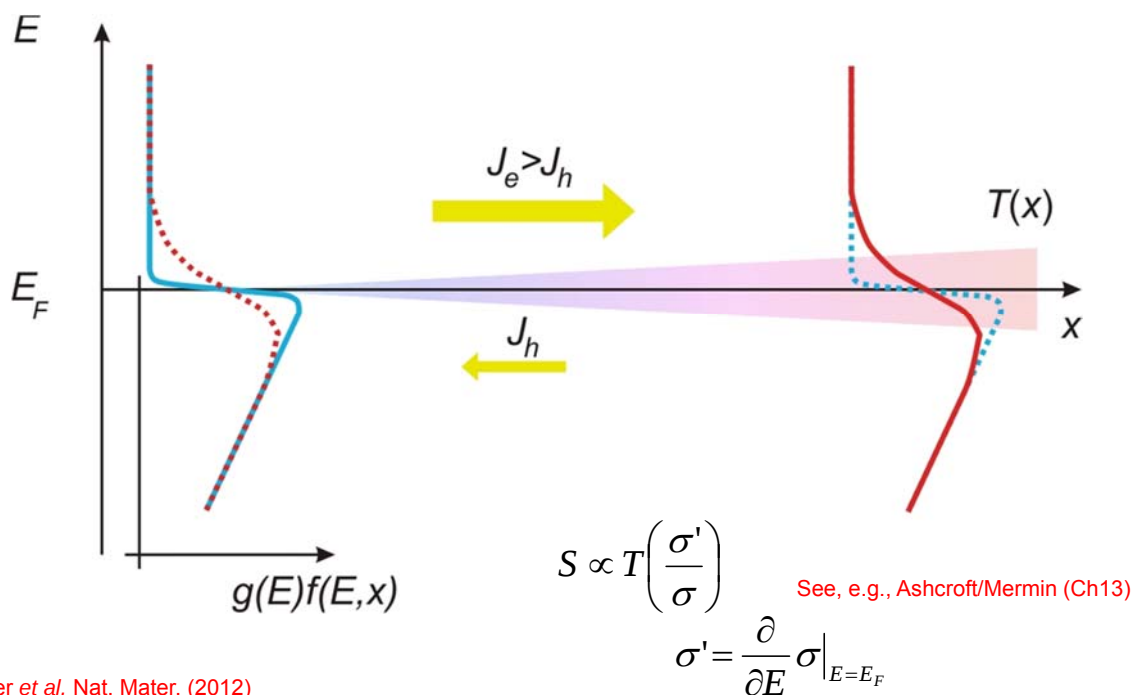
Thermoelectricity

Metal



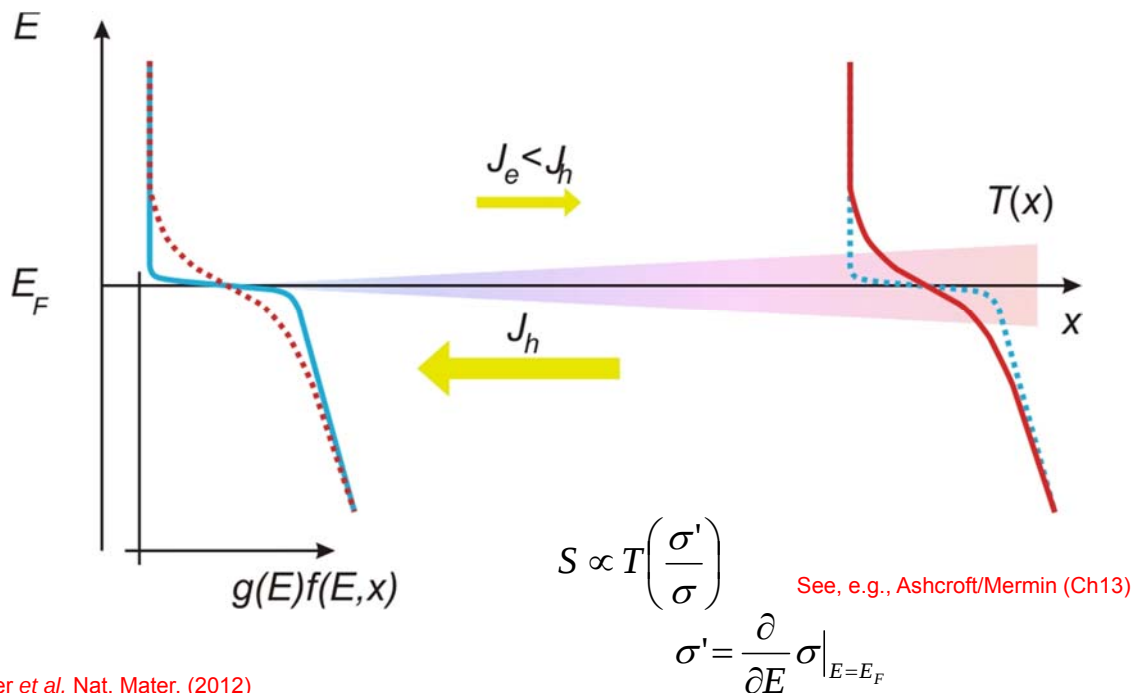
Thermoelectricity

Heat and charge transport (electron like)



Thermoelectricity

Heat and charge transport (electron like)



G. Bauer et al. Nat. Mater. (2012)

Thermoelectricity

Onsager reciprocity

The Nobel Prize in Chemistry 1968:

“for the discovery of the reciprocal relations bearing his name, which are fundamental for the thermodynamics of irreversible processes”

F_n Generalized force

J_n Generalized current

If $J_i = \sum_j L_{ij} F_j$ Linear response

and $\frac{dS}{dt} = \sum_j F_j J_j$ entropy change rate

then $L_{ij} = L_{ji}$



Lars Onsager
(1906-1976)

First law of thermodynamics: identify F_n , J_n

$F_n : \Delta V, -\Delta T/T$

then $\Pi = ST$ (Thomson relation)

$$\begin{pmatrix} J_c \\ J_Q \end{pmatrix} = \sigma \begin{pmatrix} 1 & S \\ \Pi & \kappa / \sigma \end{pmatrix} \begin{pmatrix} \Delta V \\ -\Delta T \end{pmatrix}$$

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Spin currents vs. charge currents

Charge

$$\vec{j}_e = \frac{d}{dt}(q\vec{r})$$

$$\vec{j}_e = q\vec{v}$$

Moving charge, kinetic energy and dissipation

Spin

$$\vec{j}_s = \frac{d}{dt}(\sigma\vec{r})$$

$$\vec{j}_s = \sigma\vec{v} + \dot{\sigma}\vec{r}$$

Spin currents are even under time reversal

Spins in motion (independent electron)

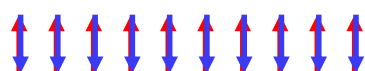


Electrons do not need to move

Nonvolatile memory

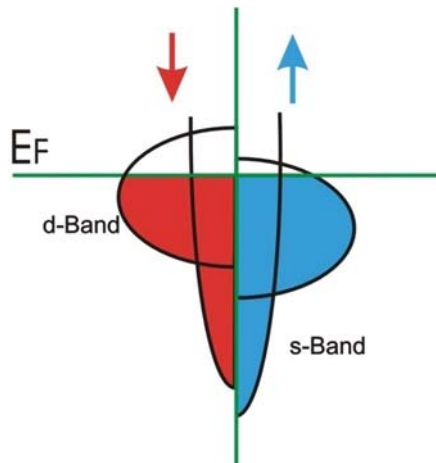
Possibility to reduce dissipation

Spin dynamics (collective)



Spin generation and spin injection

- Two spin channel model (Mott 1930)
 - Metallic ferromagnets. Spin-up and spin-down are two independent families of carriers



Spin splitting

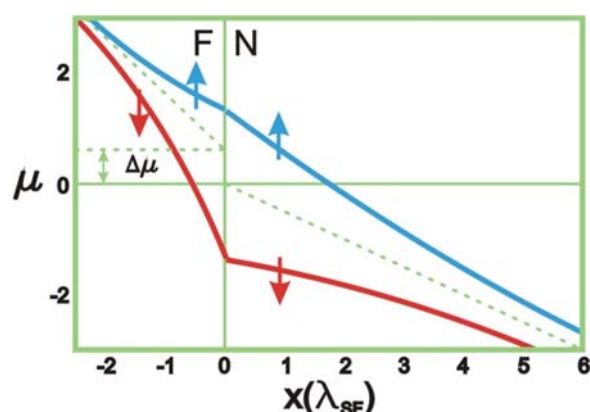
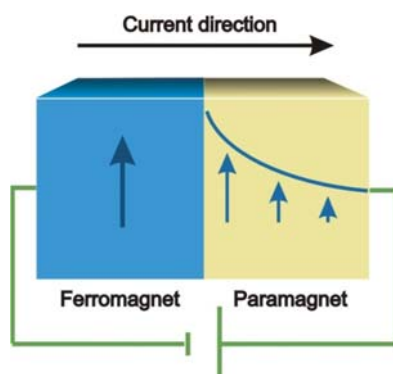
- Different density of states at the Fermi level for spin up and down carriers
- Different mobility for spin up and down carriers

$$P = \frac{N_M - N_m}{N_M + N_m} \quad -1 \leq P \leq 1$$

I.I. Mazin, PRL **83**, 1427-1430 (1999)

Spin currents and spin accumulation

- Spin polarized current in a nonmagnetic metal
- Spin accumulation decays exponentially
- Characteristic length. Spin diffusion/relaxation length λ_{sf}



$$\nabla^2(\mu^\uparrow - \mu^\downarrow) = \frac{1}{\lambda_{sf}}(\mu^\uparrow - \mu^\downarrow)$$

Johnson and Silsbee PRB **35**, 4959 (1987)
van Son et al., PRL **58**, 2271 (1987)

