

# Jan Vogel

## Institut Néel, CNRS Grenoble, France

# European School on Magnetism 2023

# Light – Matter Interaction



## Interaction between photons and magnetism

## About me :

- Ph.D. in Physics in 1994 (University of Nijmegen, the Netherlands) : study of the magnetic properties of thin films using polarized x-rays
- Postdocs in LURE (Orsay), ESRF (Grenoble) and Laboratoire Louis Néel (Grenoble)
- CNRS researcher in LLN -> Institut Néel sinds 1998
  
- Magnetic imaging (PEEM, Kerr), magnetization and domain wall dynamics in magnetic thin films and microstructures

# Outline of the lecture

## 1) Photons to probe magnetism :

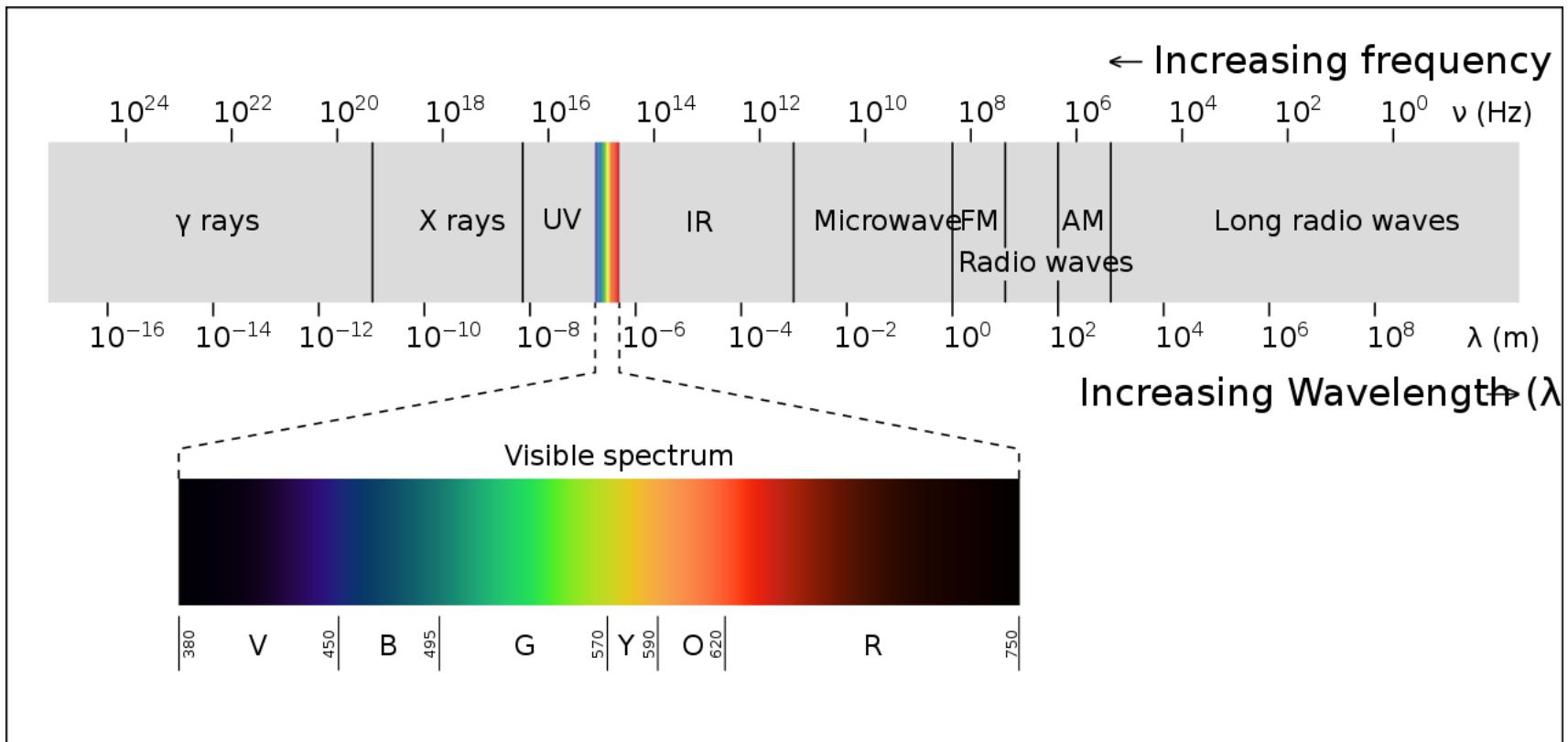
- X-ray magnetic circular and linear dichroism
- Faraday and Kerr effects
- Brillouin Light Scattering

## 2) Influence of light on magnetic properties

- Inverse Faraday effect
- Ultrafast demagnetization using light
- All-optical magnetic switching

# 1) Photons to probe magnetism

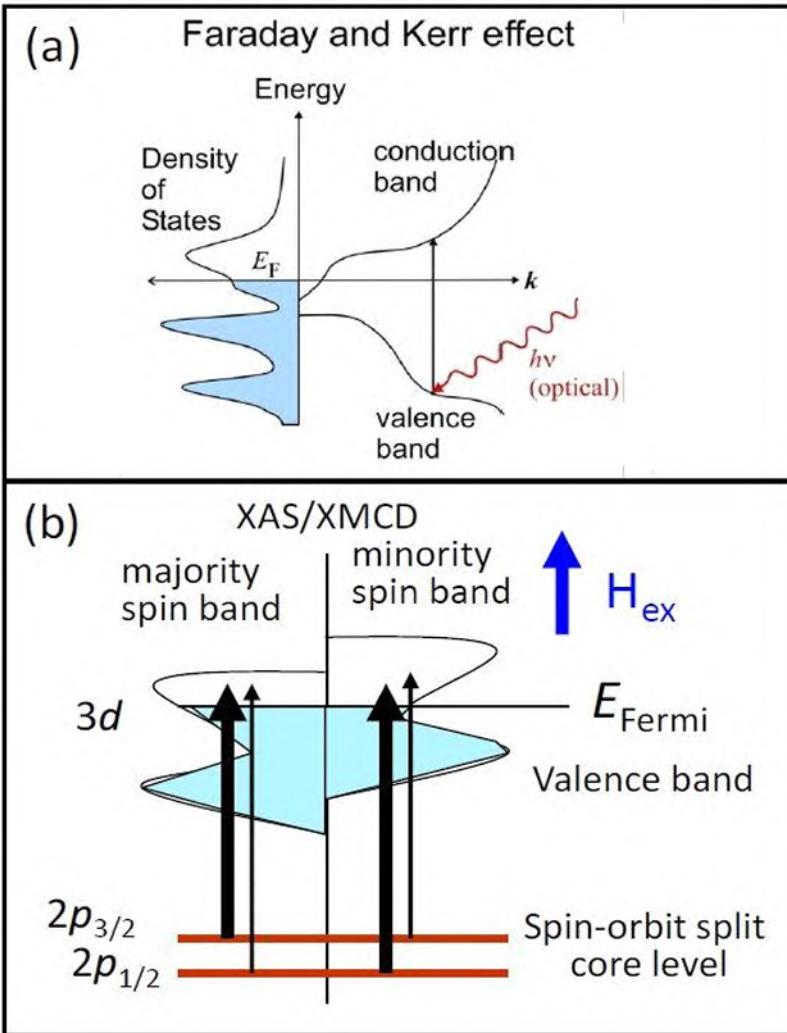
# Optical photons vs X-ray photons



$$E = h\nu = 1.6 - 3.3 \text{ eV for visible light}$$

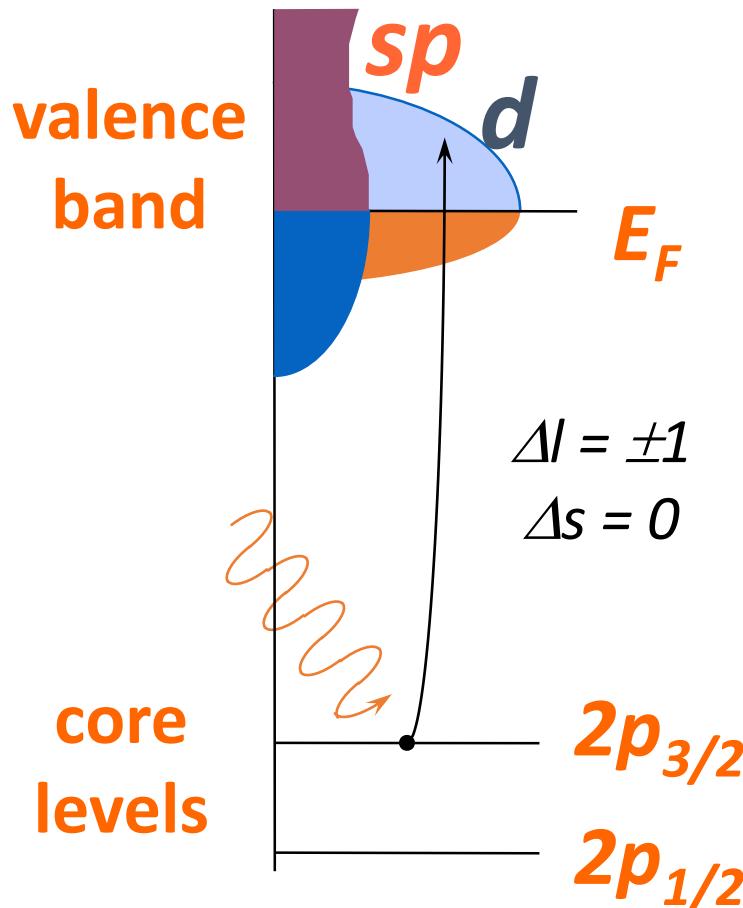
$$E = h\nu = 50 \text{ eV} - 10 \text{ keV for x-rays}$$

# Optical photons vs X-ray photons



Optical photons : transitions between valence band and conduction band → complex initial and final states

X-ray photons : transitions between core level and conduction band → well-defined initial states



To study magnetism :

L<sub>2,3</sub>-edges : 2p → 3d, 4d

M<sub>4,5</sub>-edges : 3d → 4f

**Absorption cross-section :**

$$w_{\text{abs}} = (2\pi/h) | \langle \Phi_f | T | \Phi_i \rangle |^2 \rho_f (E_{\text{hv}} - E_i) \quad \text{Fermi's Golden Rule}$$

# X-ray Absorption

**Periodic Table  
of the Elements**

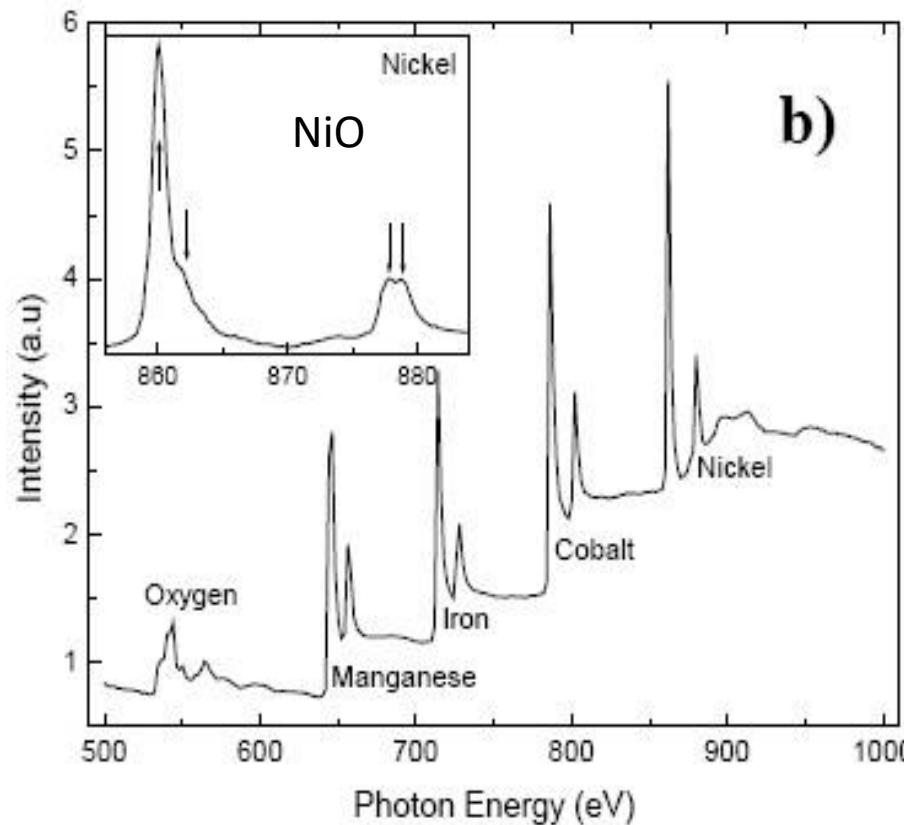
		Periodic Table of the Elements																			
1		IA																		0	
2		H	IIA																	2 He	
3		Li	Be																		
4		Na	Mg																		
5		K	Ca	Sc	Ti	V	Cr	Mn	Fe	Co	Ni	Cu	Zn	Ga	Ge	As	Se	Br	Kr		
6		Rb	Sr	Y	Zr	Nb	Mo	Tc	Ru	Rh	Pd	Ag	Cd	In	Sn	Sb	Te	I	Xe		
7		Cs	Ba	* La	Hf	Ta	W	Re	Os	Ir	Pt	Au	Hg	Tl	Pb	Bi	Po	At	Rn		
		Fr	Ra	+ Ac	Rf	Ha	106	107	108	109	110										
* Lanthanide Series		58 Ce	59 Pr	60 Nd	61 Pm	62 Sm	63 Eu	64 Gd	65 Tb	66 Dy	67 Ho	68 Er	69 Tm	70 Yb	71 Lu						
+ Actinide Series		90 Th	91 Pa	92 U	93 Np	94 Pu	95 Am	96 Cm	97 Bk	98 Cf	99 Es	100 Fm	101 Md	102 No	103 Lr						

Main magnetic elements:

- Transition Metals (Mn, Fe, Co, Ni but also Pd, Pt,...) 2p → 3d (4d, 5d)
- Rare Earths (Nd, Gd, Tb, Dy, ...) 3d → 4f (also 2p → 5d)

# X-ray Absorption

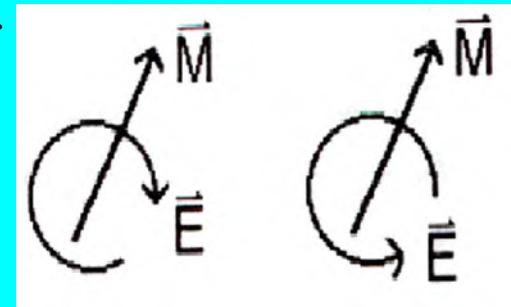
3d-transition metal L<sub>2,3</sub>-edges



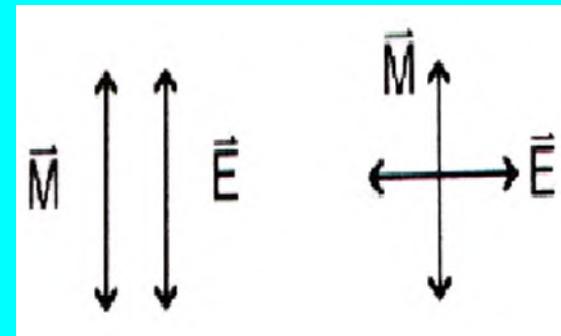
Element selective, information on electronic state

# Magnetic sensitivity

**X-ray Magnetic Circular Dichroism (XMCD):** difference in absorption for left- and right circularly polarized light.

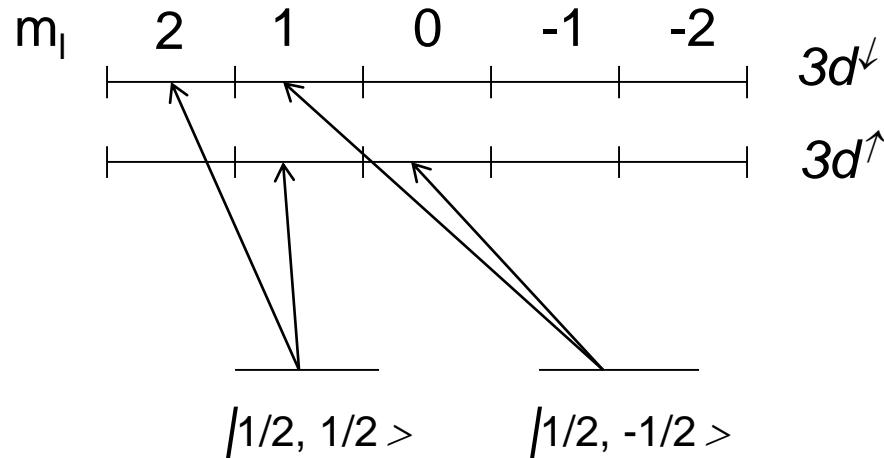


**X-ray Linear Dichroism:** difference in absorption for linearly polarized light  $\perp$  and  $\parallel$  to quantization axis.