Spin dependent transport

In general the spin conductivities, $\sigma^\uparrow$, $\sigma^\downarrow$ are not equal

$$\begin{pmatrix} J^\uparrow \\ J^\downarrow \end{pmatrix} = \begin{pmatrix} \sigma^\uparrow & 0 \\ 0 & \sigma^\downarrow \end{pmatrix} \begin{pmatrix} \Delta\mu^\uparrow / e \\ \Delta\mu^\downarrow / e \end{pmatrix}$$

thus

$$\begin{pmatrix} J_c \\ J_s \end{pmatrix} = \sigma \begin{pmatrix} 1 & P \\ P & 1 \end{pmatrix} \begin{pmatrix} \Delta\mu_c / e \\ \Delta\mu_s / e \end{pmatrix}$$

where

$$J_c = J^\uparrow + J^\downarrow \quad \text{Charge current}$$

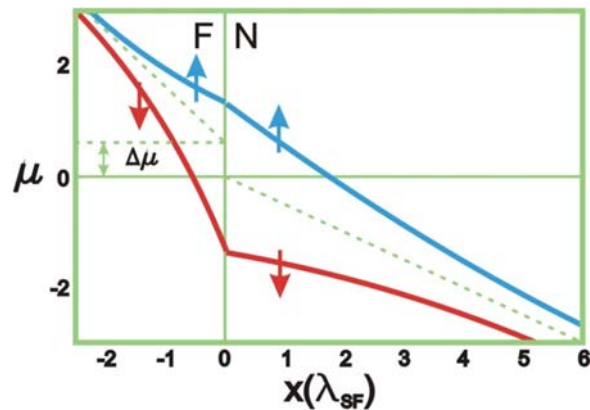
$$J_s = J^\uparrow - J^\downarrow \quad \text{Spin current}$$

$$P = \frac{\sigma^\uparrow - \sigma^\downarrow}{\sigma^\uparrow + \sigma^\downarrow} \quad \text{Spin polarization}$$

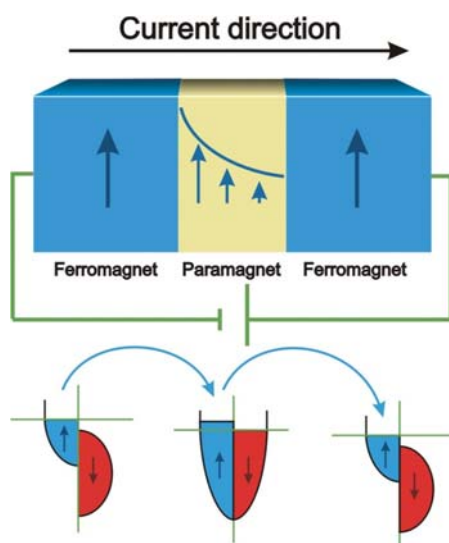
$$\mu_c = \frac{\mu^\uparrow + \mu^\downarrow}{2} \quad \text{Charge chemical potential}$$

$$\mu_s = \mu^\uparrow - \mu^\downarrow \quad \text{Spin accumulation}$$

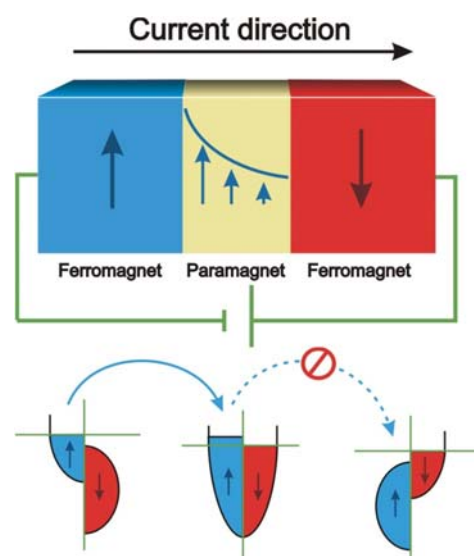
$$\sigma = \sigma^\uparrow + \sigma^\downarrow \quad \text{Charge conductance}$$



Two terminal spin electronics (GMR)

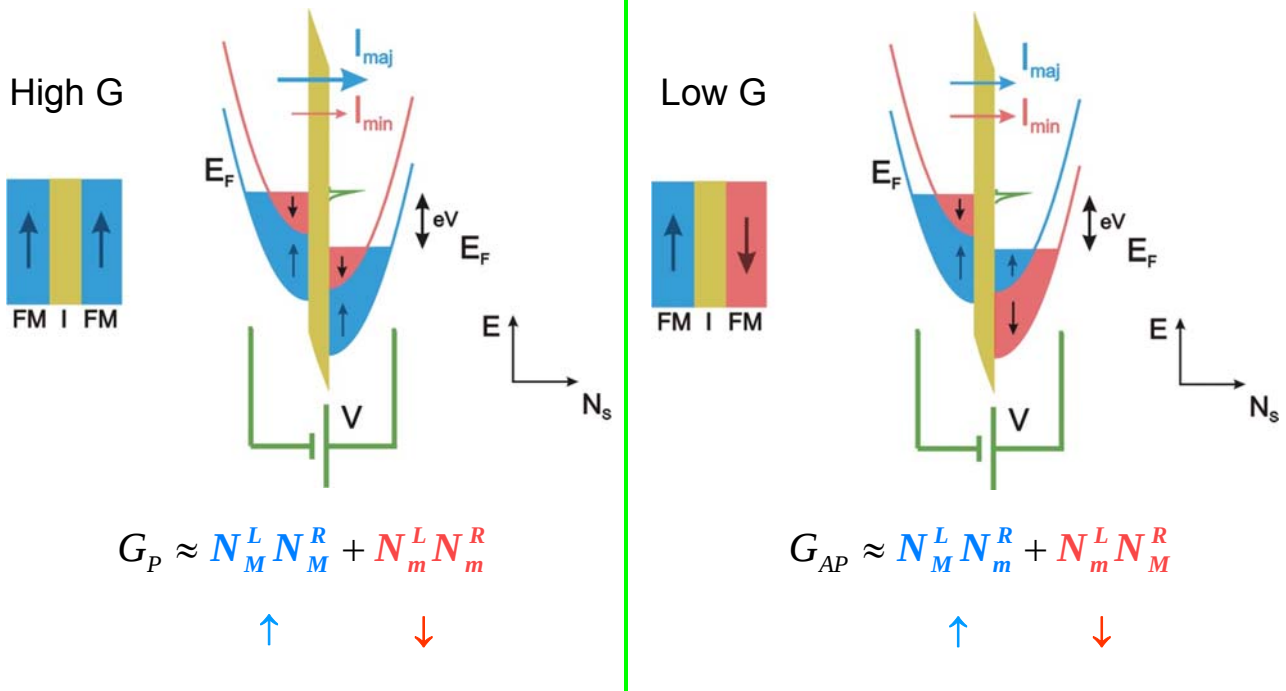


Parallel Magnetization, $\uparrow\uparrow$
High conductance



Antiparallel Magnetization, $\uparrow\downarrow$
Low conductance

Two terminal spin electronics (TMR)



Two terminal spin electronics

- Magnetic Tunnel Junctions (MTJ).

$$TMR = \frac{R_{AP} - R_P}{R_P} = \frac{G_P - G_{AP}}{G_{AP}} = \frac{2P_L P_R}{1 - P_L P_R}$$

$$G_P \approx N_M^L N_M^R + N_m^L N_m^R$$

$$G_{AP} \approx N_M^L N_m^R + N_m^L N_M^R$$

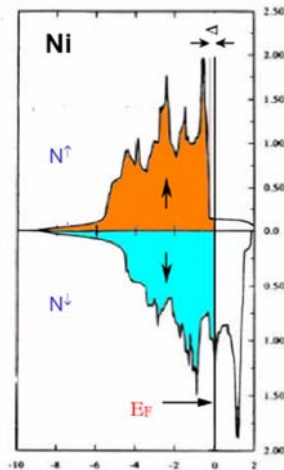
$$P_{L,R} = \frac{N_M^{L,R} - N_m^{L,R}}{N_M^{L,R} + N_m^{L,R}}$$

Jullière, Phys. Lett. **54A**, 225 (1975)

Spin dependent tunneling. Density of states

How is P defined ?

$$TMR = \frac{R_{AP} - R_P}{R_P} = \frac{G_P - G_{AP}}{G_{AP}} = \frac{2P_L P_R}{1 - P_L P_R}$$



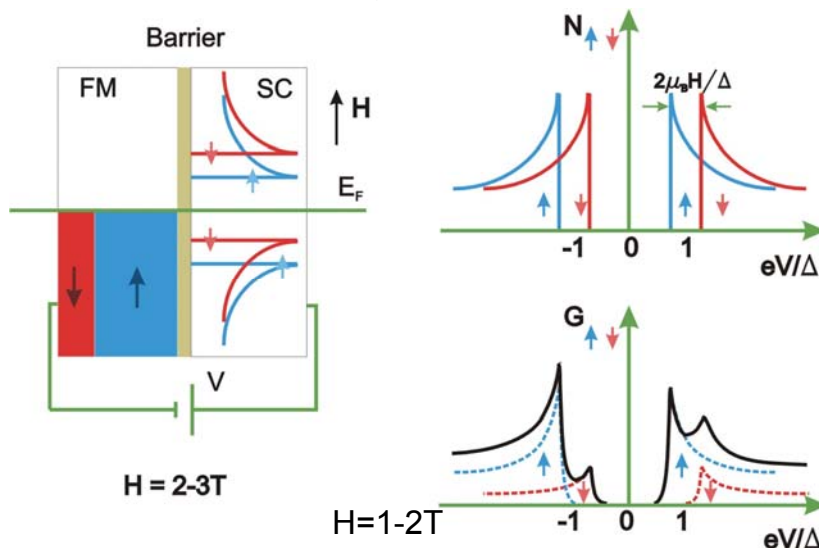
- Julliere: P equals the spin polarization of the electrons at the Fermi level
- From energy band calculations $P < 0$ for Ni and $P > 0$ for Fe
- Experimentally $TMR > 0$ for Ni-I-Fe and bias dependent
- Tunneling probability has to be taken into account (tunneling matrix)
- Current and tunneling mediated by free-like electrons
- Interface. Symmetry states

Stearns, J. Magn. Magn. Mat. **5**, 167 (1977)

Polarization measurement

Meservey-Tedrow

- Meservey-Tedrow technique. Superconductor with Zeeman split density of states as a spin detector
 - P is obtained at high field, low temperatures and zero bias



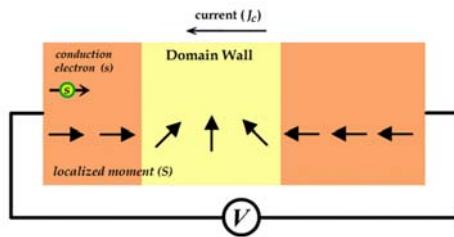
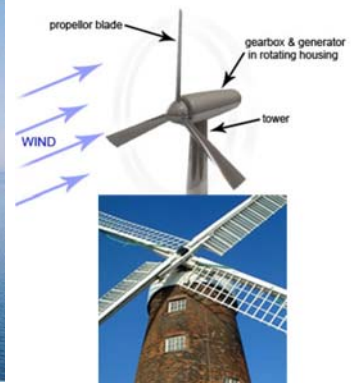
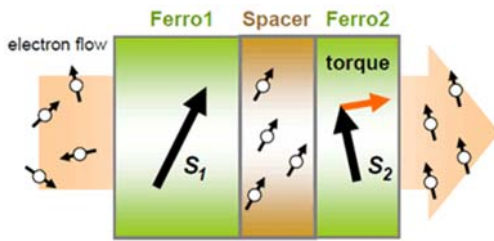
Partially polarized materials:
Fe, Co, Ni
($P \sim 25-45\%$)

Meservey and Tedrow, Phys. Rep. **238**, 173 (1994)

Spin Torque vs spin pumping

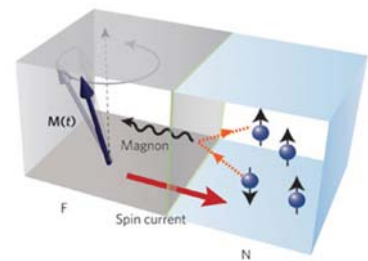
Onsager reciprocity

Michel Viret talk (yesterday)



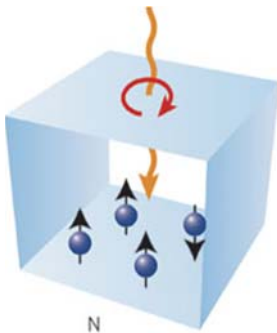
S. Maekawa (Tohoku and JAEA)

Spin pumping

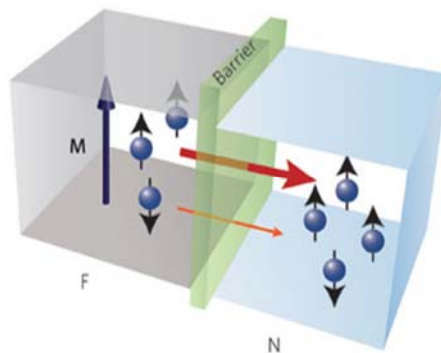


Spin current generation and accumulation

Optical Orientation



Spin Injection



Spin Pumping

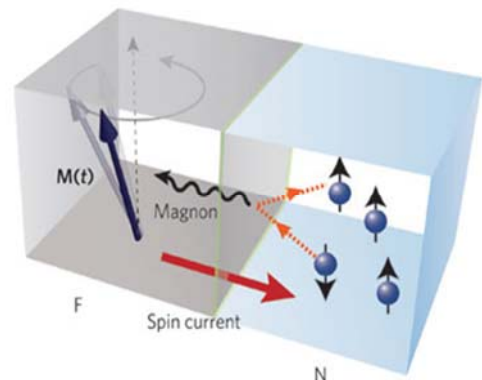
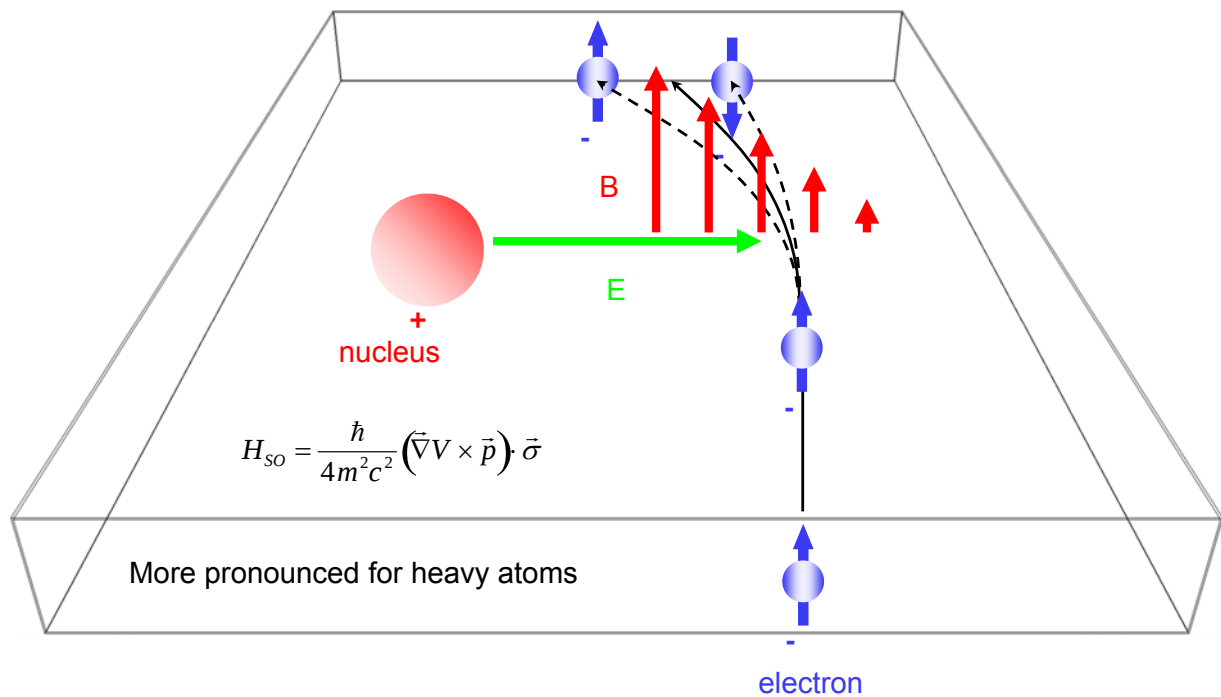


Figure from Zutic and Dery Nat. Mater (2011)

Spin Hall effect

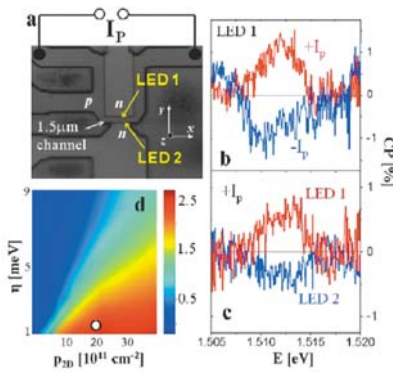
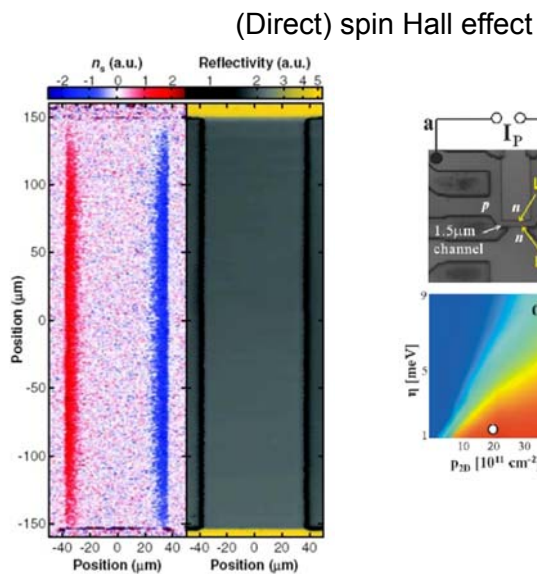
Skew Scattering



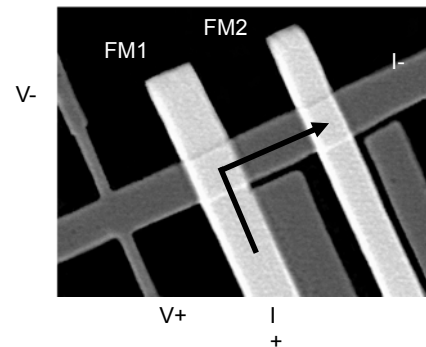
Axel Hoffmann, Argonne National Laboratory, US.

Spin Hall Effect

Pure spin currents



Inverse spin Hall effect
(follows for Onsager reciprocity)



Valenzuela and Tinkham Nature (2006)

Y.K. Kato *et al.* Science (2004); J. Wunderlich, PRL (2005).

V. Sih *et al.* Nature Physics (2005).

Spin currents and spin caloritronics

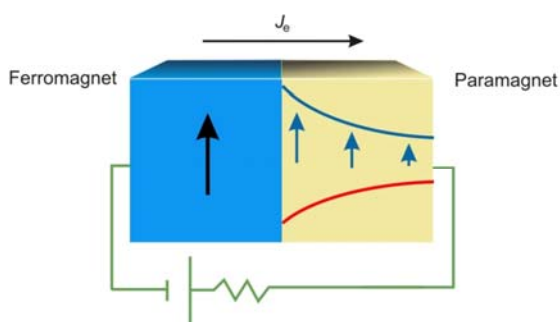
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Spin currents and spin caloritronics