# XMCD : origin

One electron picture: transitions from  $2p$  to  $3d$  band, split by exchange in  $3d^\uparrow$  and  $3d^\downarrow$



$L_2$  edge - left polarisation (  $\Delta m_l = +1$  )

$$I^\uparrow = \sum | \langle f | P_1 | i \rangle |^2 = (1/3 | \langle 2,1 | P_1 | 1,0 \rangle |^2 + 2/3 | \langle 2,0 | P_1 | 1,-1 \rangle |^2) R^2$$

$$R = \int R_{nl}^*(r) R_{n'l'}(r) r^3 dr$$

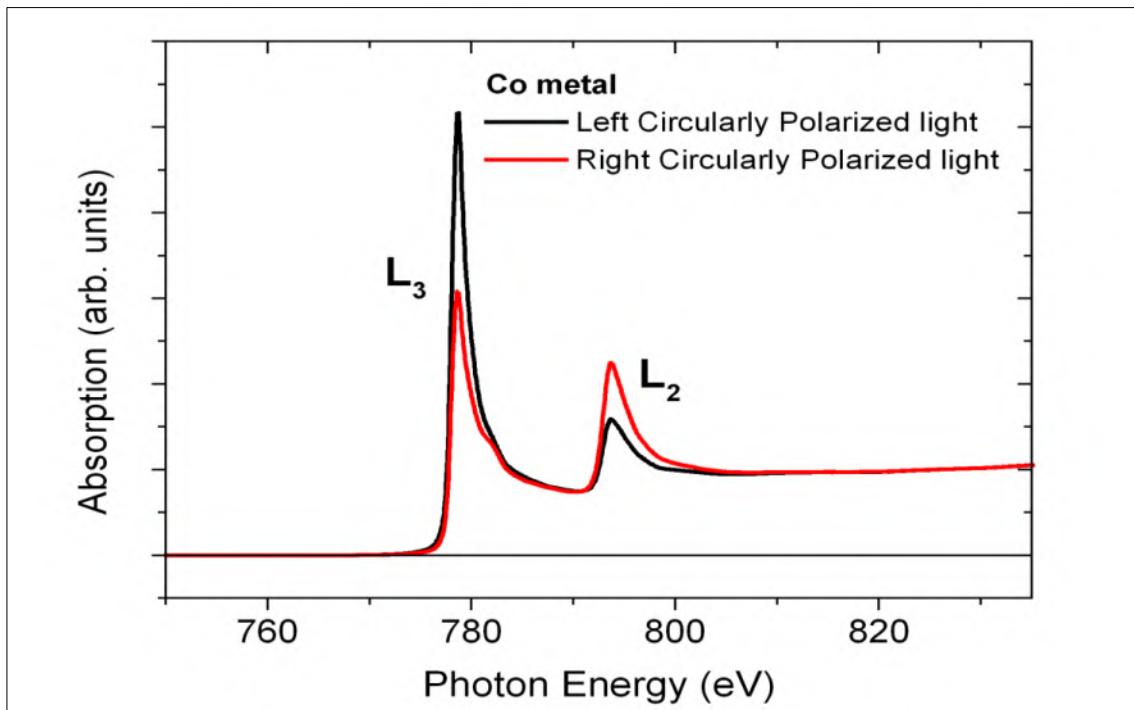
$$I^\downarrow = \sum_{i,f} | \langle f | P_1 | i \rangle |^2 = (2/3 | \langle 2,2 | P_1 | 1,1 \rangle |^2 + 1/3 | \langle 2,1 | P_1 | 1,0 \rangle |^2) R^2$$

J.L. Erskine and E.A.Stern, Phys.Rev.B 12, 5016 (1975).

# XMCD : origin

**Step 1** : spin-polarised electrons emitted by the spin-orbit split 2p band  
75% spin down and 25% spin up electrons at the  $L_2$ -edge with LCP light  
37.5% spin down and 62.5% spin up electrons at the  $L_3$ -edge with LCP light

**Step 2**: the exchange split  $d$ -band acts as spin-detector.



Sum rules relate dichroism and total absorption to the ground-state orbital and spin magnetic moment of the probed element and shell:

L<sub>2,3</sub>-edges of Fe → Fe 3d-moments.

## Orbital moment sum rule:

$$\langle L_Z \rangle = [2l(l+1)(4l+2-n)]/[l(l+1)+2 - c(c+1)] \bullet$$

$$[\int_{j_+ + j_-} d\omega (\mu^+ - \mu^-) / \int_{j_+ + j_-} d\omega (\mu^+ + \mu^- + \mu^0)]$$

***l*** = orbital quantum number of the valence state

***c*** = orbital quantum number of the core state

***n*** = number of electrons in the valence state

$\mu^+$  ( $\mu^-$ ) = absorption spectrum for left (right) circularly polarized light.

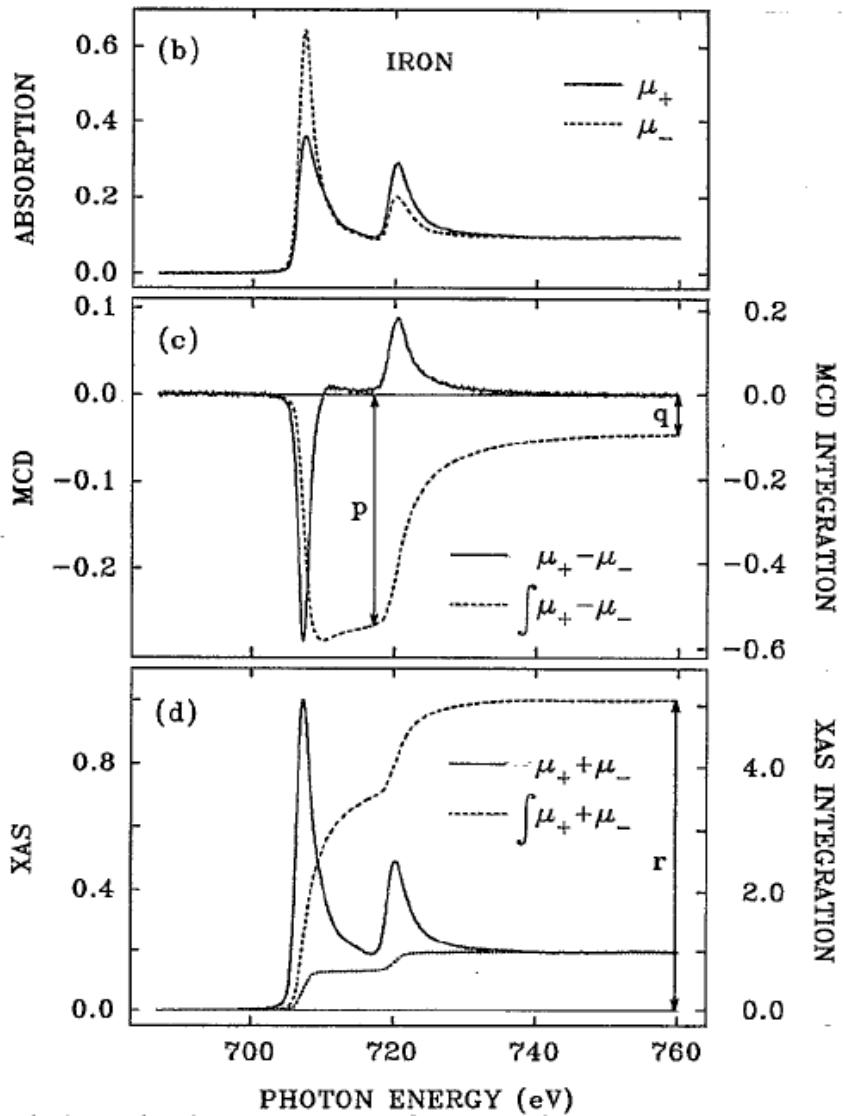
$\mu^0$  = absorption spectrum for linearly polarized light, with polarization parallel to quantization axis.

$j^+ (j^-)$  =  $(l + 1/2)$  resp.  $(l - 1/2)$  absorption (ex. 2p<sub>3/2</sub>, 2p<sub>1/2</sub>)

B.T.Thole *et al.*, Phys.Rev.Lett. 68, 1943 (1992)

M.Altarelli, Phys.Rev.B 47, 597 (1993)

# Sum rule for orbital moment



For L<sub>2,3</sub>-edges  $c = 1$  ( $2p$ ),  $l = 2$  ( $d$ ):

$$\langle L_Z \rangle = 2(10-n) \bullet (\Delta L_3 + \Delta L_2)$$

$$/ \int_{L_3+L_2} d\omega (\mu^+ + \mu^- + \mu^0)]$$

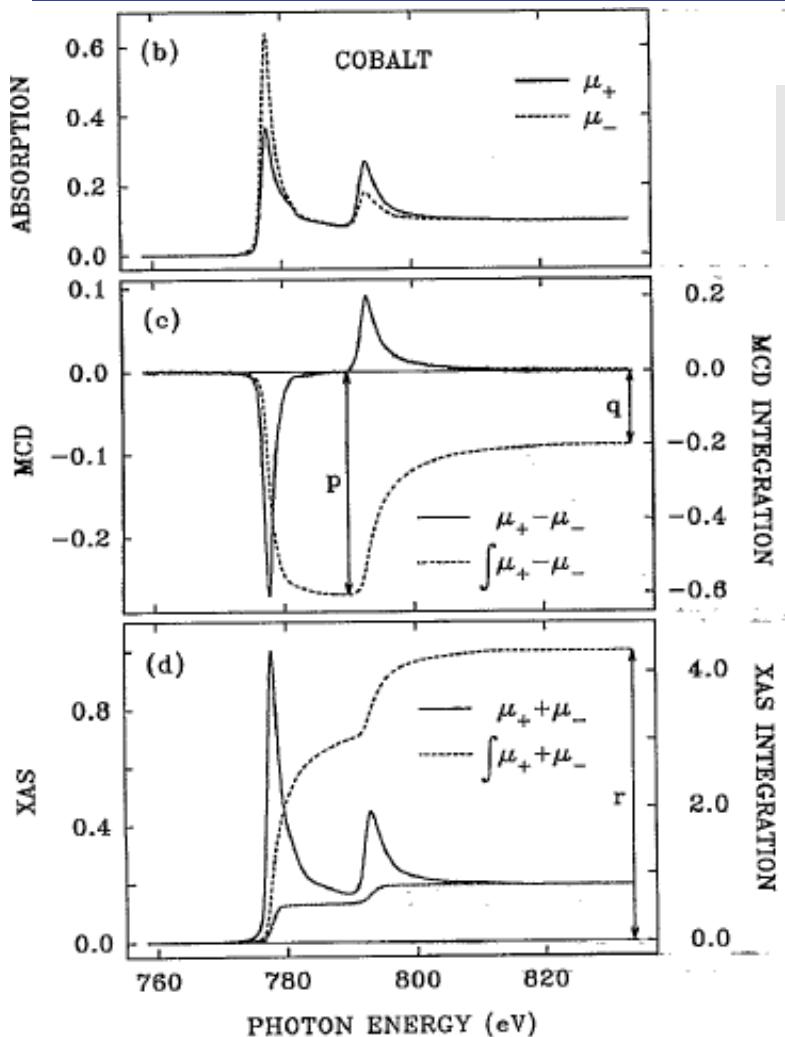
$$q = \Delta L_3 + \Delta L_2$$

$$r = \mu^+ + \mu^- = (2/3)(\mu^+ + \mu^- + \mu^0)$$

$$\langle L_Z \rangle = 4q (10-n) / 3r$$

C.T.Chen *et al.*, PRL 75, 152 (1995)

# Sum rule for spin moment



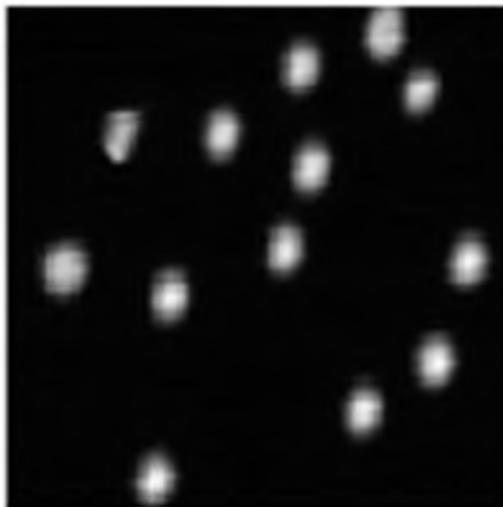
$$\langle S_z \rangle + (7/2) \langle T_z \rangle = \\ (3/2)(10-n)[(\Delta L_3 - 2\Delta L_2)/ \int_{L_3+L_2} d\omega (\mu_+ + \mu_- + \mu_0)]$$

$$= (3/2)(10-n)(p - 2(q-p))/(3/2)r = \\ = (3p - 2q)(10-n)/r$$

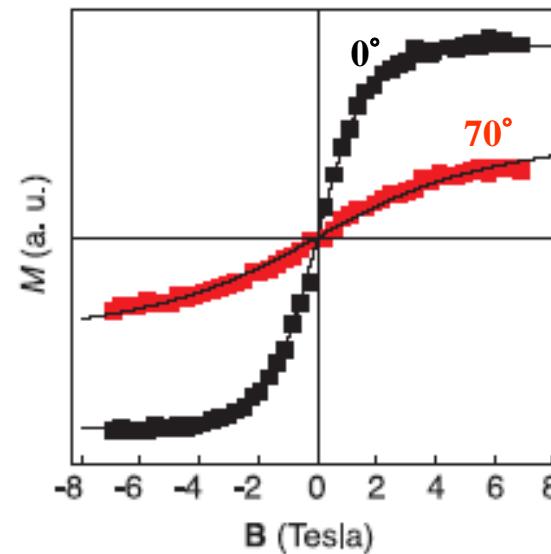
C.T.Chen *et al.*, PRL 75, 152 (1995)

**Single Co adatoms and particles MBE deposited on Pt(111) surfaces**P. Gambardella *et al.*, Science 300, 1130 (2003)

STM image of isolated Co adatoms  
(8.5 nm x 8.5 nm)

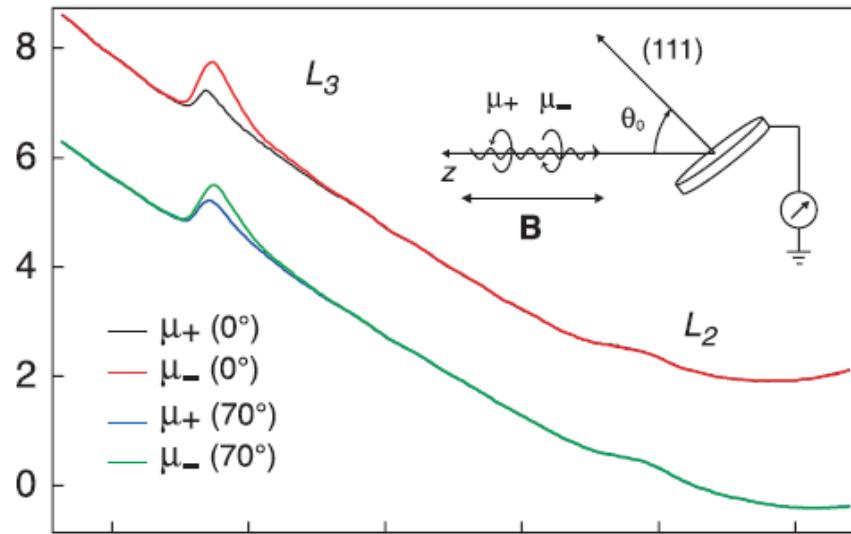


Large difference of in-plane and out-of-plane saturation field : very large Magnetic Anisotropy Energy

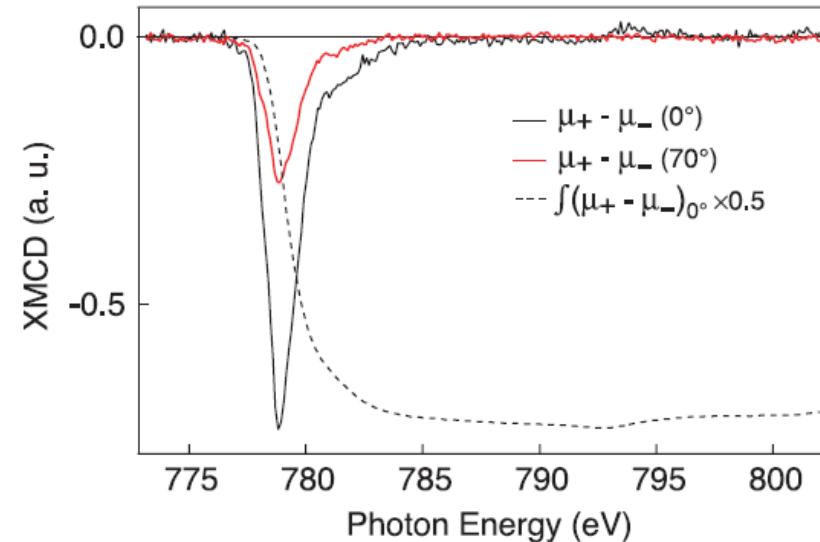


## Single Co adatoms and particles MBE deposited on Pt(111) surfaces

P. Gambardella *et al.*, Science 300, 1130 (2003)



Co L<sub>2,3</sub> X-ray absorption spectra for  $\theta = 0^\circ$  and  $70^\circ$



vanishing L<sub>2</sub> XMCD: very large orbital magnetism

**Single Co adatoms and particles MBE deposited on Pt(111) surfaces**

P. Gambardella *et al.*, Science 300, 1130 (2003)

**Sum rules :**  $\langle L \rangle = 1.1 \pm 0.1 \mu_B$  for isolated Co adatoms ( $L = 0.15 \mu_B$  Co-hcp)  
 $(L = 0.29 \mu_B$  1ML Co/Pt)

Reduced coordination of isolated atoms on top of a flat surface → d-electron localisation, increase of atomic character

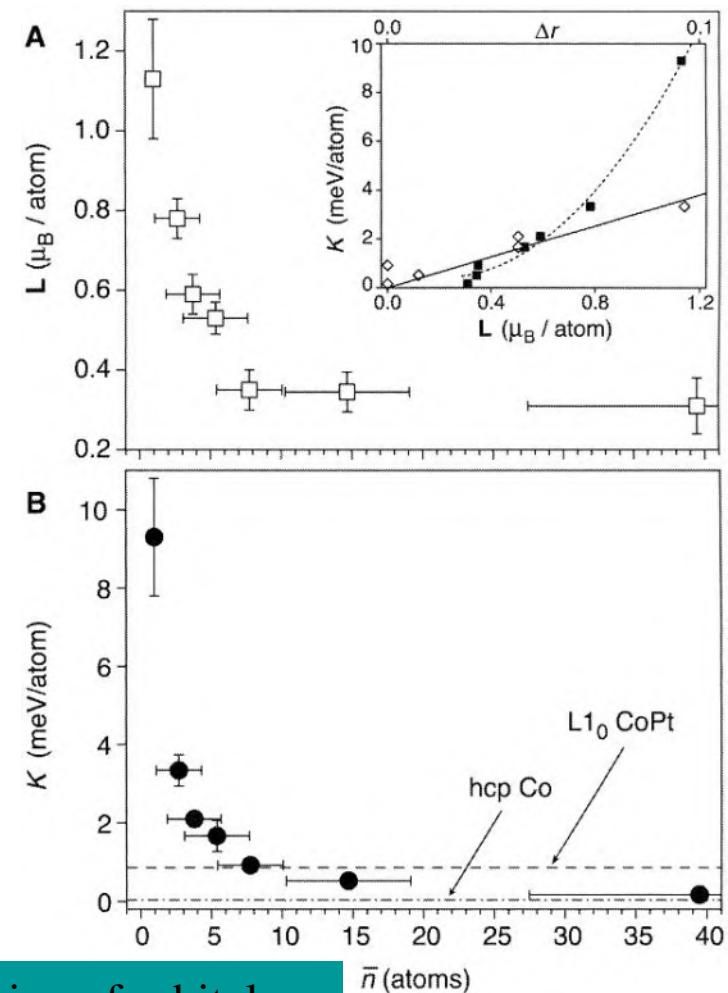
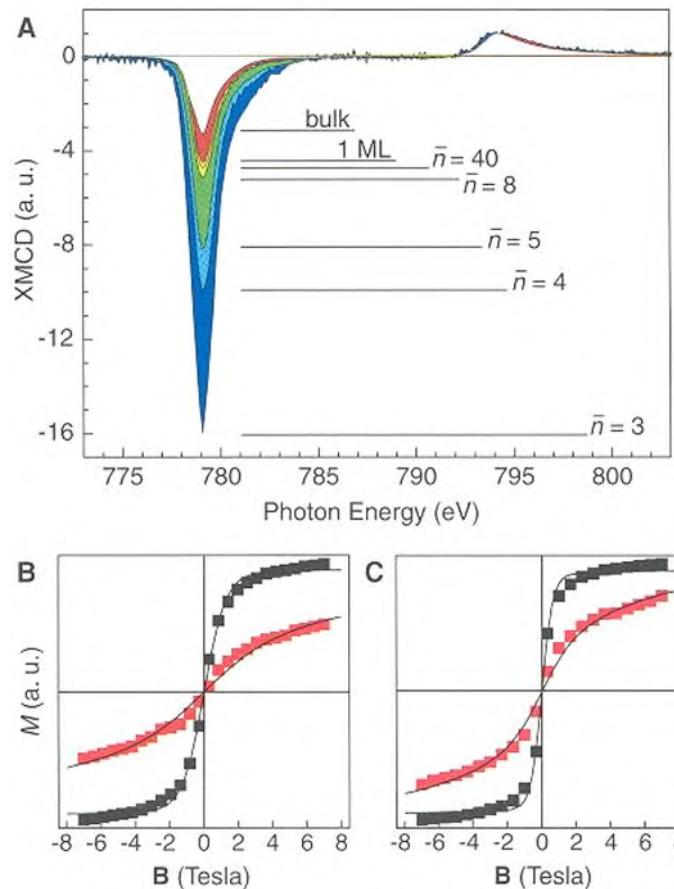
From element-selective XMCD magnetization curves (up to 7 Tesla):  
very large magnetic anisotropy energy (MAE)

$K = 9.3 \pm 1.6 \text{ meV/atom}$

( $K = 1.8 \text{ meV/Co atom in SmCo}_5$ )  
( $K = 0.3 \text{ meV/atom in Pt/Co multilayers}$ )

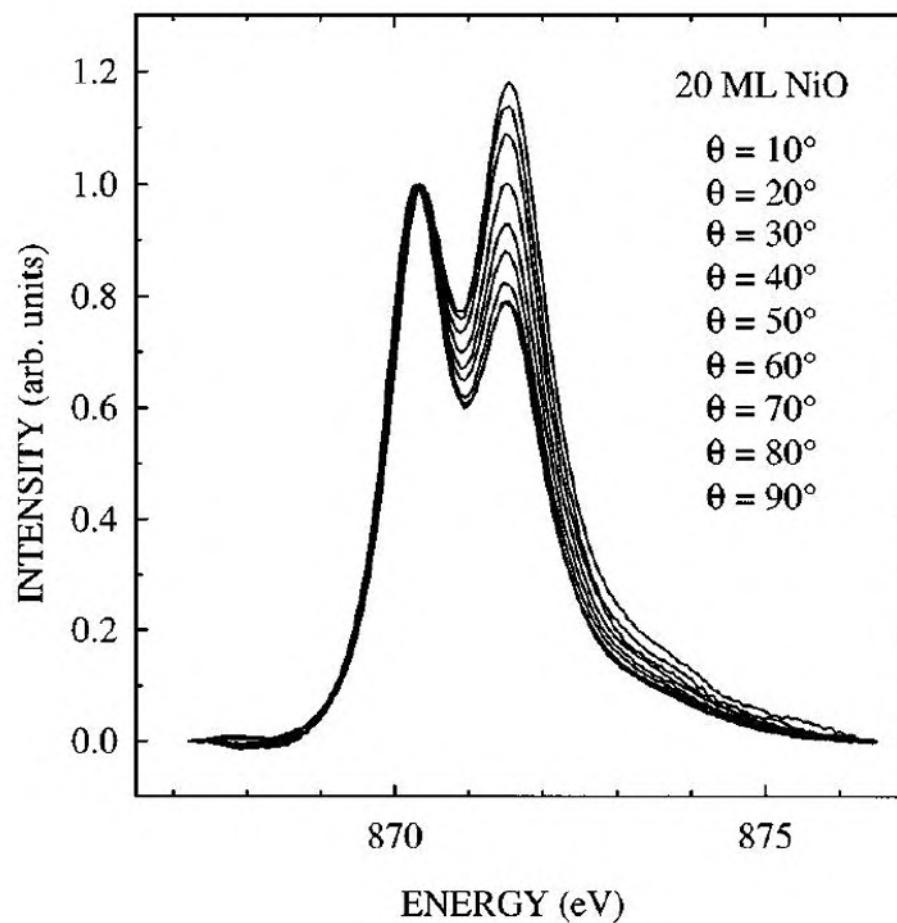
# XMCD sum rules : Orbital magnetic moments and anisotropy

P. Gambardella *et al.*, Science 300, 1130 (2003)



Increase of particle size : progressive quenching of orbital moment and consequent decrease of MAE

## X-ray Linear Dichroism



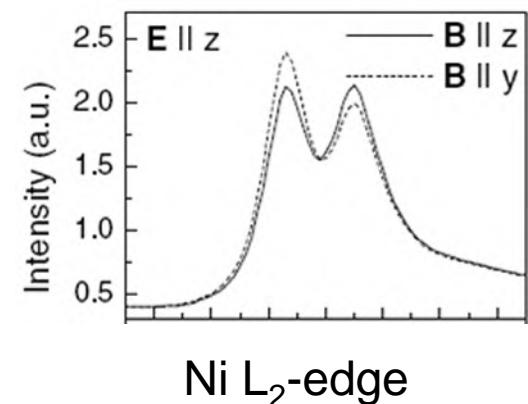
Dependence of Ni L2 spectrum in NiO on the incidence angle of linearly polarized light  
D. Alders et al., Phys. Rev. B 57, 11623 (1998).

# X-ray Linear Dichroism

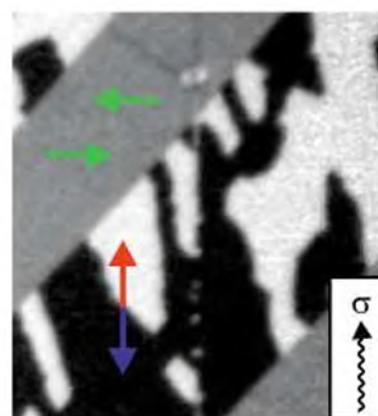
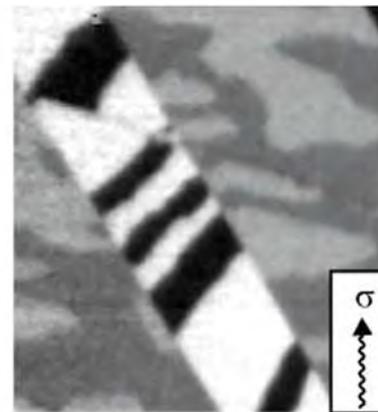
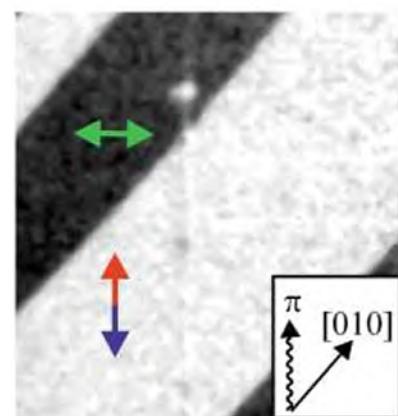
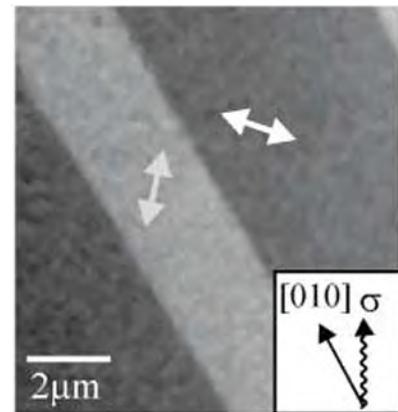
Ni XMCD

Co XMCD

8 ML Co/NiO(001)

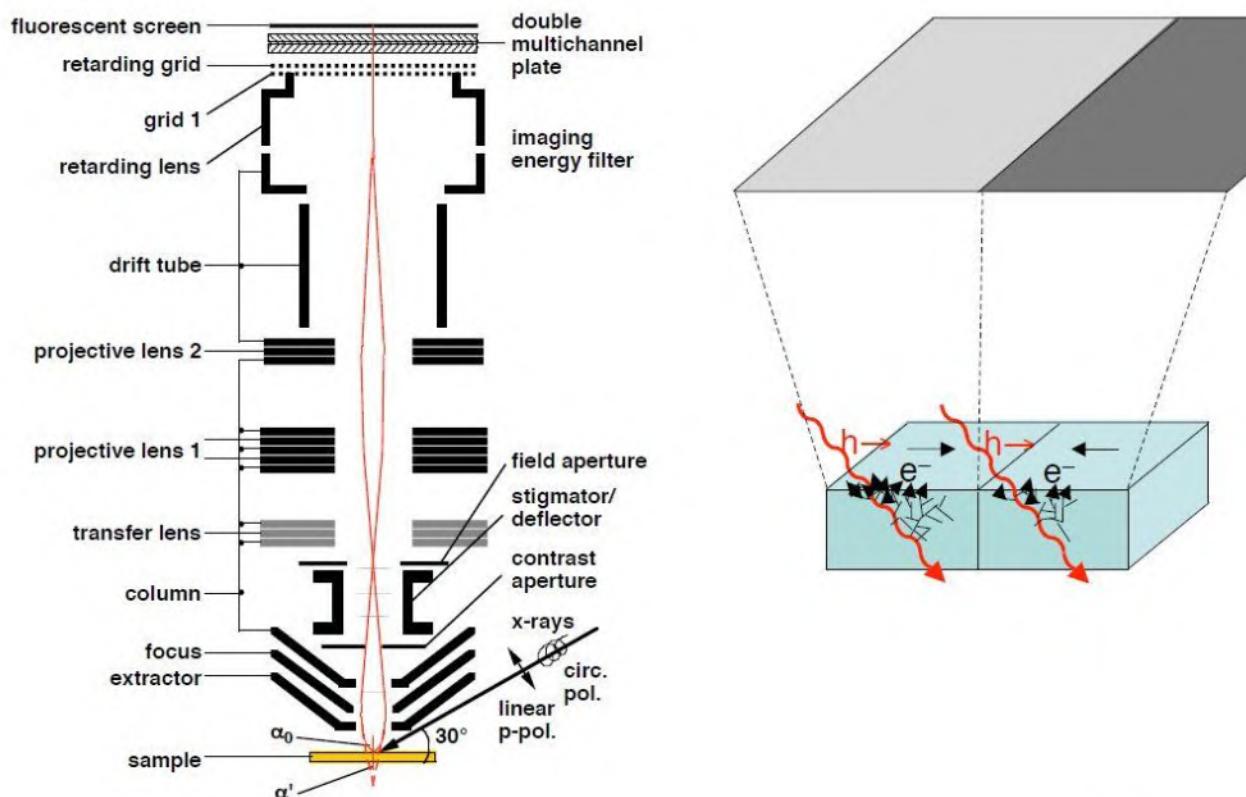


Ni  $L_2$ -edge



H. Ohldag et al., PRL 86 (2001) 2878

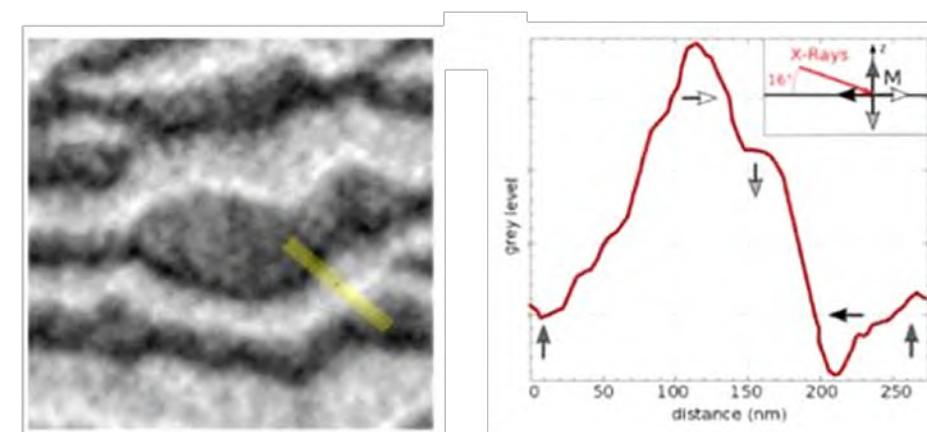
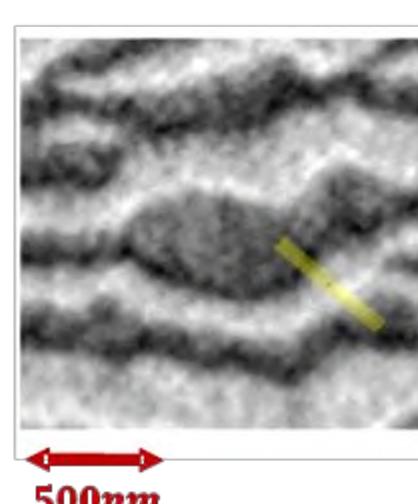
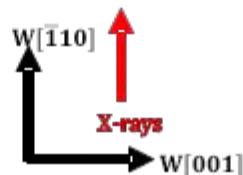
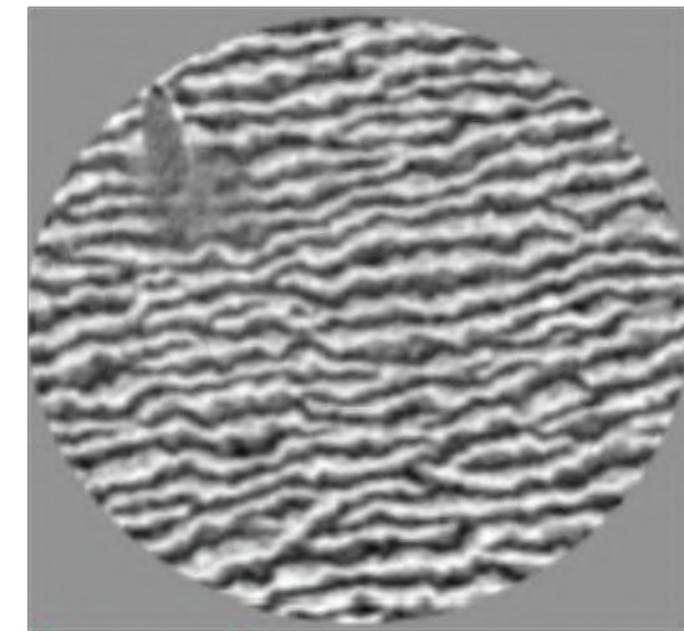
# Magnetic Imaging using X-rays