Thermal spin injection

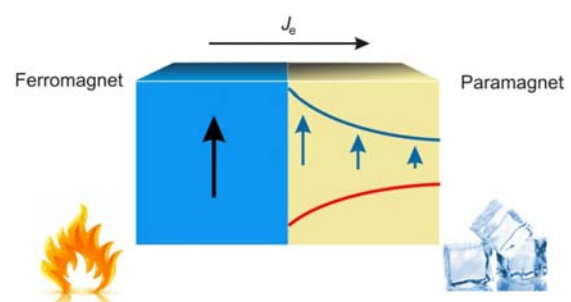


$$\begin{pmatrix} J_c \\ J_s \end{pmatrix} = \sigma \begin{pmatrix} 1 & P \\ P & 1 \end{pmatrix} \begin{pmatrix} \Delta\mu_c / e \\ \Delta\mu_s / e \end{pmatrix}$$

$$\begin{pmatrix} J_c \\ J_Q \end{pmatrix} = \sigma \begin{pmatrix} 1 & S \\ \Pi & \kappa / \sigma \end{pmatrix} \begin{pmatrix} \Delta V \\ -\Delta T \end{pmatrix}$$

$$S \propto T \left(\frac{\sigma'}{\sigma} \right) \quad \sigma' = \frac{\partial}{\partial E} \sigma \Big|_{E=E_F}$$

$$\Pi = ST \text{ (Thomson relation)}$$



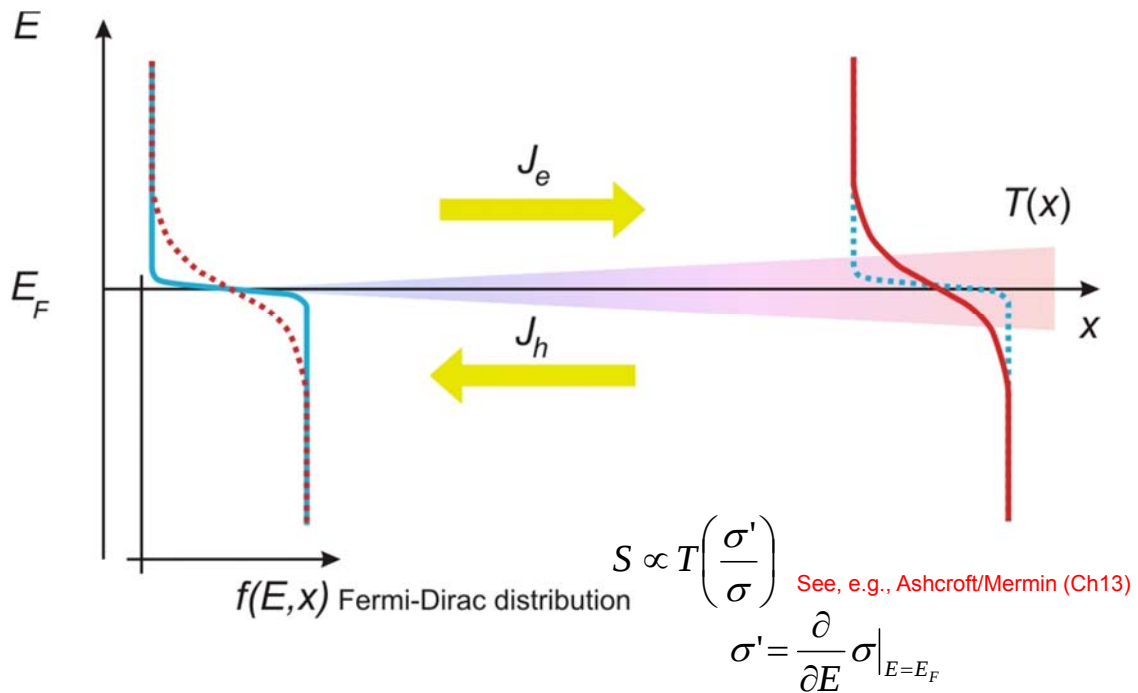
$$\begin{pmatrix} J_c \\ J_s \\ J_Q \end{pmatrix} = \sigma \begin{pmatrix} 1 & P & S \\ P & 1 & P'S \\ ST & P'ST & \kappa / \sigma \end{pmatrix} \begin{pmatrix} \Delta\mu_c / e \\ \Delta\mu_s / e \\ -\Delta T \end{pmatrix}$$

$$P = \frac{\sigma^\uparrow - \sigma^\downarrow}{\sigma^\uparrow + \sigma^\downarrow} \quad P' = \frac{\partial_\epsilon \sigma^\uparrow - \partial_\epsilon \sigma^\downarrow}{\partial_\epsilon \sigma^\uparrow + \partial_\epsilon \sigma^\downarrow} \Big|_{\epsilon_F}$$

M. Johnson and R. H. Silsbee, Phys. Rev. B **35**, 4959 (1987)

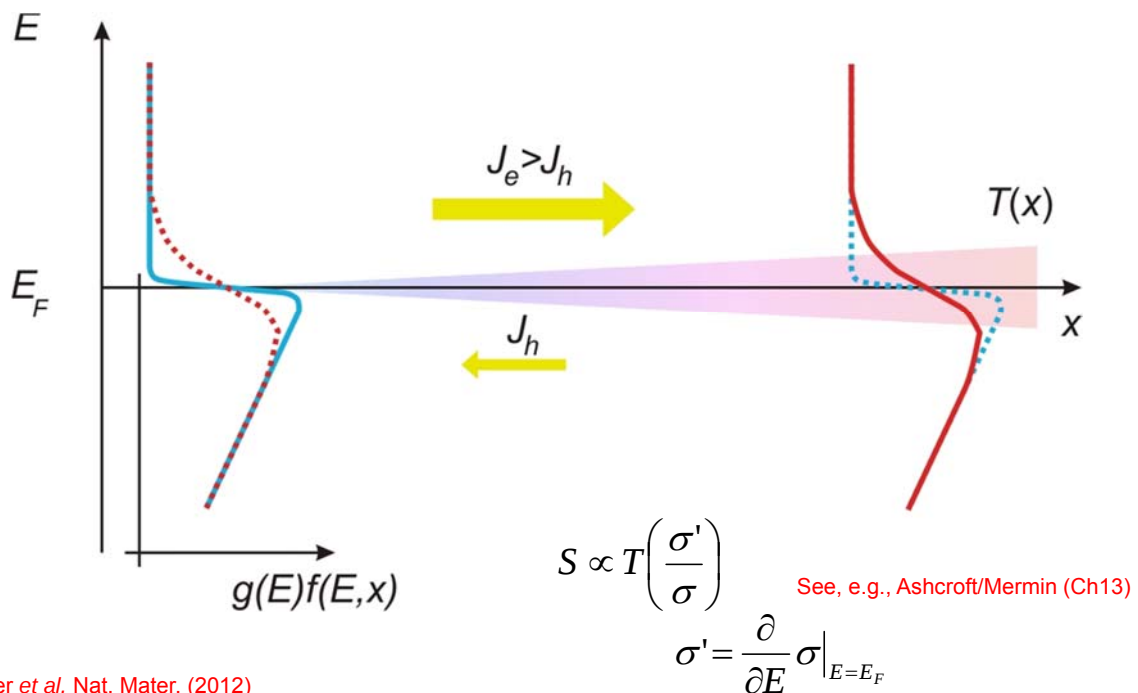
Thermoelectricity

Metal



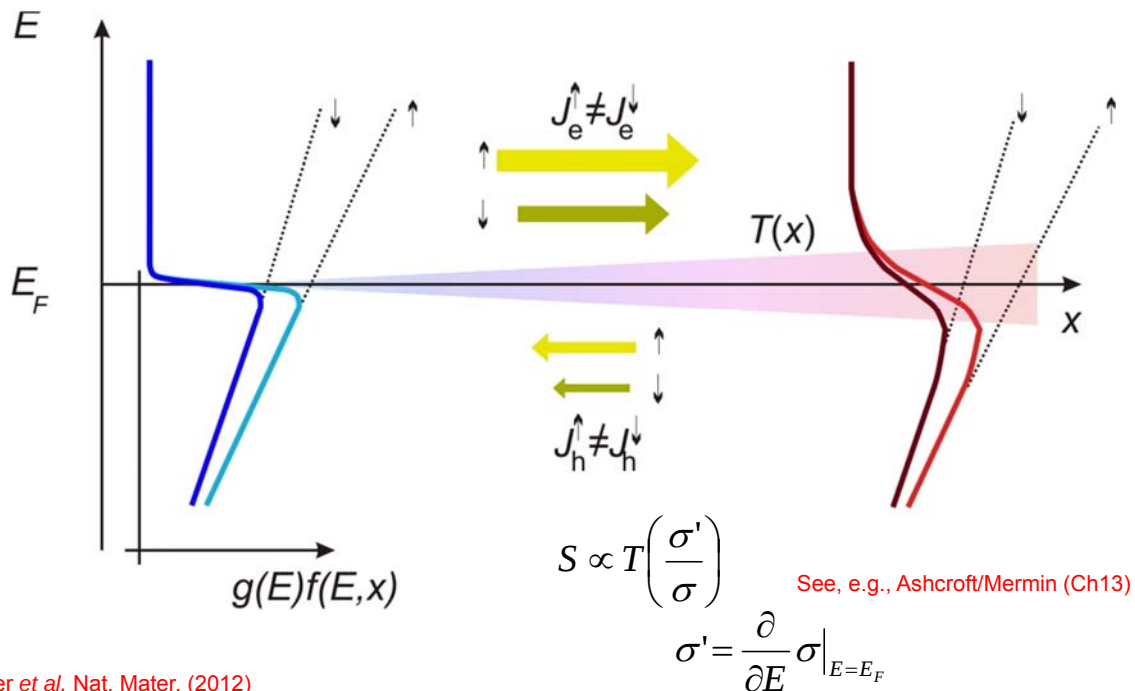
Thermoelectricity

Heat and charge transport (electron like)



Thermoelectricity + Spin

Heat, spin and charge transport (electron like)



G. Bauer et al. Nat. Mater. (2012)

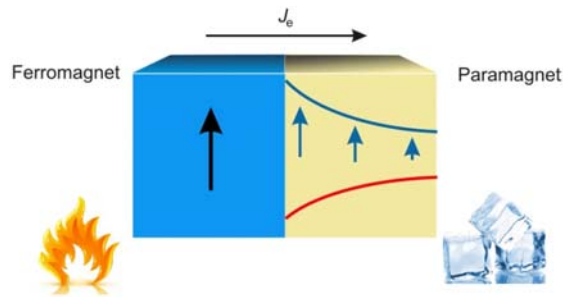
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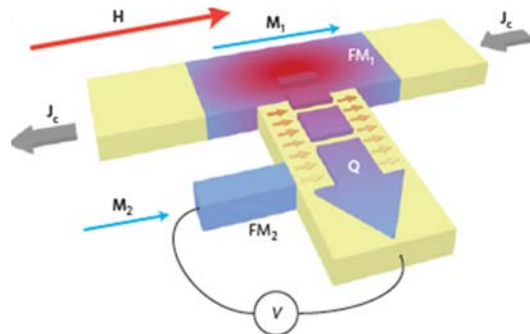
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Spin-dependent Seebeck effect



$$\begin{pmatrix} J_c \\ J_s \\ J_Q \end{pmatrix} = \sigma \begin{pmatrix} 1 & P & S \\ P & 1 & P'S \\ ST & P'ST & \kappa/\sigma \end{pmatrix} \begin{pmatrix} \Delta\mu_c/e \\ \Delta\mu_s/e \\ -\Delta T \end{pmatrix}$$



Ferromagnetic injector: Permalloy (NiFe)

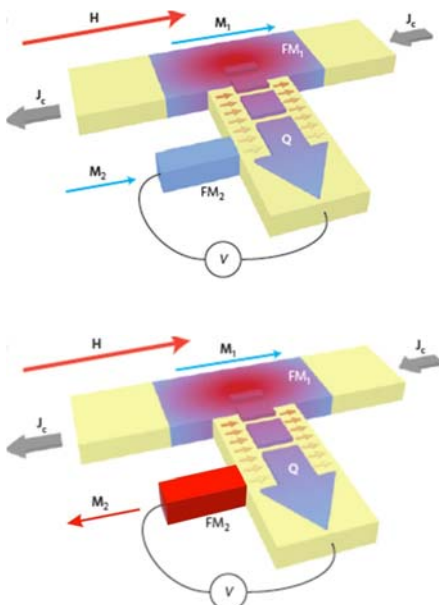
Paramagnetic material: Cu

$$S_{Py} = -20 \text{ uW/K}$$

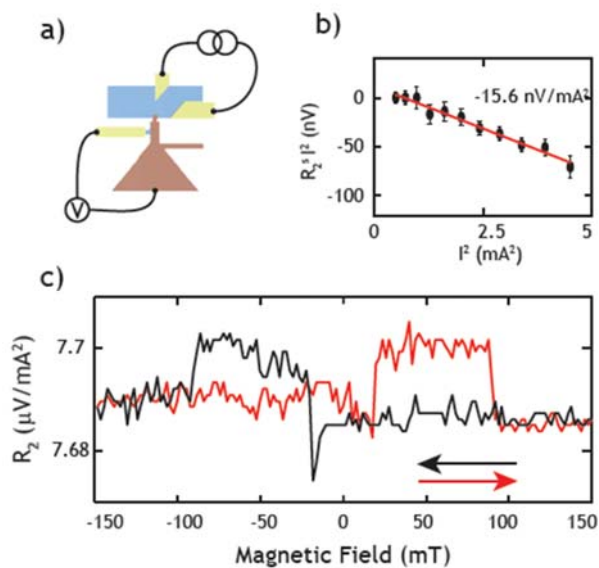
$$S_{Cu} = 1.6 \text{ uW/K}$$

Slachter *et al.* Nature Phys. (2010)

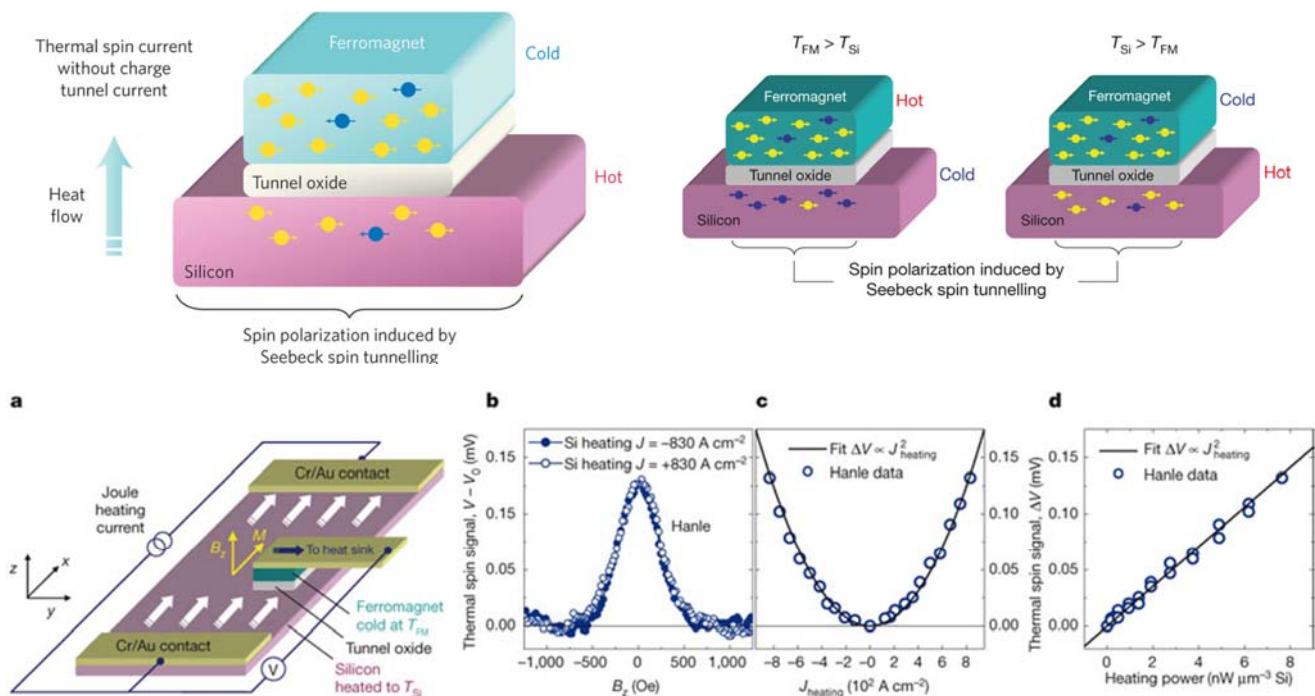
Spin-dependent Seebeck effect



Slachter *et al.* Nature Phys. (2010)



Spin-dependent Seebeck tunneling

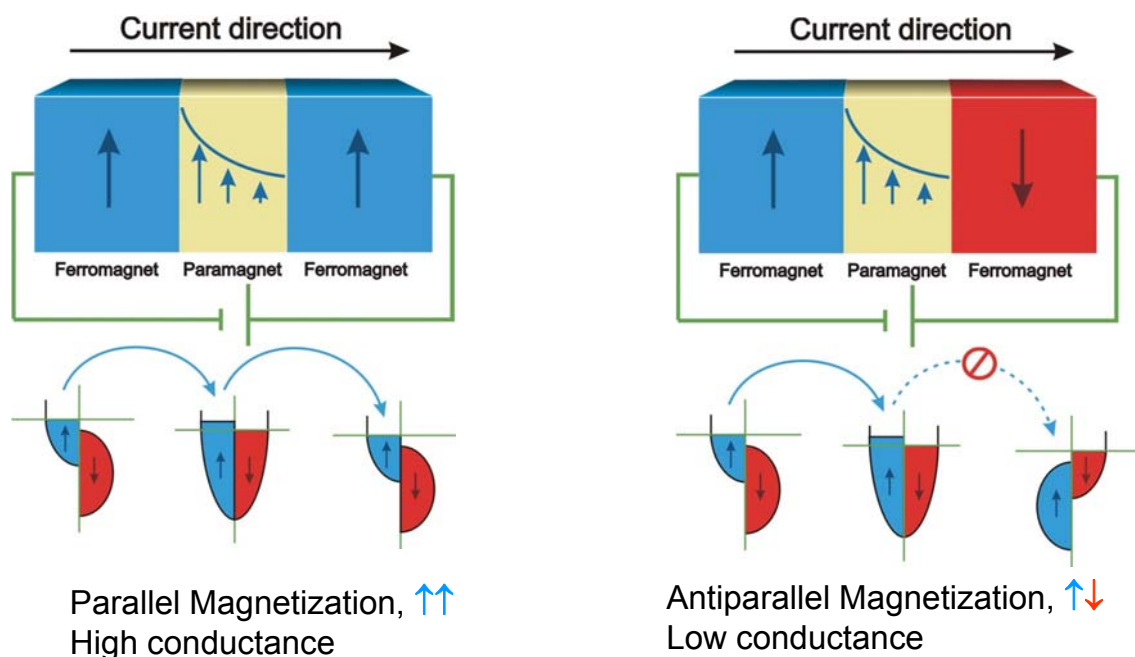


Le Breton et al. Nature (2011)

European School on Magnetism 2013, Cargèse, Corsica

Spin Currents and Spin Caloritronics

Two terminal spin electronics (GMR)



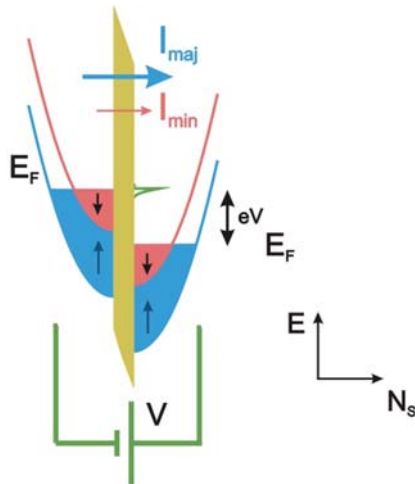
Fert, Grünberg. Nobel Prize 2007

European School on Magnetism 2013, Cargèse, Corsica

Spin Currents and Spin Caloritronics

Two terminal spin electronics (TMR)