© W. Kuch 2009

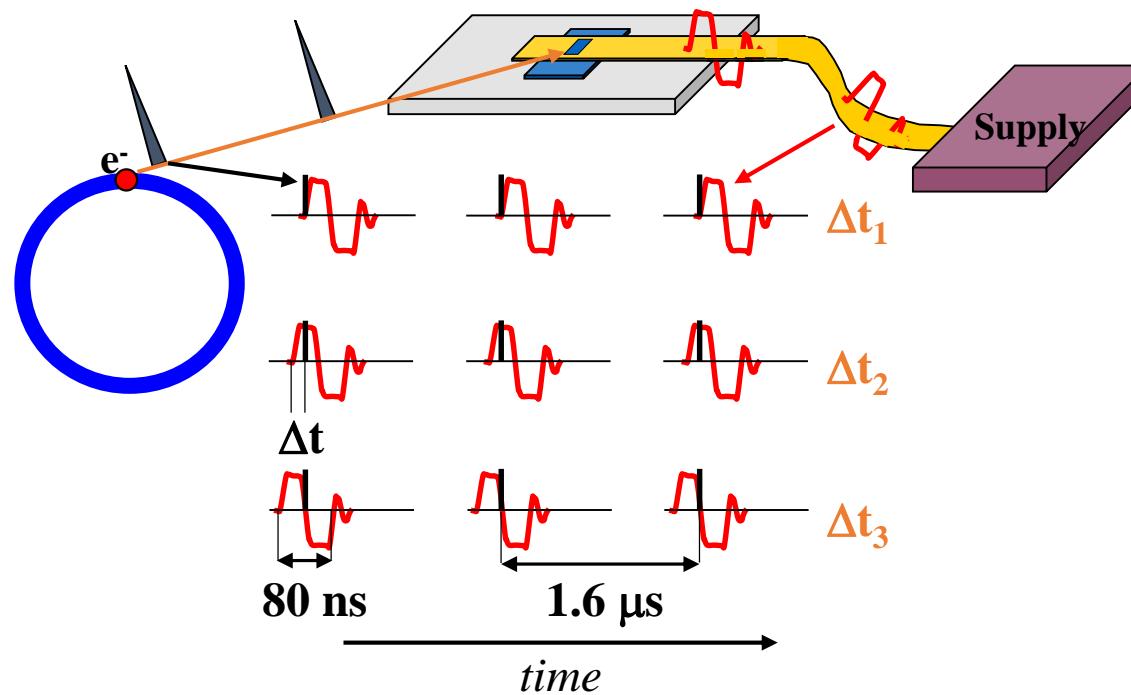
PhotoEmission Electron Microscope + XMCD

## XMCD-PEEM: 4 ML Co



L. Camosi, J. Peña Garcia et al., New. J. Phys. 23, 013020 (2021)

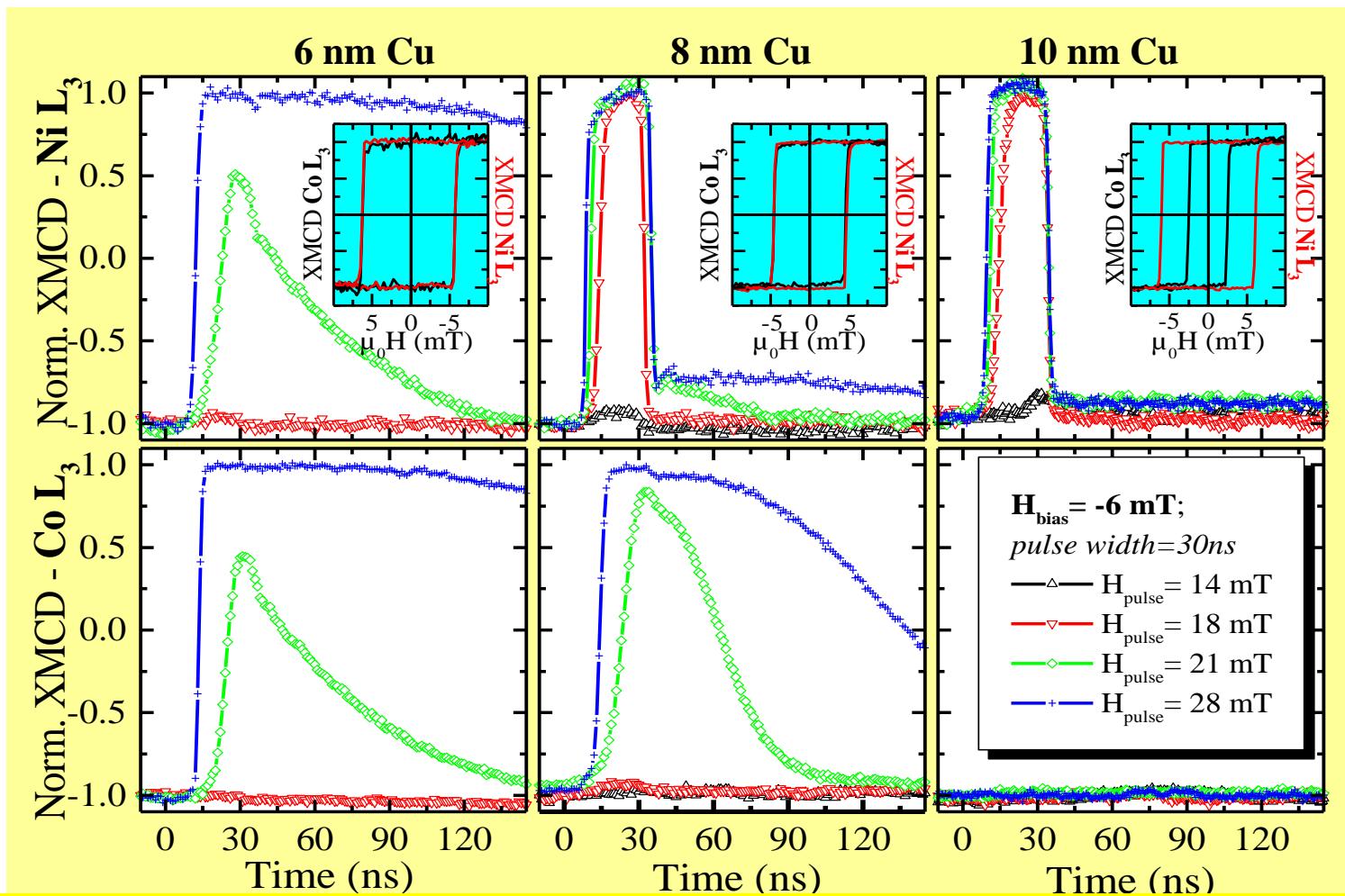
## Pump-probe measurements with pulsed x-ray bunches



M. Bonfim et al., Phys. Rev. Lett. 86, 3646 (2001); J. Vogel et al., Appl. Phys. Lett. 82, 2299 (2003)

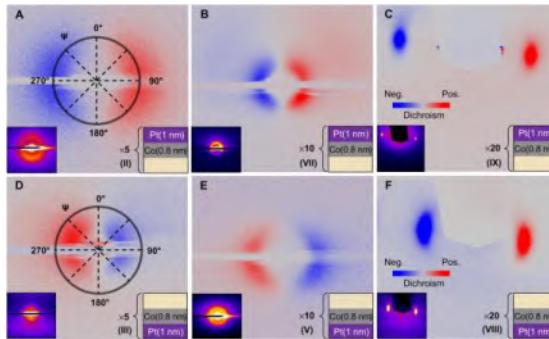
# Time-resolved XMCD

FeNi/Cu/Co



The coupling between the soft and hard layers in the spin-valve systems is different in the static and dynamic regimes

# Other x-ray techniques

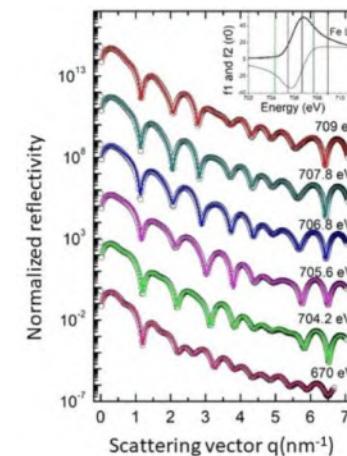


X-ray resonant magnetic scattering :  
periodic magnetic structures, chirality

W. Legrand et al., Sci. Adv. 2018;4:eaat041

X-ray resonant magnetic reflectivity :  
depth resolved magnetization profile

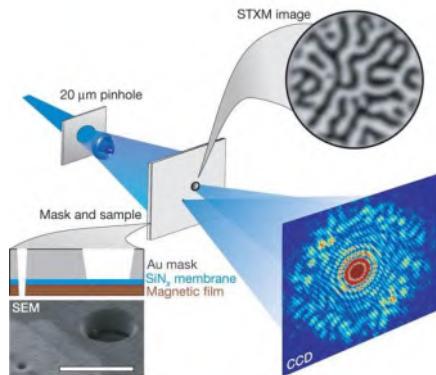
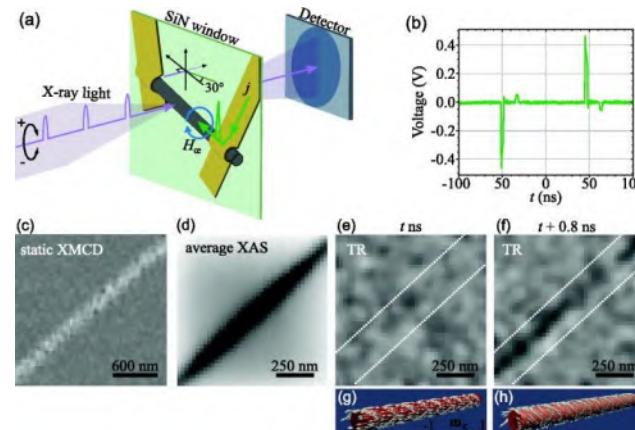
J.-M. Tonnerre et al., AIP Adv. 12,  
035129 (2022).



# Other x-ray techniques

## Scanning Transmission X-ray Microscopy

O. Fruchart et al., APL 2022

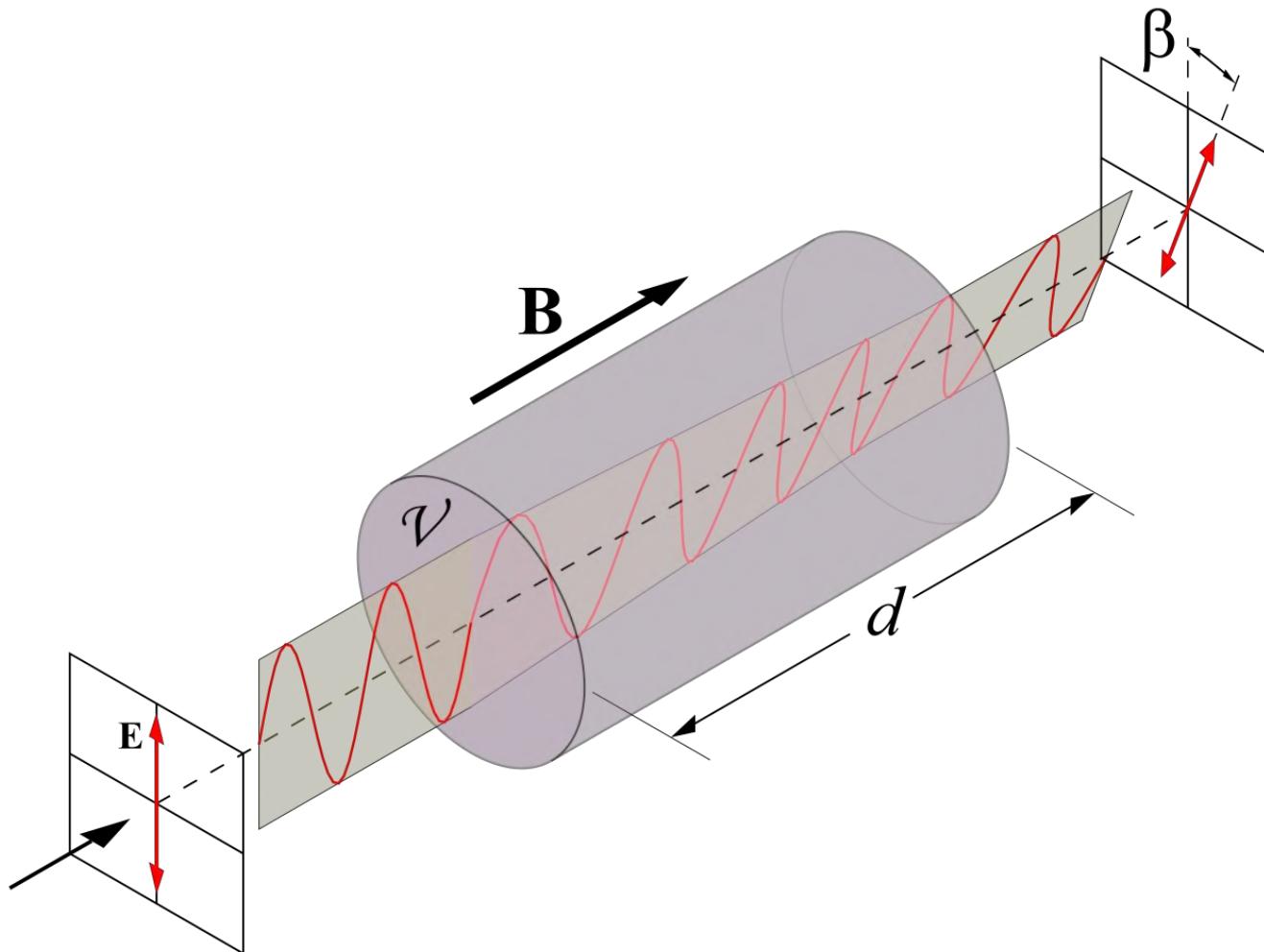


## Magnetic X-ray Holography

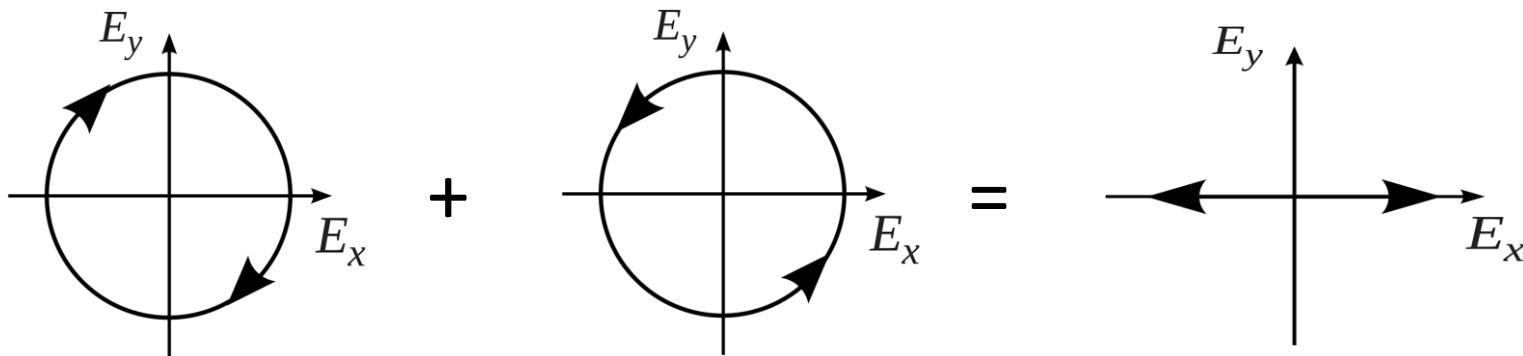
S.Eisebitt et al., Nature **432**, 885 (2004)

# Magneto-optics

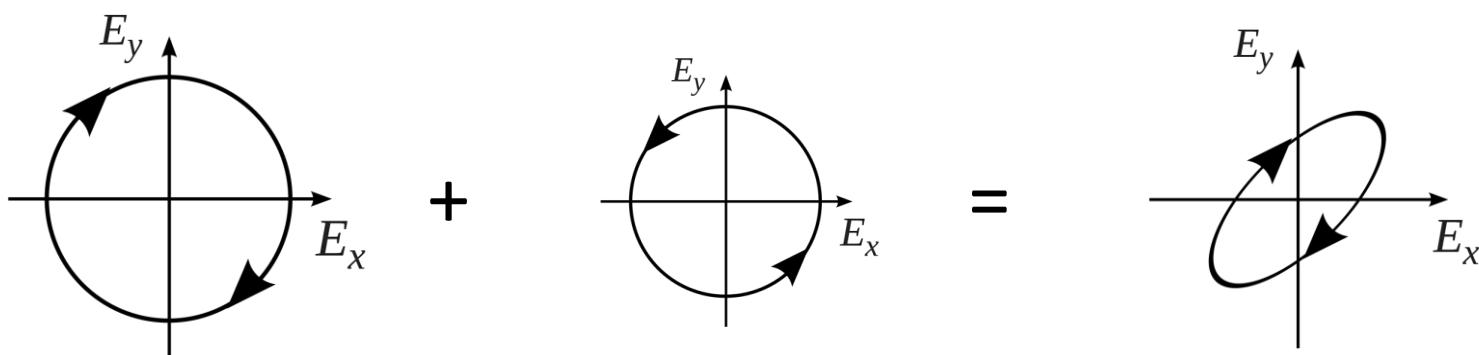
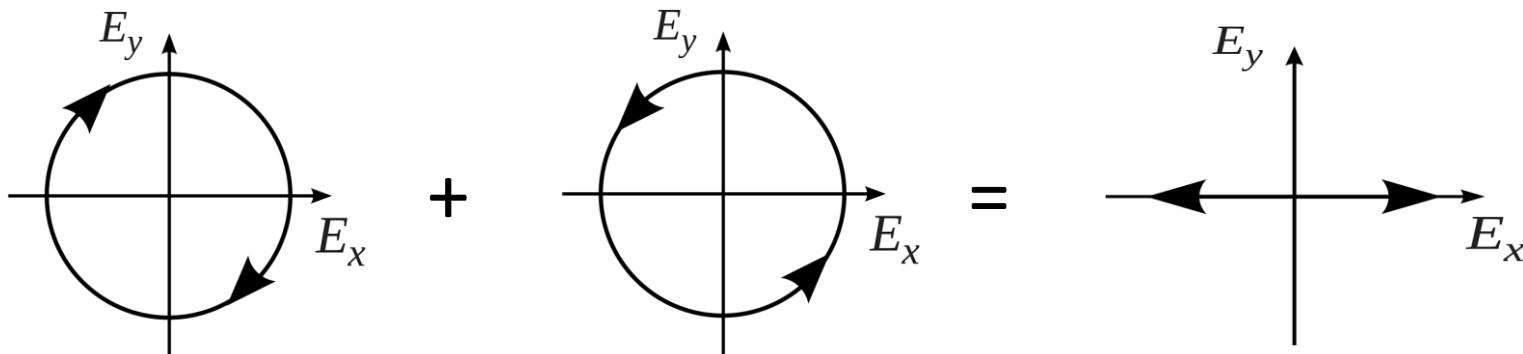
# Faraday effect



# Light polarization

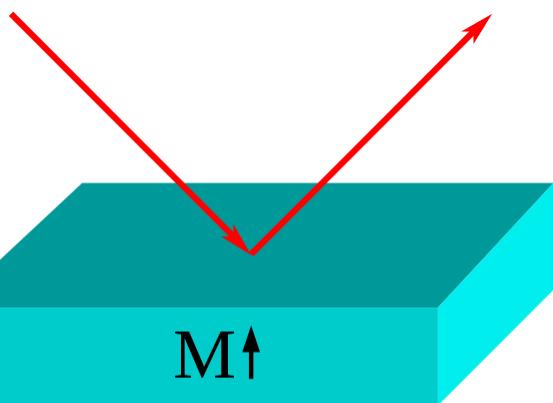


# Light polarization

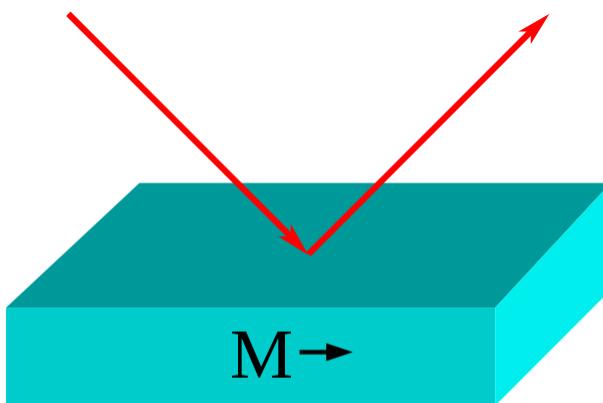


Different absorption left- and right circularly polarized light → ellipticity  
+ rotation polarization axis if different phases

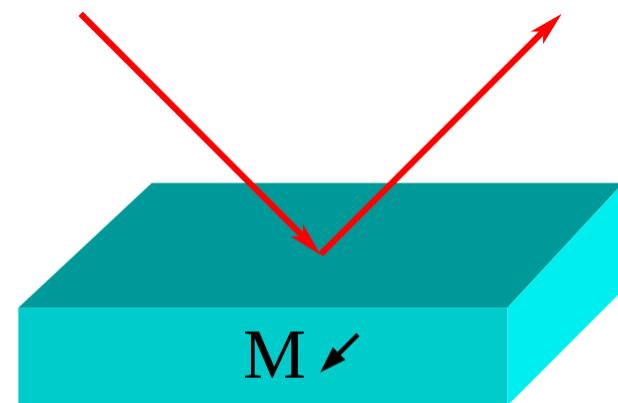
# Magneto-optical Kerr effect (MOKE)



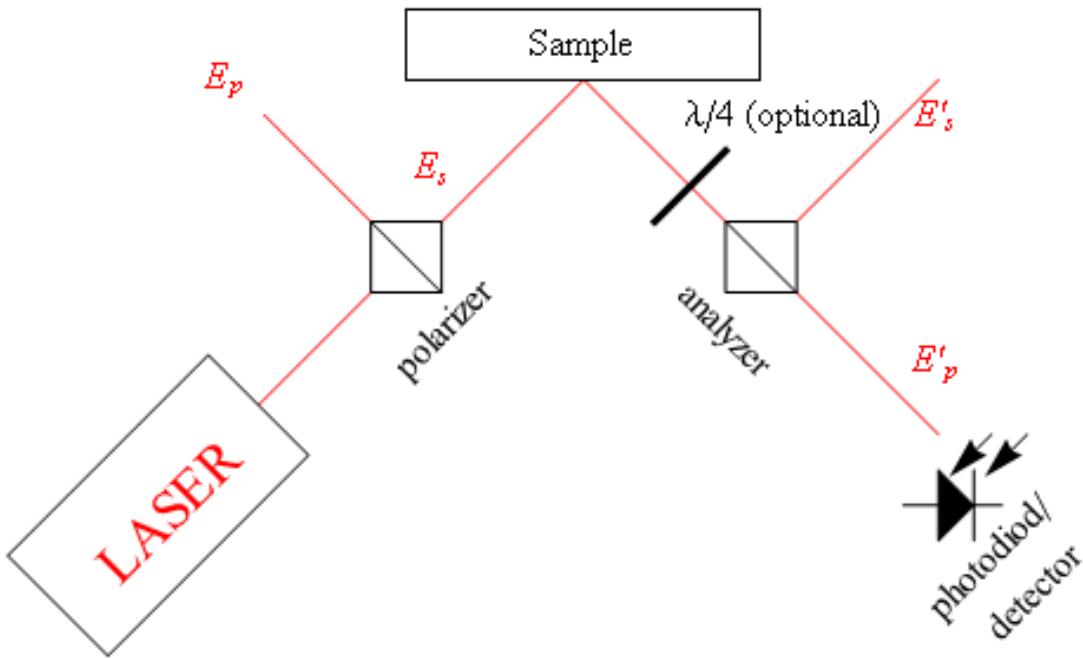
Polar



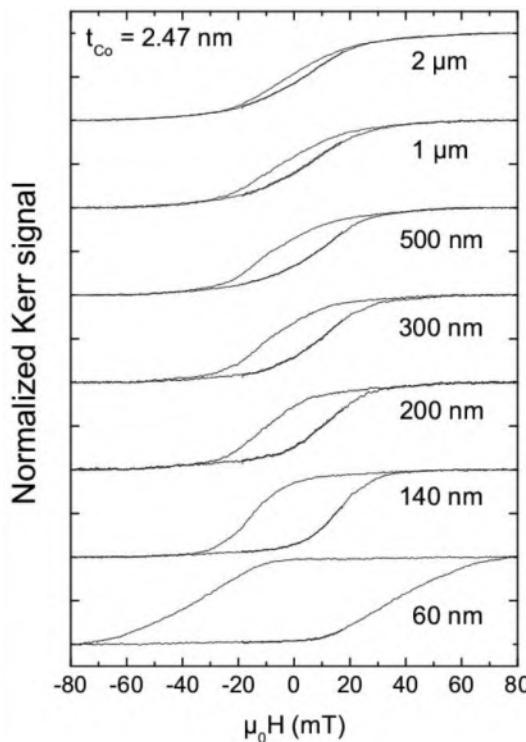
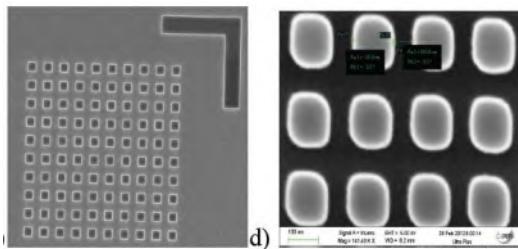
Longitudinal



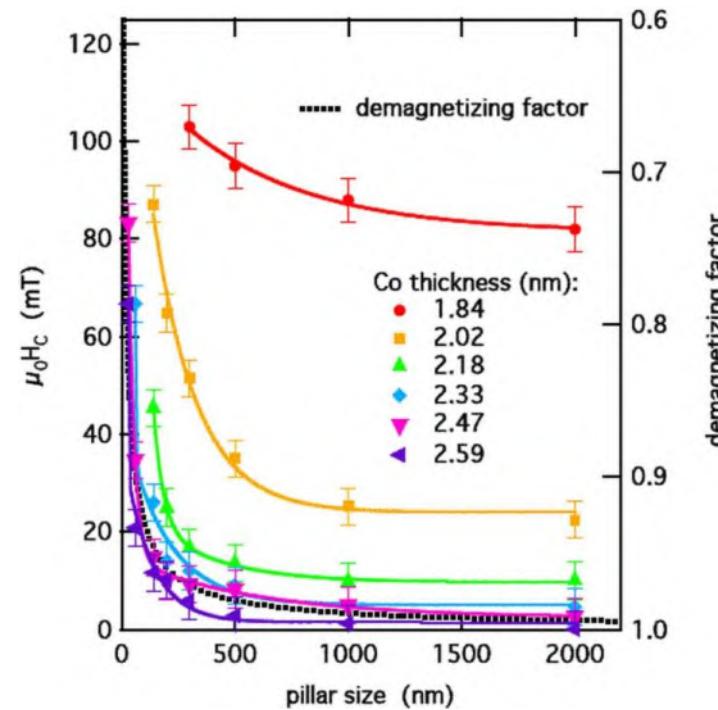
Transversal MOKE



# Focused Kerr



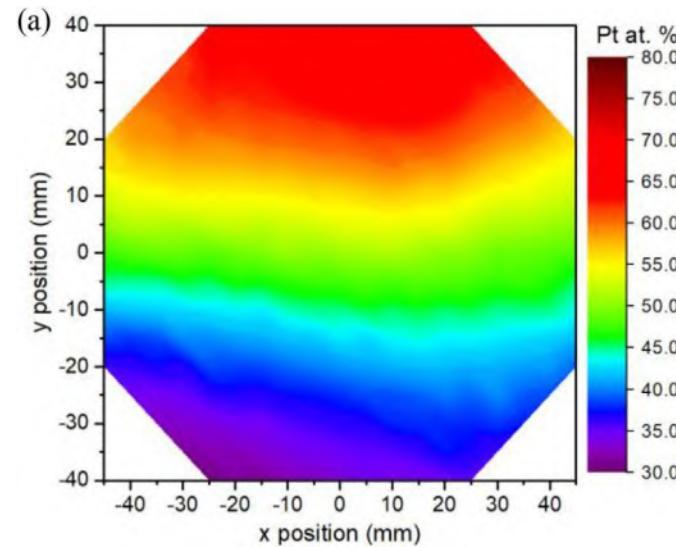
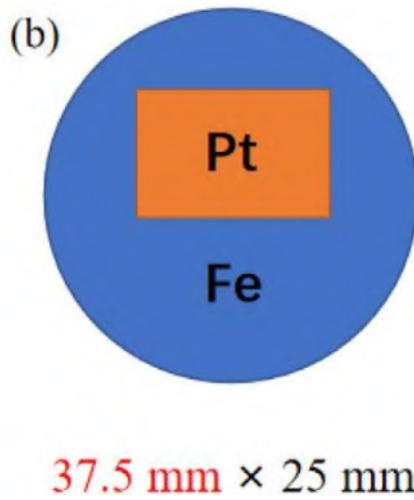
Pt/Co/MgO pillars with different size and Co thickness



W. Stefanowicz et al., Appl. Phys. Lett. 104, 012404 (2014)

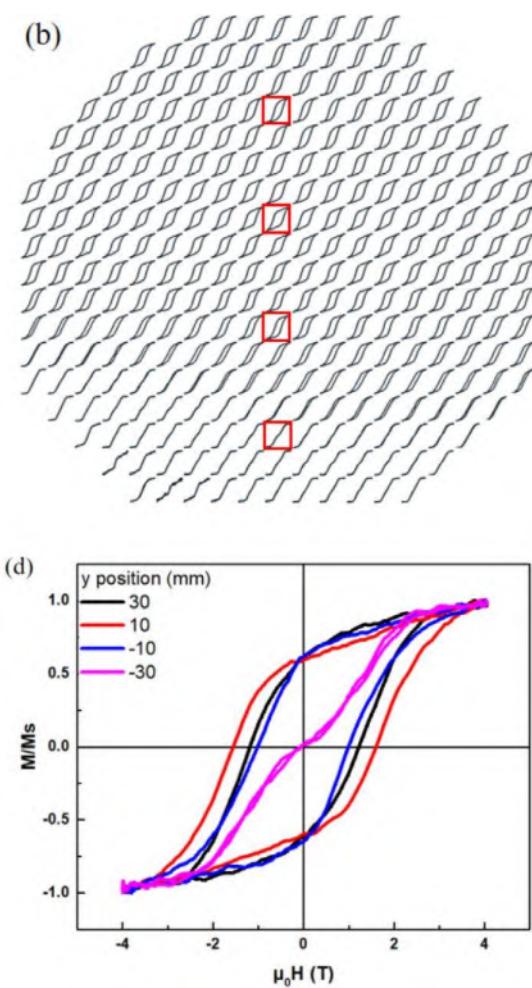
# Scanning Kerr

A high throughput study of compositionally graded FePt thin films using scanning MOKE