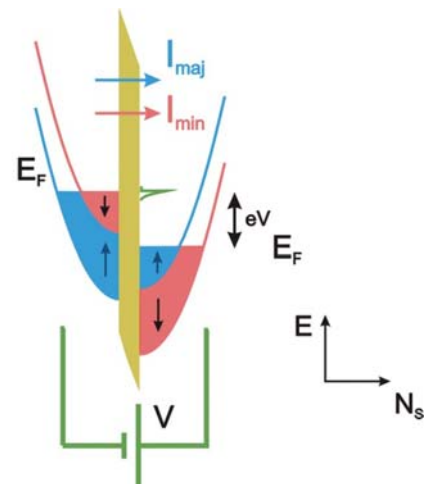
High G



$$G_P \approx N_M^L N_m^R + N_m^L N_m^R$$



Low G

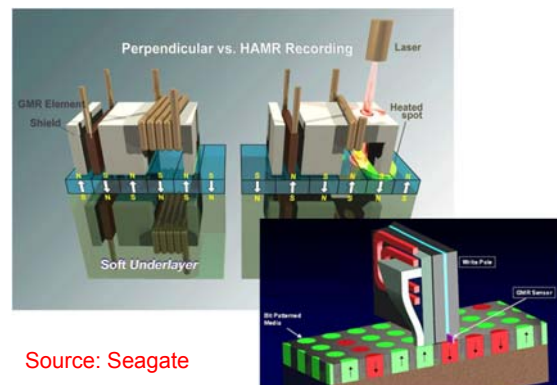
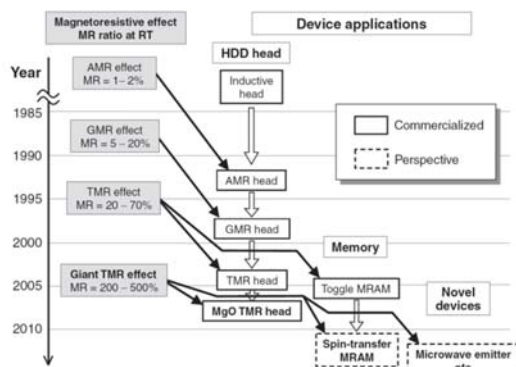


$$G_{AP} \approx N_M^L N_m^R + N_m^L N_M^R$$



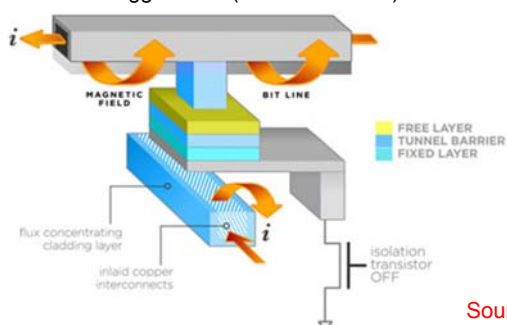
Spin valves

Magnetic Access Random Memory (MRAM) Current developments



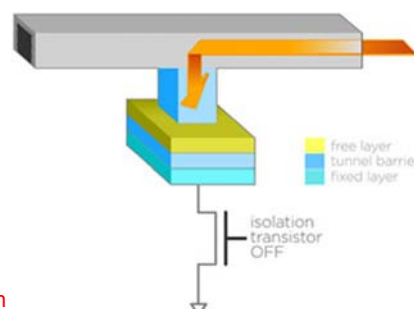
Source: Seagate

Toggle RAM (commercialized)



Source: Everspin

Spin Transfer Torque RAM (under development)

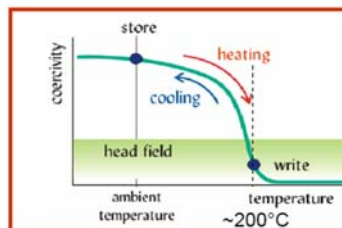


Spin valves

Magnetic Access Random Memory (MRAM) Current developments



Thermally Assisted Switching (TAS) MRAM



• Use temperature-dependence of switching field

- Write at elevated temperature
- Store / read at room temperature

• Same basic concept as in Heat Assisted Magnetic Recording.
• Here heating produced by Joule dissipation around tunnel barrier.

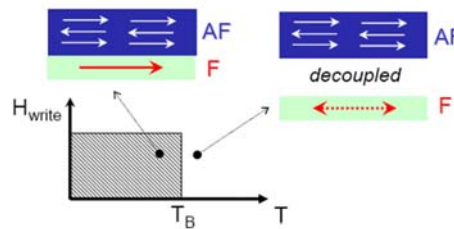
❖ Use **exchange biased storage bilayer**

to "lock" the stored data:

→ **high stability**

❖ Locally **heat the magnetic cell** above its "blocking temperature" to "unlock" the storage layer :

→ **selectivity**
reduced write power



8 Nov 2010

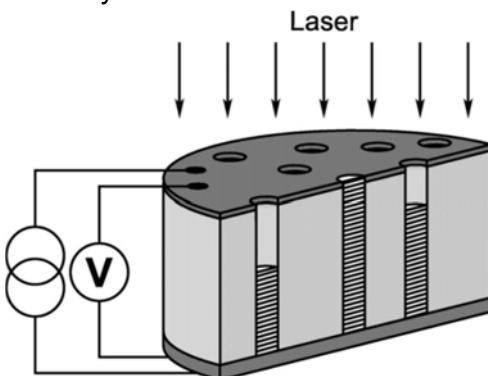
Bernard Dieny (SPINTEC/CROCUS)



Spin-dependent thermoelectric effects in spin valves (2006)

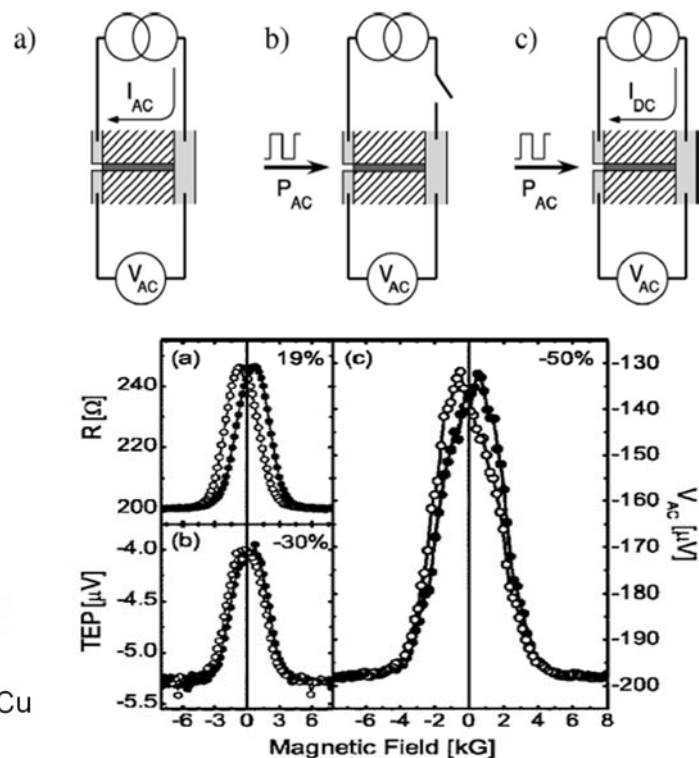
J. Shi, K. Pettit, *et al.* Phys. Rev. B (1996)
McCann, E. & Fal'ko, Appl. Phys. Lett. (2002).

thermoelectric power and thermal conductivity on granular and multilayer GMR systems

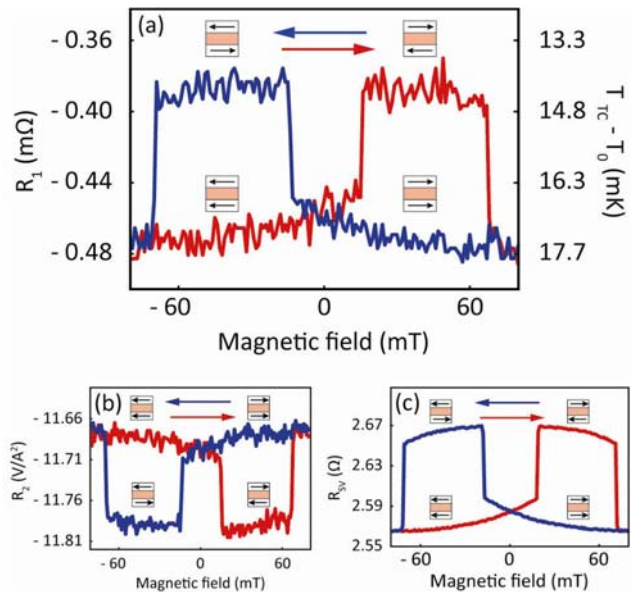
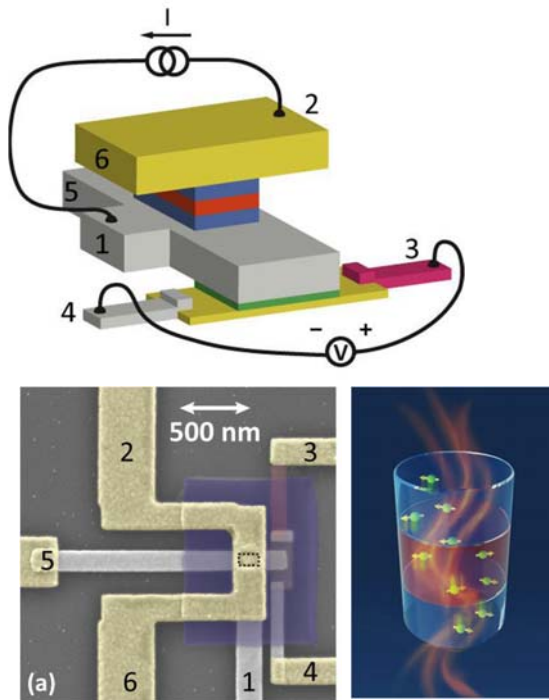


L. Gravier, S. Serrano-Guisan, *et al.*, Phys. Rev. B (2006)

Spin-dependent heat and charge transport perpendicular to the plane of magnetic Co/Cu multilayers. Peltier effect.