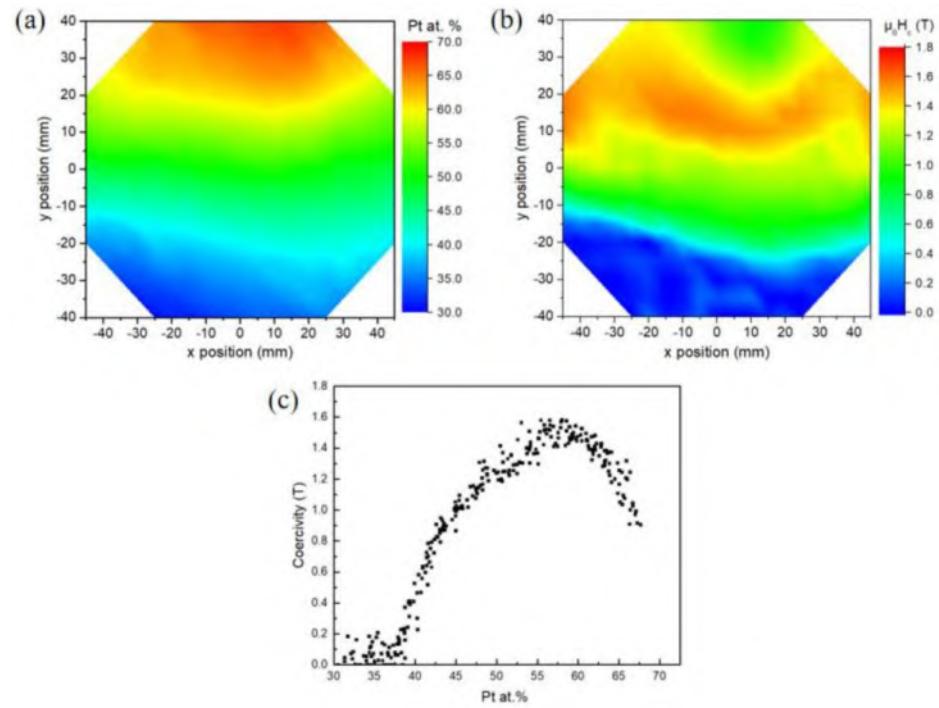


Y. Hong et al., J. Mater. Res. and Techn. 18, 1245 (2022).

# Scanning Kerr



Automatic scanning over 100mm diameter wafer with field pulses up to 3T

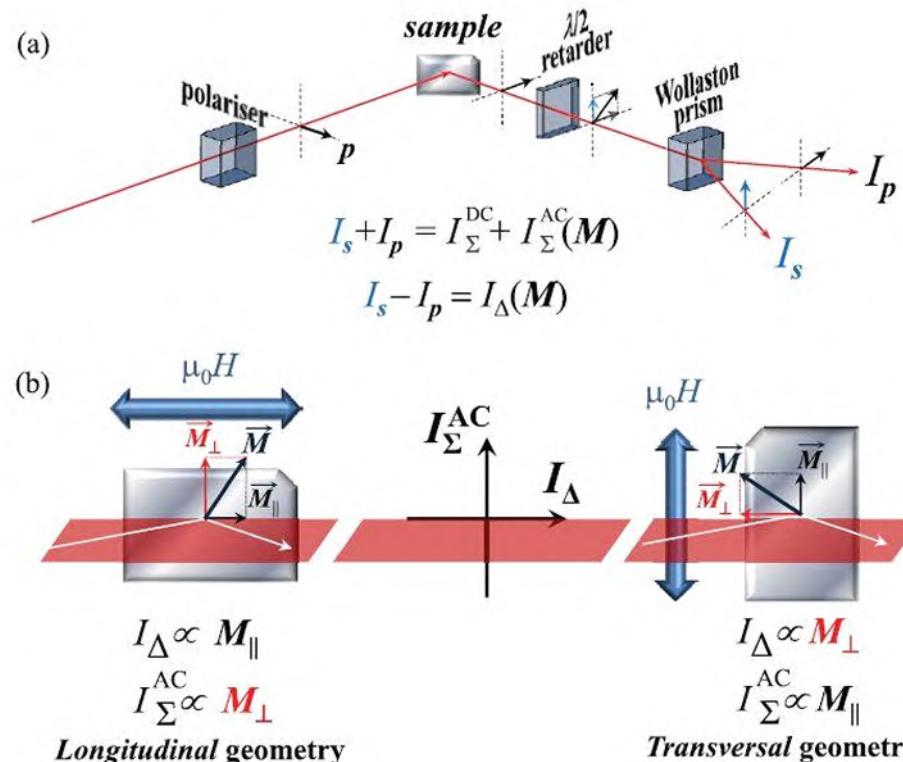


Y. Hong et al., J. Mater. Res. and Techn. 18, 1245 (2022).

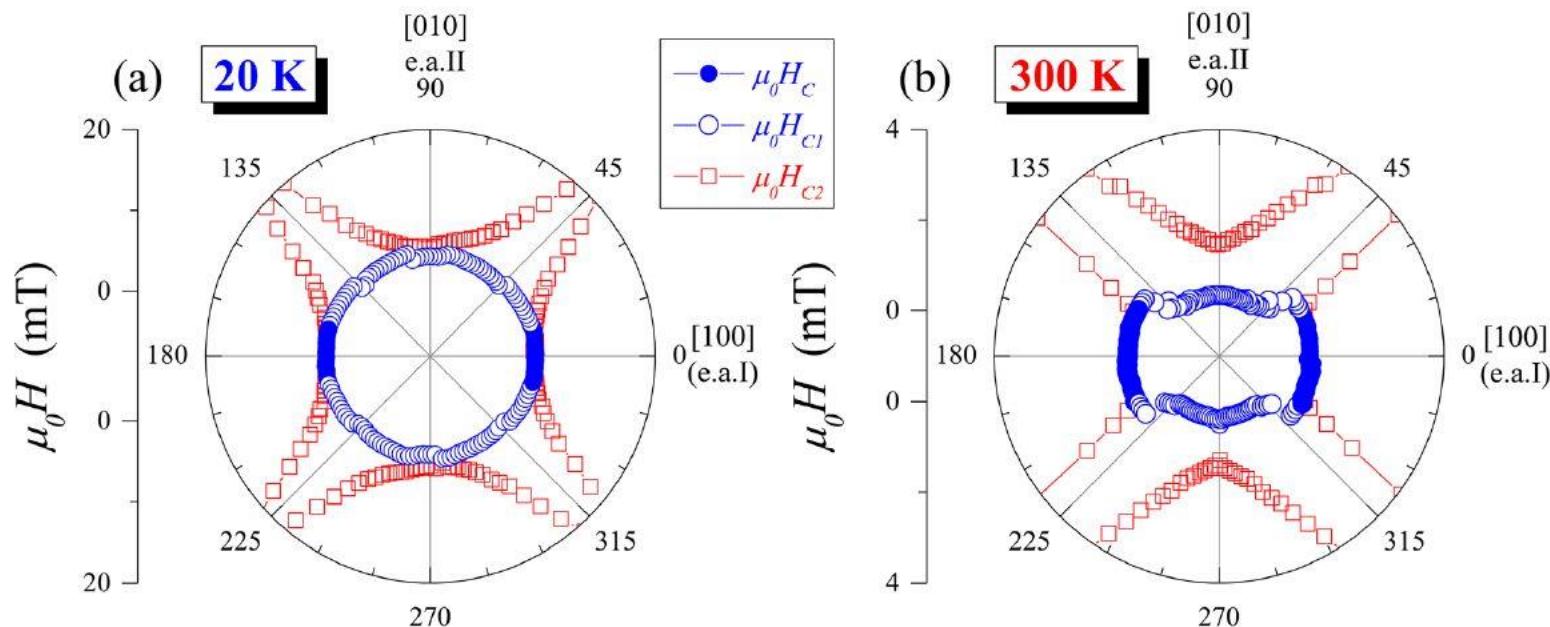
# Vectoriel Kerr

Measuring at the same time longitudinal and transverse magneto-optical Kerr effect : vectoriel resolved MOKE magnetometry

Jose Luis F Cuñado, Javier Pedrosa, Fernando Ajejas, Paolo Perna, Rodolfo Miranda  
and Julio Camarero, J. Phys.: Condens. Matter **29** (2017) 405805



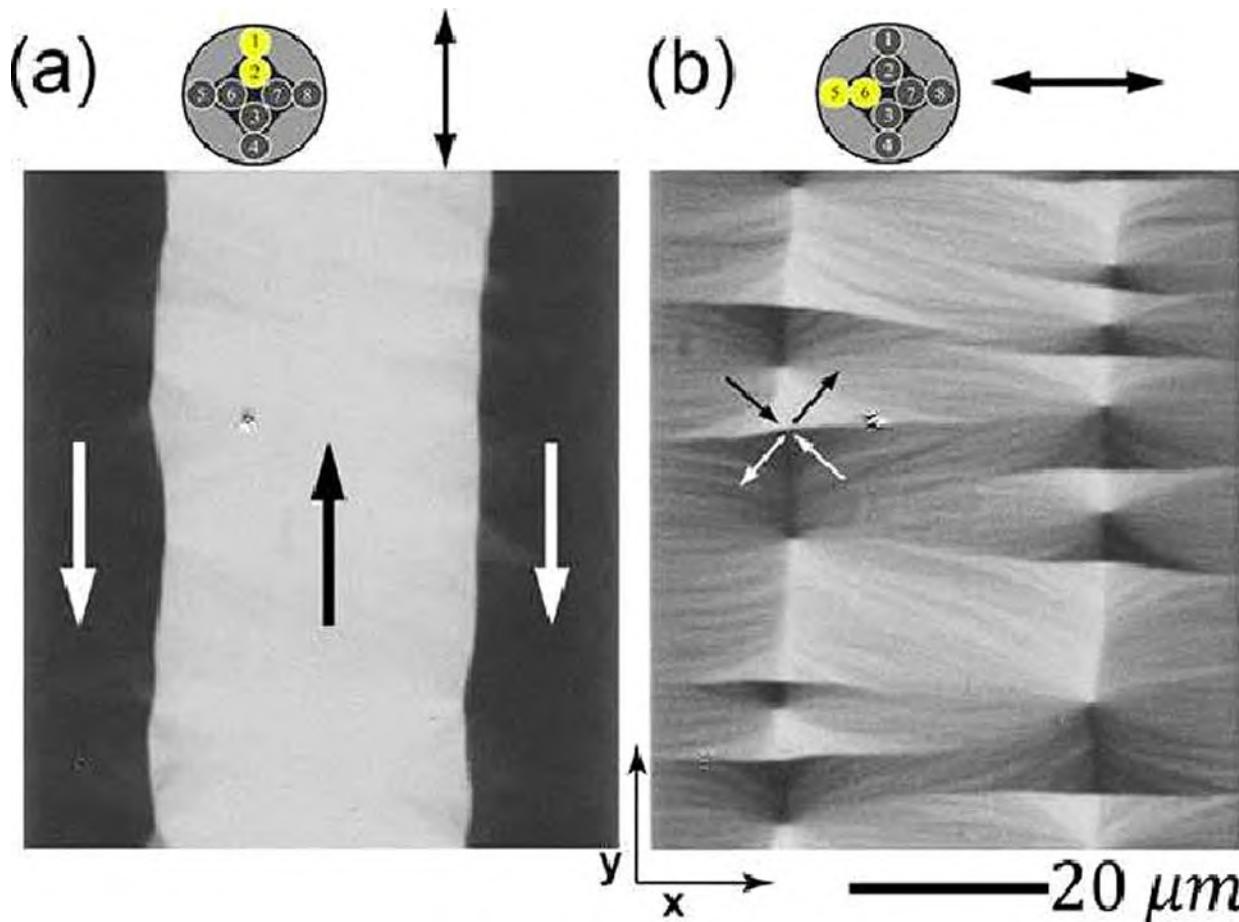
Fe(100) : fourfold in-plane symmetry



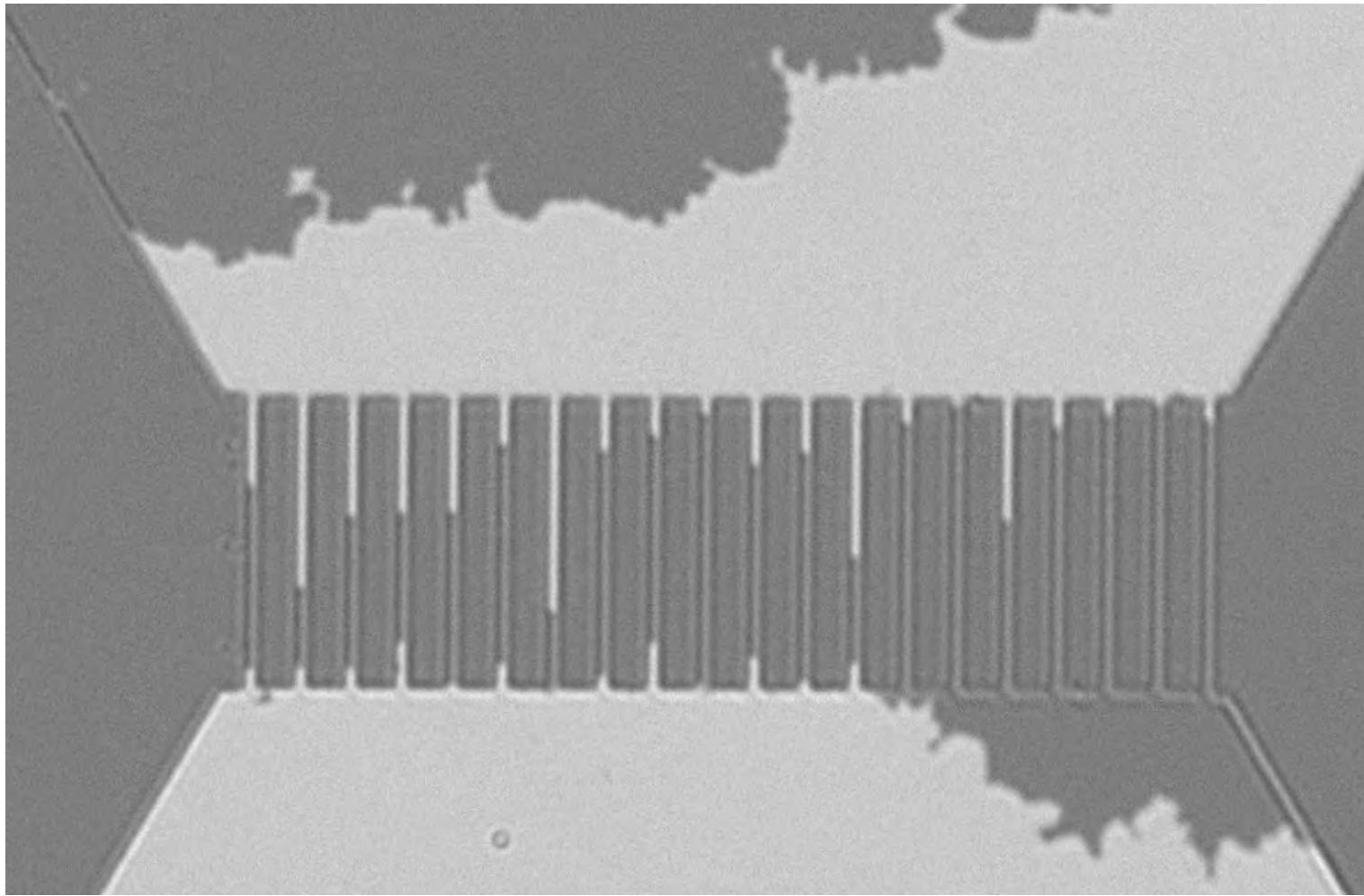
Practical by Jose Luis Coñado

# Kerr microscopy

Magnetic Domains : The Analysis of Magnetic Microstructures  
A. Hubert and R. Schäfer (Springer)

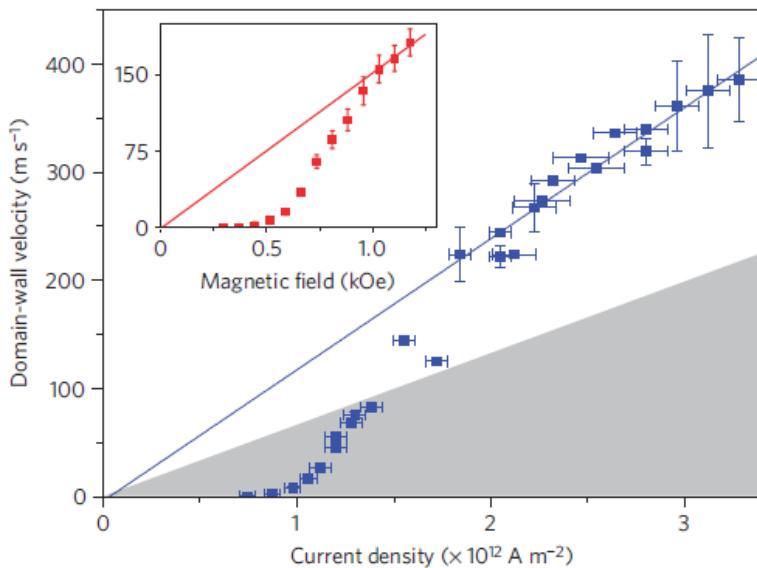
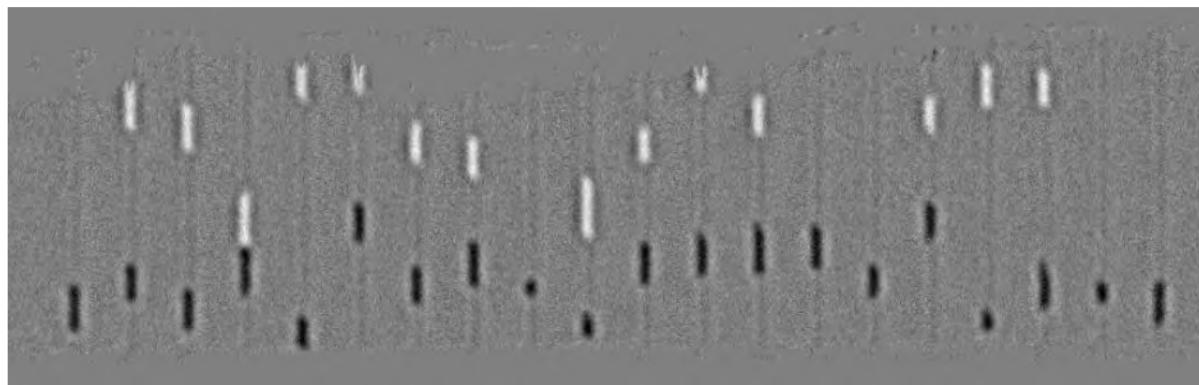


# Kerr microscopy



Magnetic domain walls in Pt/Co/AlOx with perpendicular anisotropy

# Kerr microscopy



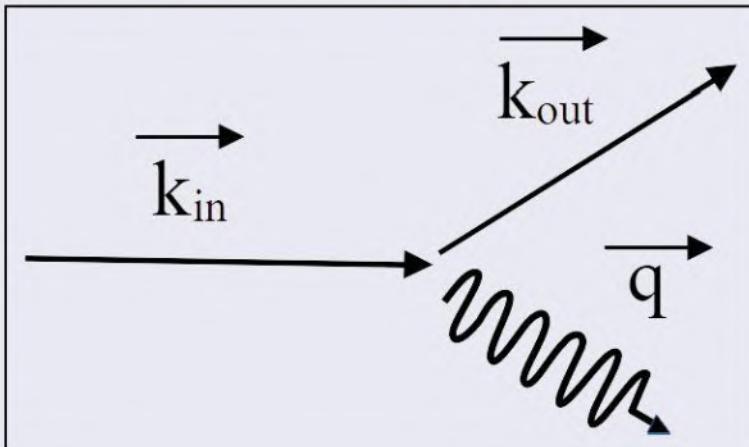
Fast current-induced domain wall motion by spin-orbit torque + Dzyaloshinskii-Moriya interaction

I.M. Miron et al., Nature Mater. **10**, 419 (2011).

# Probing magnetic excitations (magnons) with light

Photon IN – Photon OUT : scattering by Phonons and Magnons

## Inelastic Scattering



The diagram illustrates inelastic scattering. A horizontal arrow labeled  $\vec{k}_{in}$  enters from the left. A second horizontal arrow labeled  $\vec{k}_{out}$  exits to the right. A third arrow labeled  $\vec{q}$  is shown as a wavy line connecting the tip of  $\vec{k}_{in}$  to the tip of  $\vec{k}_{out}$ , representing the momentum transfer.

- Conservation of Energy
- Conservation of Wavevector
- $\vec{k}_{in} = \vec{k}_{out} \pm \vec{q}$

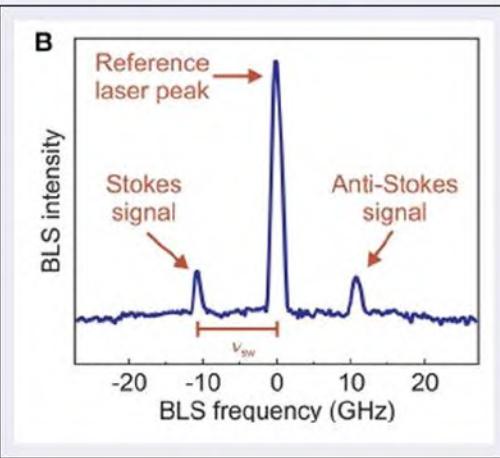
$$k = 4\frac{\pi}{\lambda}$$

Visible light (HeNe – laser) :  
 $\lambda = 532 \text{ nm} \rightarrow k = 23.6 \mu\text{m}^{-1}$   
 $v = 563 \text{ THz}$

©Laurent Ranno

BLS : spectrometry of the scattered light around the elastic line

## Inelastic Scattering



- Elastic mode : Rayleigh scattering
- Emission of excitation : Stokes mode
- Absorption of excitation : Anti-Stokes mode

## E(q) dispersion

- Acoustic Phonons :  $E = \hbar c q$
- Magnons :  $E = D q^2$
- Dispersion  $E(q)$  by rotating the sample
- Need for stable and high resolution spectrometer

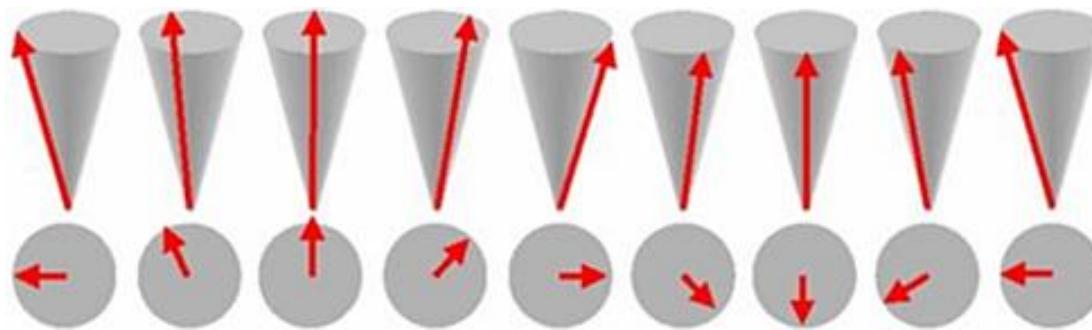
$E(q) \Rightarrow$  Sound velocity,  
Mechanical properties  
(contactless)

$E(q) \Rightarrow A_{ex}, DMI, K_{anis}, M_s,$   
relaxation, excitation distribution

...

©Laurent Ranno

## Spin waves (magnons)



Frequency (energy) of magnons depends on  $M_s$ ,  $H$ ,  $K$ ,  $A_{ex}$ , DMI...

Lecture by Oksana Chubykalo-Fesenko on Friday

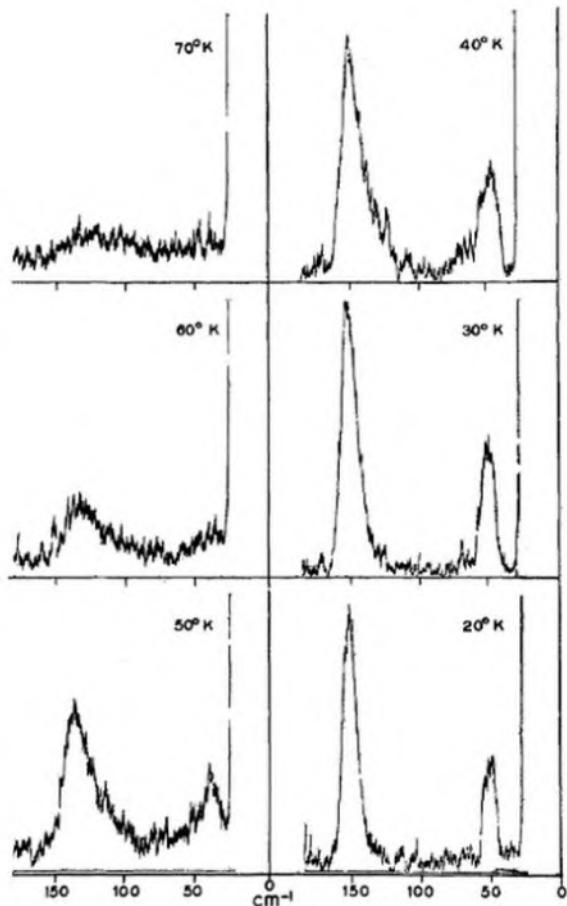


FIG. 1. Recorder traces showing frequency shifts of Stokes scattered light in the (*zy*) experimental geometry for various temperatures in  $\text{FeF}_2$ . The lines at  $\sim 52$  and  $\sim 154 \text{ cm}^{-1}$  are due to photons scattered by one and two magnons, respectively.

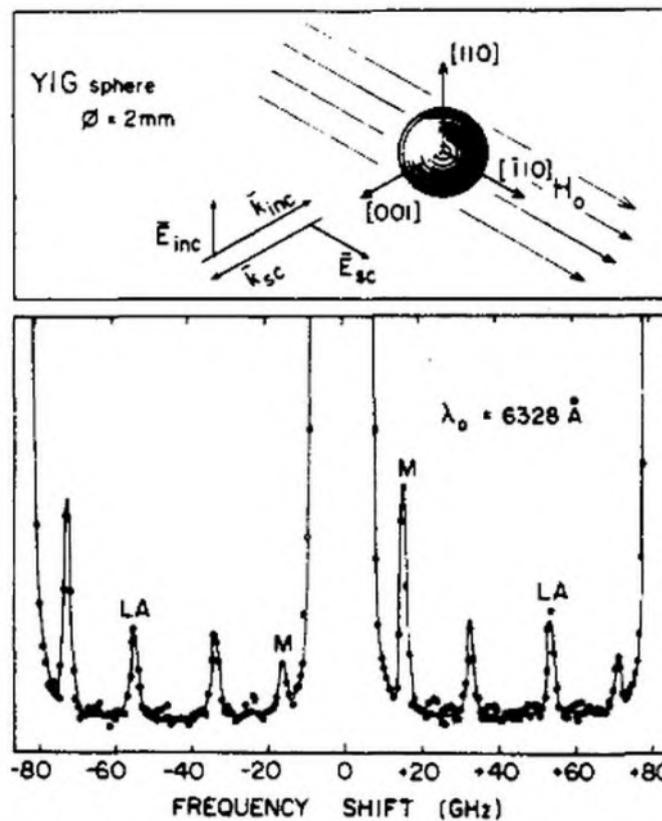
## First BLS measurement on magnons

### $\text{FeF}_2$

- Néel Temp = 78K
- Polarisation analyses
- Scattering mechanism : one-magnon process ( $52 \text{ cm}^{-1}$ )
- Scattering mechanism : 2-magnon process ( $154 \text{ cm}^{-1}$ )

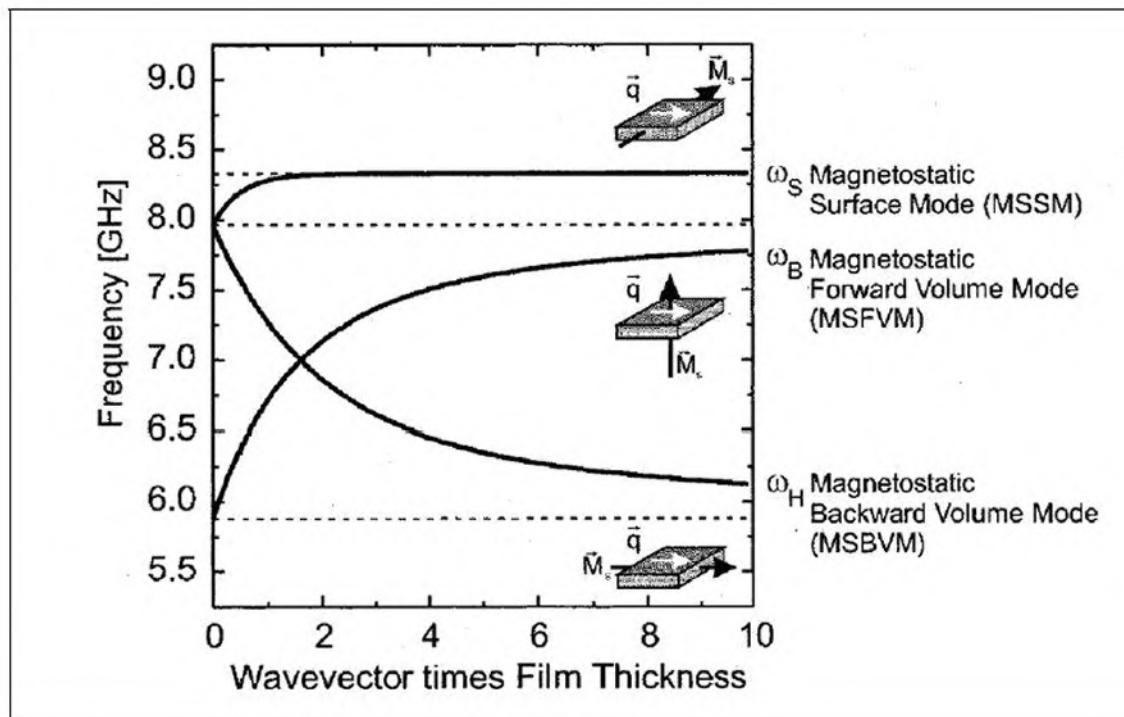
P.A. Fleury et al, Phys. Rev. Lett. 18, 84 (1966)

## Ytrrium Iron Garnet

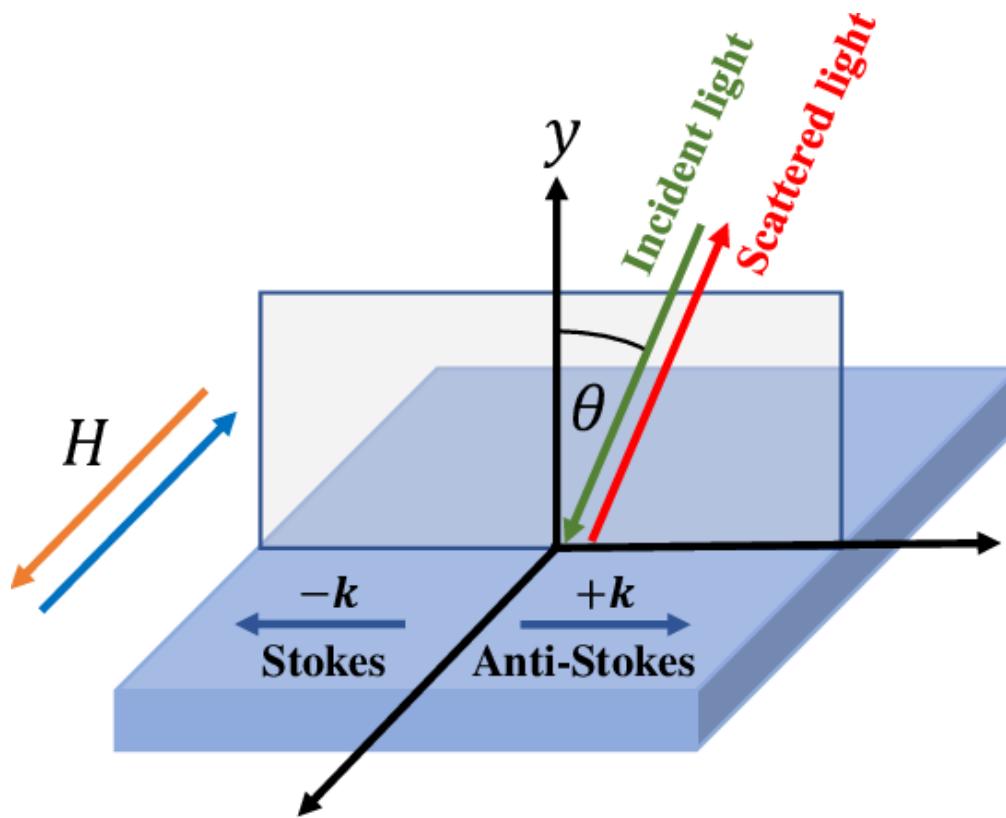


Very narrow scattering peaks  $\rightarrow$  very small damping constant

## Excited spin-waves depend on measuring geometry



Most measurements on thin films in Damon-Eshbach geometry



Three informations : Line energy, line intensity, linewidth

- Scattering amplitude related to magnon population  
Low temperature : Little Anti-Stokes ! (no excitation to absorb)
- Linewidth : damping mechanisms or excitation distribution (+ resolution).

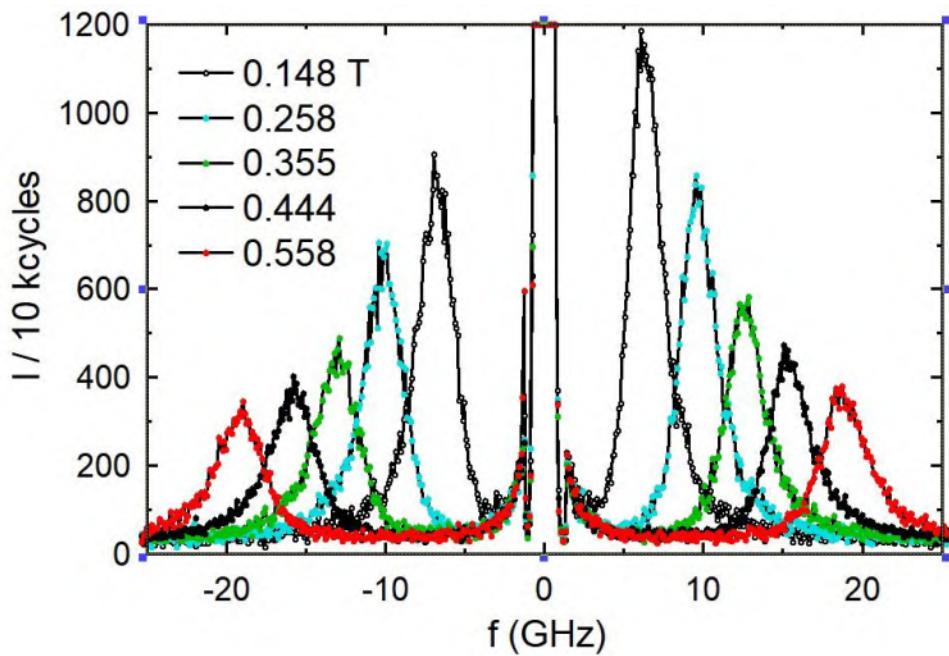
Frequency of Stokes and Anti-Stokes in presence of DMI :

$$f = \frac{\gamma}{2\pi} \sqrt{B_x B_z} \pm \frac{\gamma}{\pi M_s} Dk \equiv f_0 \pm f_D$$