Spin-dependent Peltier effect

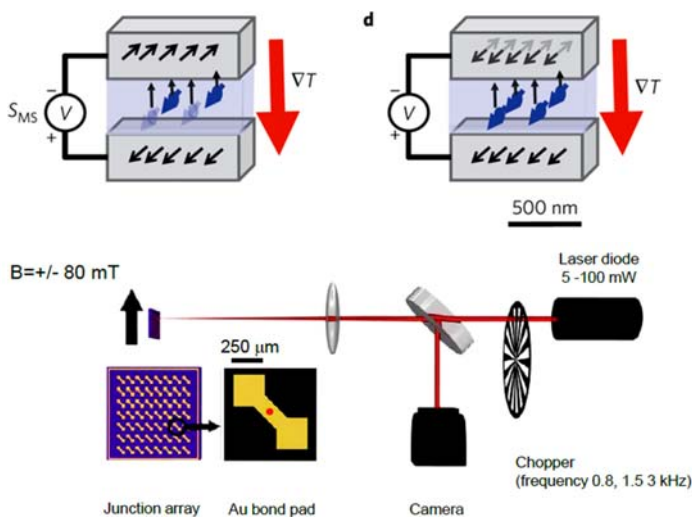


Flipse *et al.* Nature Nanotechnol. (2012)

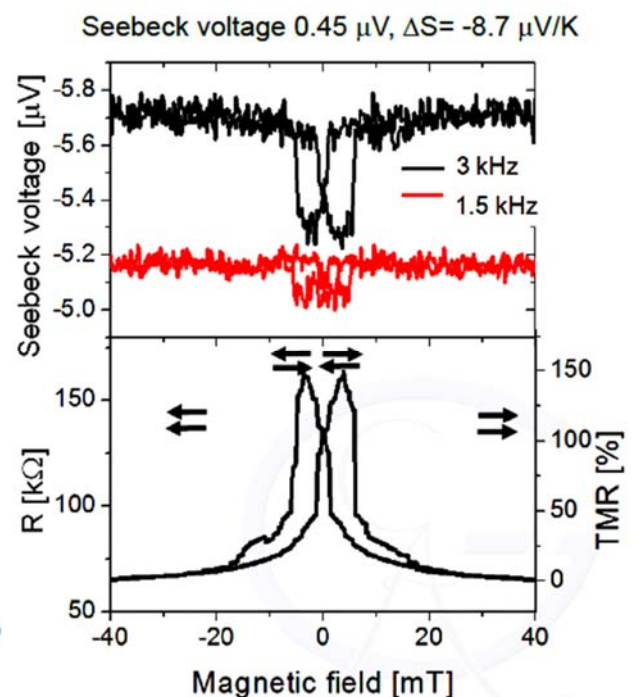
Magneto Seebeck effect tunnel junctions

Magneto Seebeck effect -8.8%

$$S_{MS} = \frac{S_P - S_{AP}}{\min(S_P, S_{AP})}$$



Walter *et al.* Nature Mater. (2011)



Spin currents vs. charge currents

Charge

$$\vec{j}_e = \frac{d}{dt}(q\vec{r})$$

$$\vec{j}_e = q\vec{v}$$

Moving charge, kinetic energy and dissipation

Spin

$$\vec{j}_s = \frac{d}{dt}(\sigma\vec{r})$$

$$\vec{j}_s = \sigma\vec{v} + \dot{\sigma}\vec{r}$$

Spin currents are even under time reversal

Spins in motion (independent electron)

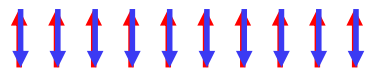


Electrons do not need to move

Nonvolatile memory

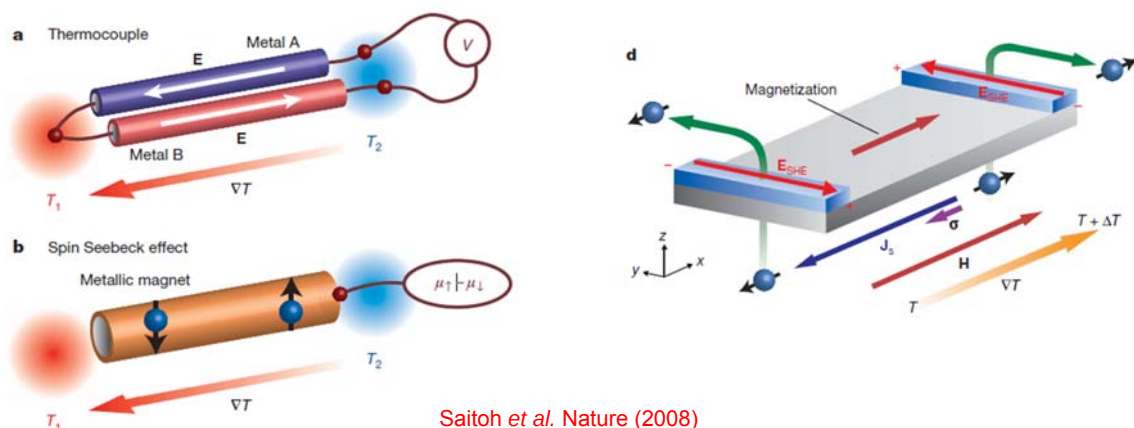
Possibility to reduce dissipation

Spin dynamics (collective)



J. Shi, et al., Phys. Rev. Lett. **96**, 076604 (2006).

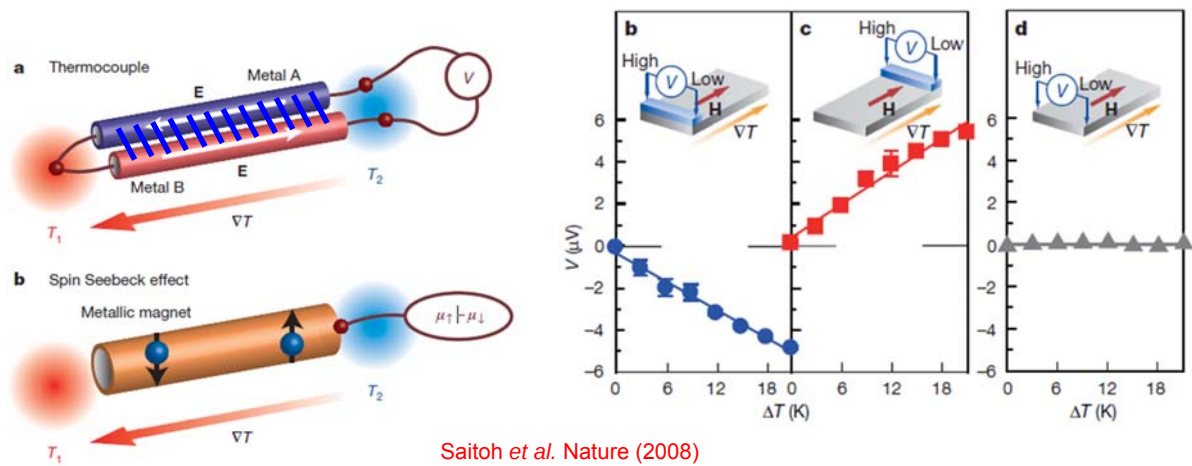
Spin Seebeck effect



Saitoh et al. Nature (2008)

Original idea: two spin channels acting as the two distinct materials in a thermocouple. A temperature gradient should result in a spin voltage proportional to the temperature difference Which can be detected by inverse spin Hall effect in Pt

Spin Seebeck effect



Original idea: two spin channels acting as the two distinct materials in a thermocouple. A temperature gradient should result in a spin voltage proportional to the temperature difference Which can be detected by inverse spin Hall effect in Pt

Reviews: Bauer, MacDonald, Maekawa, Solid State Commun. (2010); Bauer in Spin Current (Oxford University Press, 2012)

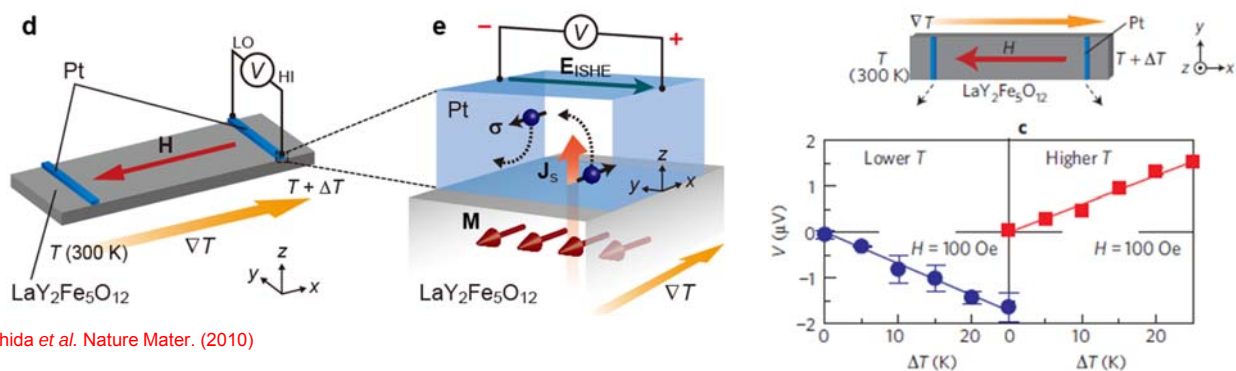
European School on Magnetism 2013, Cargèse, Corsica

Spin Currents and Spin Caloritronics

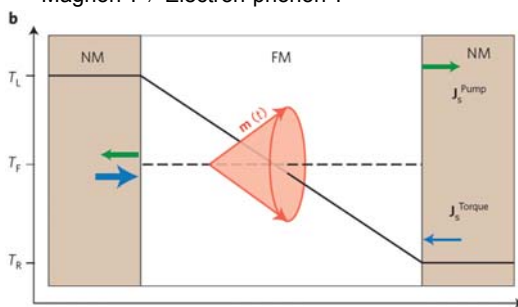
Spin Seebeck effect

Basic mechanism

Spin Seebeck in insulators



Magnon $T \neq$ Electron-phonon T



Uchida et al. Nature Mater. (2010); Xiao et al. , Phys. Rev. B (2010); Bauer et al. Nature Mater. (2012)

Spin pumping (SP) vs. Johnson-Nyquist noise

$$V_{ISHE} \propto J_S$$

$$J_S = J_S^{SP} - J_S^{J-N} = C(T_F^M - T_N^e)$$

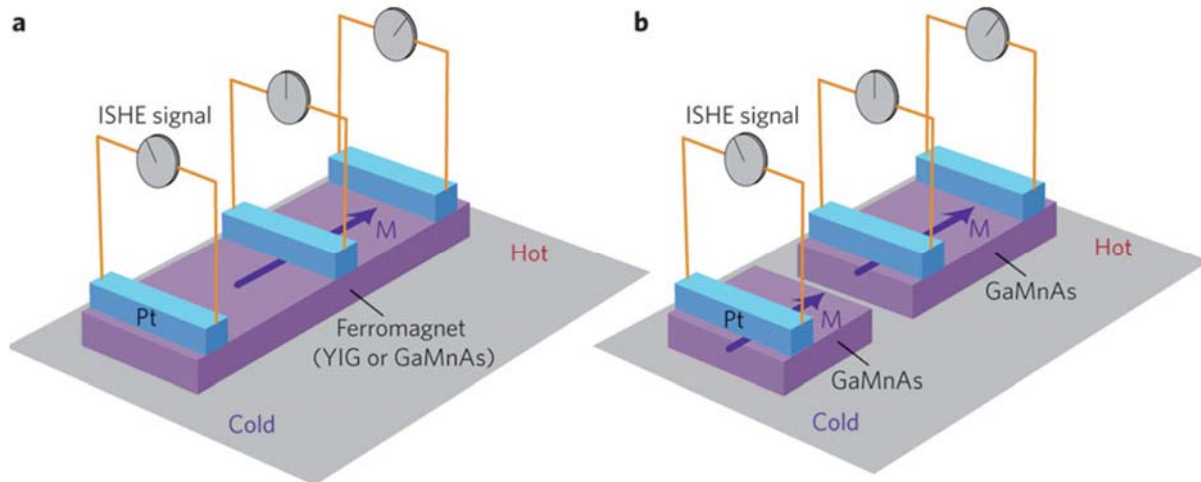
European School on Magnetism 2013, Cargèse, Corsica

Spin Currents and Spin Caloritronics

Spin Seebeck effect

Basic mechanism

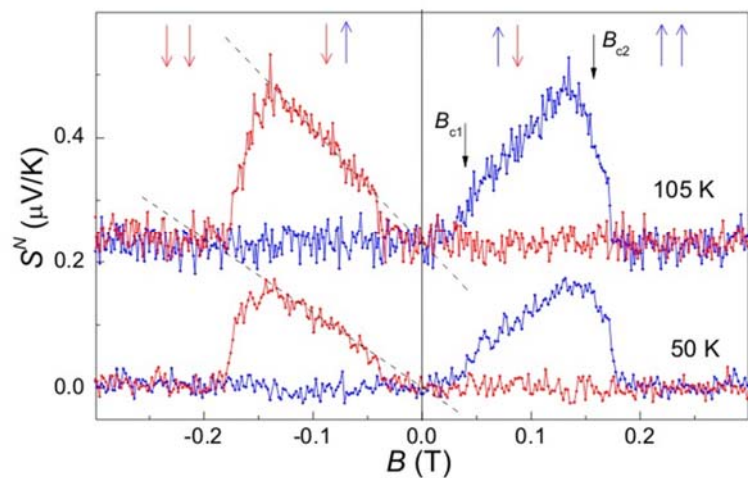
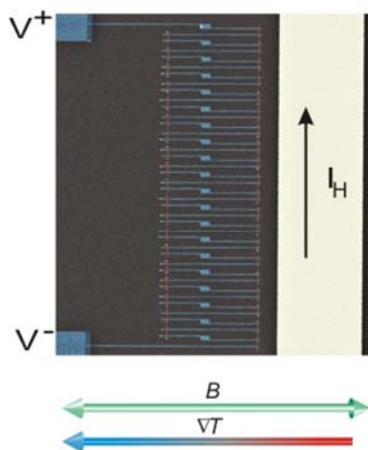
Spin Seebeck in semiconductors



Jaworski *et al.* Nature Mater. (2010); Sinova Nature Mater. (2010)

Thermopile

Thermopower



$$S^N = (V^+ - V^-) / \Delta T$$

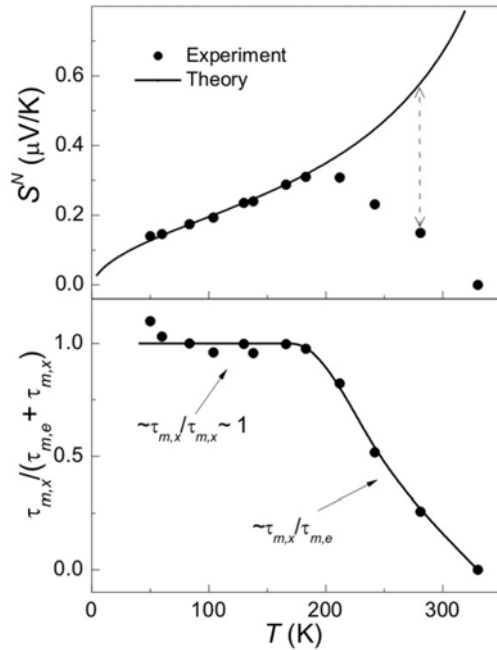
Signal only observed in the antiparallel configuration

Linear in B extrapolates to zero at $B = 0$

M.V. Costache, G. Bridoux, I. Neumann and SOV, Nature Mater. (2012)

Thermopile

Thermopower modeling



$$S^{MD} \propto P_{m,e}/(P_{m,x} + P_{m,e}) \propto \tau_{m,x}/(\tau_{m,x} + \tau_{m,e})$$

$P_{m,e}$: Probability magnon-electron interaction

$P_{m,x} + P_{m,e}$ Probability magnon interaction

Phonons: Ziman in *Electrons and phonons* (OUP, 1960)

$$S^{MD}(B, T) = \frac{1}{n'_e} \frac{k_B^{5/2} T^{3/2} F(y)}{6\pi^2 D^{3/2}} \left(\frac{\tau_{m,x}}{\tau_{m,x} + \tau_{m,e}} \right)$$

$$y = (g\mu_B B_{int}/k_B T) \quad n'_e = n_e/\alpha$$

$F(y)$ Quenching function

Grannemann and Berger, *Phys. Rev. B* (1976)

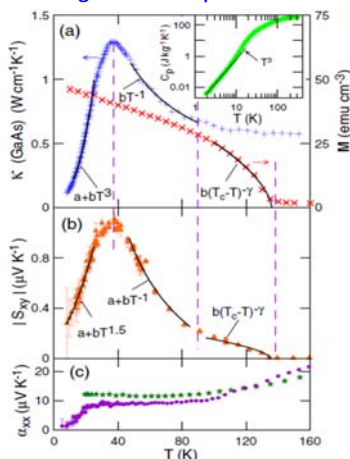
For low T : $\tau_{m,x}/(\tau_{m,x} + \tau_{m,e}) \approx 1$

$$S^N = V/\Delta T = N[S^{MD}(T, -B) - S^{MD}(T, +B)].$$

Spin Seebeck effect

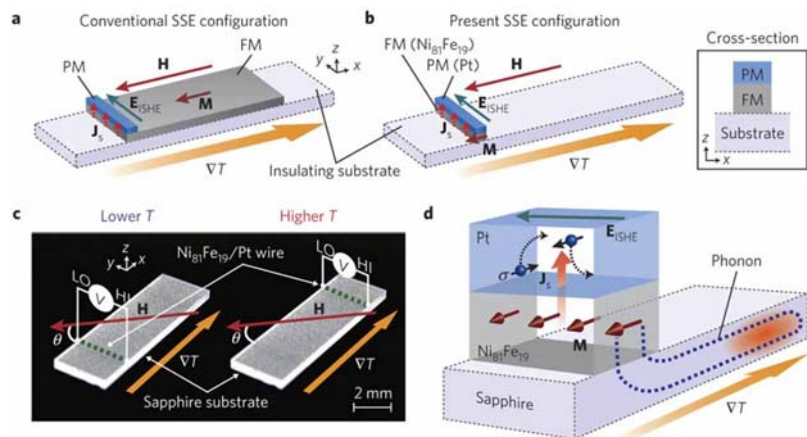
Basic mechanism

Magnons and phonons

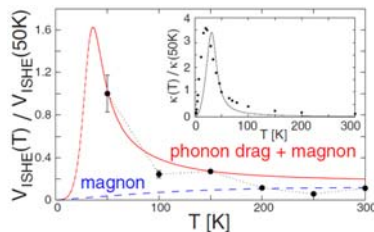


Jaworski et al., *Phys. Rev. Lett* (2011)

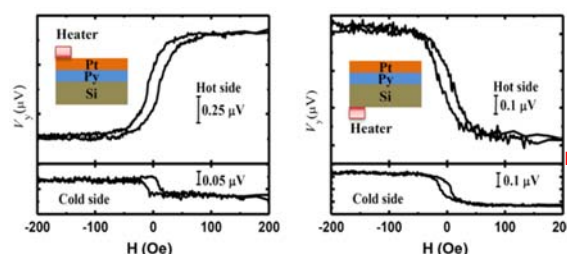
Acoustic spin pumping



Uchida et al. *Nature Mater.* (2011)



Adachi et al., *Appl. Phys. Lett* (2010)



Anomalous Nernst effect

Huang et al., *Phys. Rev. Lett* (2011)

Spin caloritronics

Spin caloritronics: interaction between spin and heat transport

Spin-dependent conductivities: Cause spin-dependent thermoelectric effects

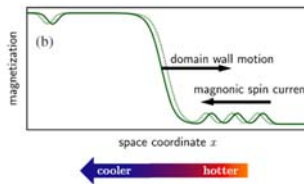
Collective spin dynamics: Associated to new thermomagnetic phenomena such as de spin Seebeck effect (strikingly in insulators)

Thermomagnetic effects: new means to control heat flow and harvest energy

Need to know: Electron-phonon, **electron-magnon**, **phonon-magnon** ... interactions. Independent determination of the time scales of these interactions

Look at: Resistivity effects, thermoelectric effects: e.g. magnetoresistance, phonon- and magnon-drag

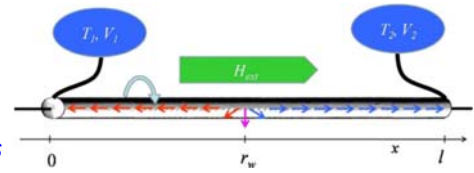
Be careful with: artifacts and mimicking effects, e.g. AMR, phonon drag...



Domain wall motion by magnonic spin currents

Hinzke, Nowak, Phys. Rev. Lett. (2011)

Heat engines and motors



Bauer et al., Phys. Rev. B (2010)