with  $B_x = \mu_0 H + P(kt) \mu_0 M_s + \frac{2A}{M_s} k^2$

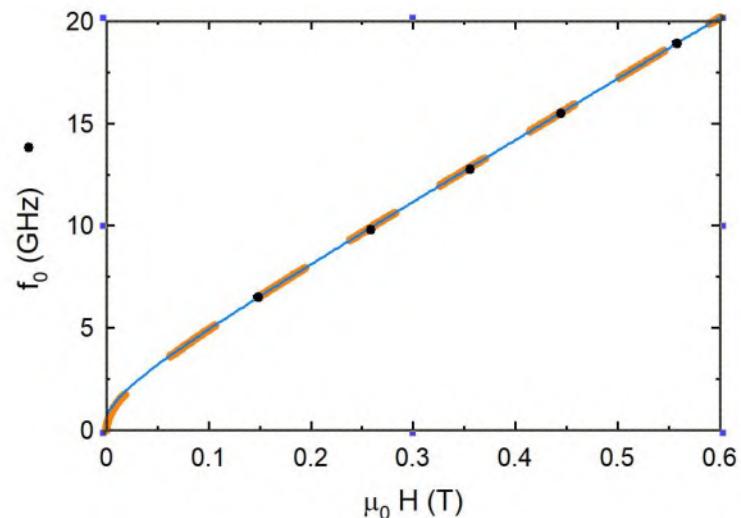
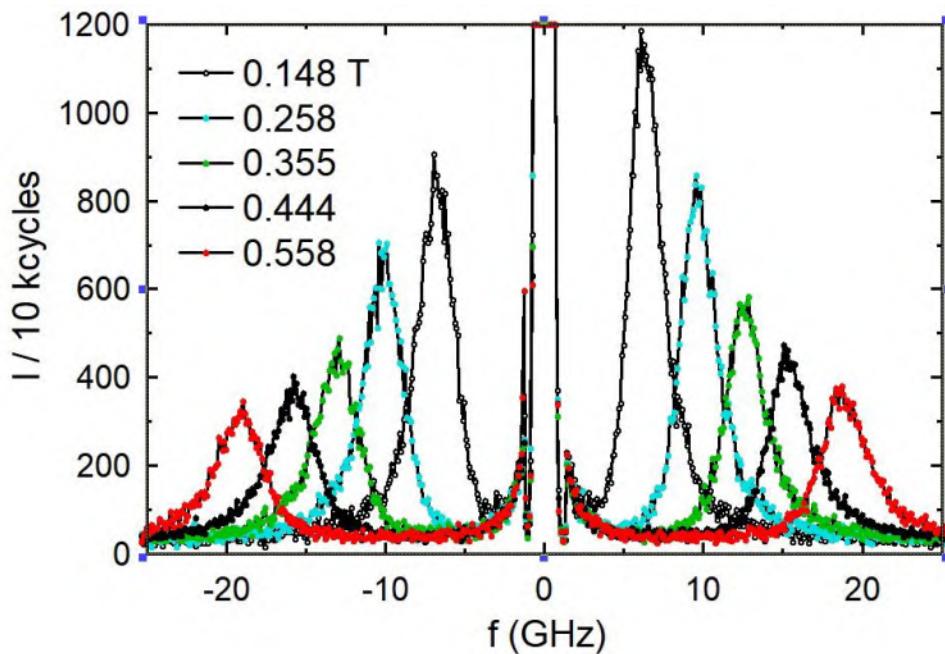
$$B_z = \mu_0 (H - H_K) - P(kt) \mu_0 M_s + \frac{2A}{M_s} k^2$$

BLS spectra for Pt/Co(1.1nm)/AlOx  
as a function of field, 10° incidence



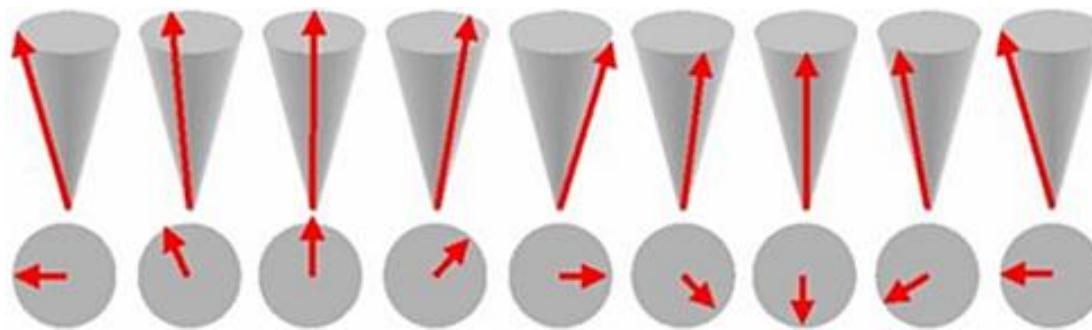
# BLS to determine anisotropy field

BLS spectra for Pt/Co(1.1nm)/AlOx  
as a function of field, 10° incidence



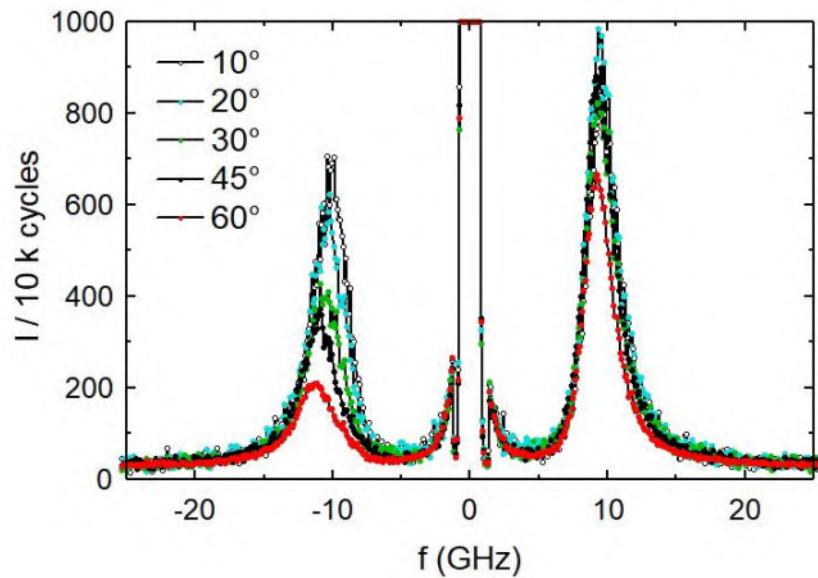
$$\mu_0 H_K = 0.176 \pm 0.009 \text{ T}$$

Measurements by A. Thiaville  
C. Balan et al., Small 19, 2302039 (2023)

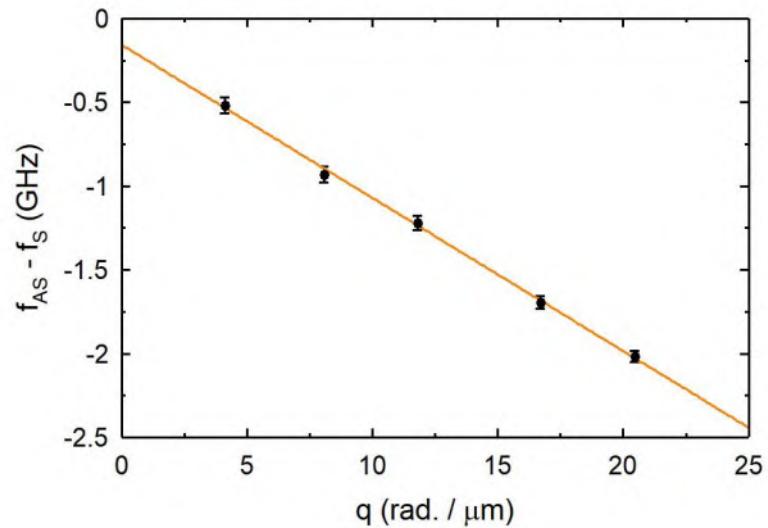
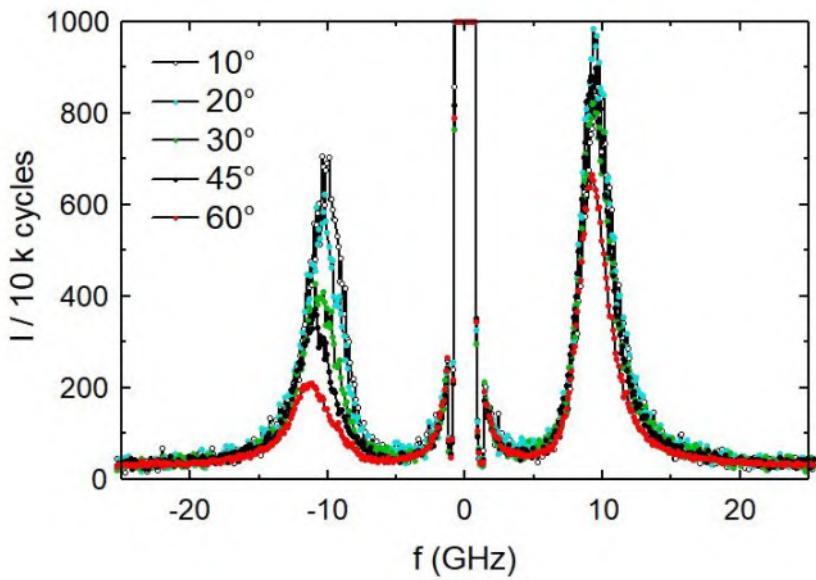
$\rightarrow k$  $\leftarrow k$ 

In the presence of DMI, magnons with  $k \leftarrow$  and  $k \rightarrow$   
do not have the same energy

## BLS spectra for Pt/Co(1.1nm)/AlOx



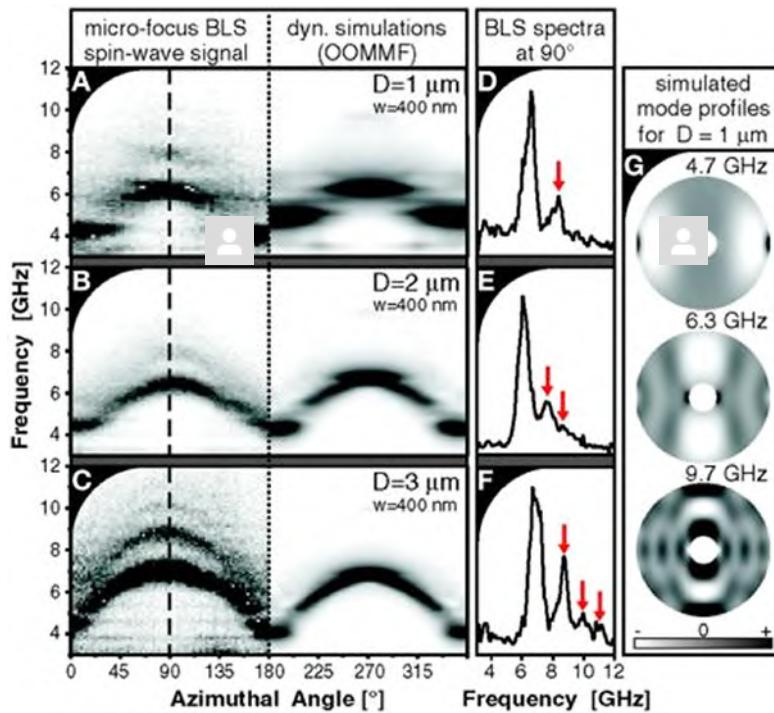
## BLS spectra for Pt/Co(1.1nm)/AlOx



$$D = -0.827 \pm 0.017 \text{ mJ/m}^2$$

Measurements by A. Thiaville  
C. Balan et al., Small 19, 2302039 (2023)

Micro-BLS can be used to study spin wave spectrum in small structures  
Loss of  $k$  (divergent beam)

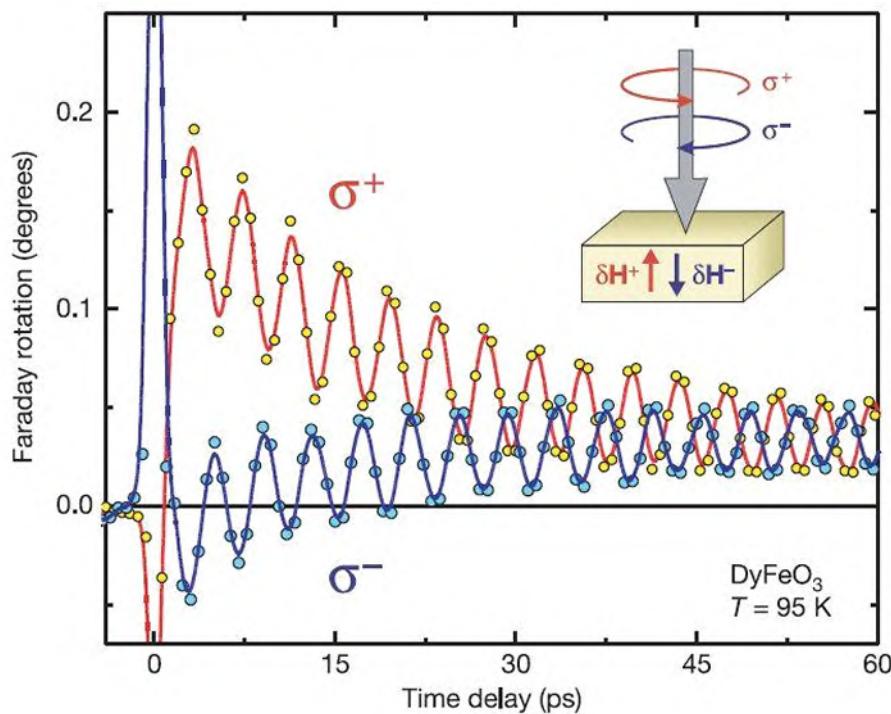


**Micro-focused Brillouin light scattering:  
imaging spin waves at the nanoscale**  
T. Sebastian, K. Schultheiss, B. Obry, B.  
Hillebrands, H. Schultheiss  
Frontiers in Physics 3 (2015)

## 2) Influence of light on magnetic properties

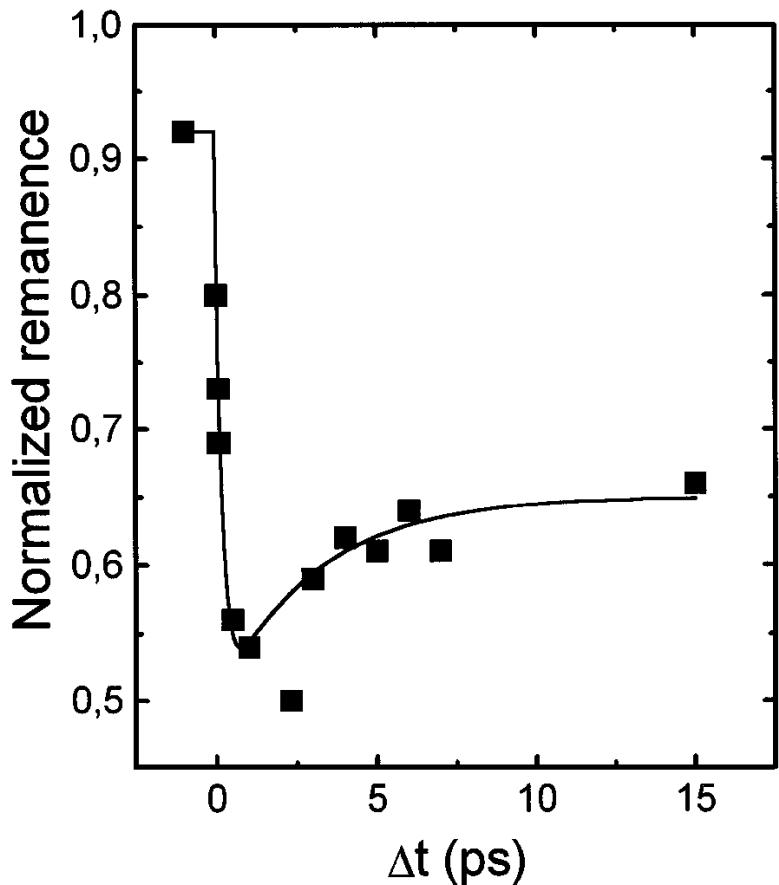
# Influence of light on magnetic properties

A.V. Kimel et al., Nature 435, 655 (2005)



Magnetic excitations in  $\text{DyFeO}_3$  probed by the magneto-optical Faraday effect  
 Left- and right-circularly polarized photons induced opposite (small) magnetic moments  
 → Inverse Faraday effect

## Ultrafast demagnetization induced by fs laser pulses



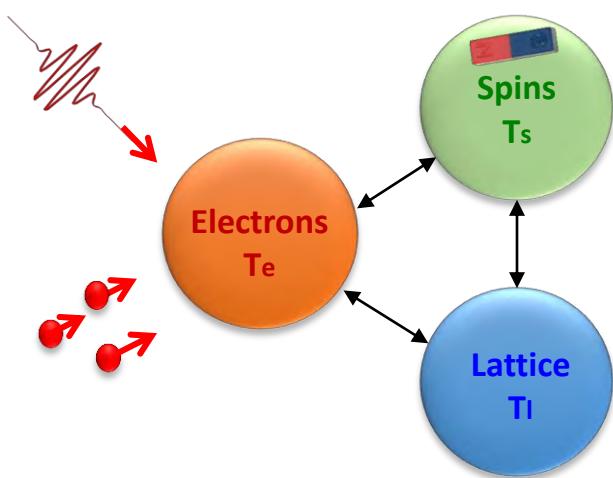
20nm Ni film  
Excitation by 60fs laser pulse  
620nm, 7 mJ/cm<sup>2</sup>

Magnetic response measured using  
longitudinal Kerr in pump-probe mode

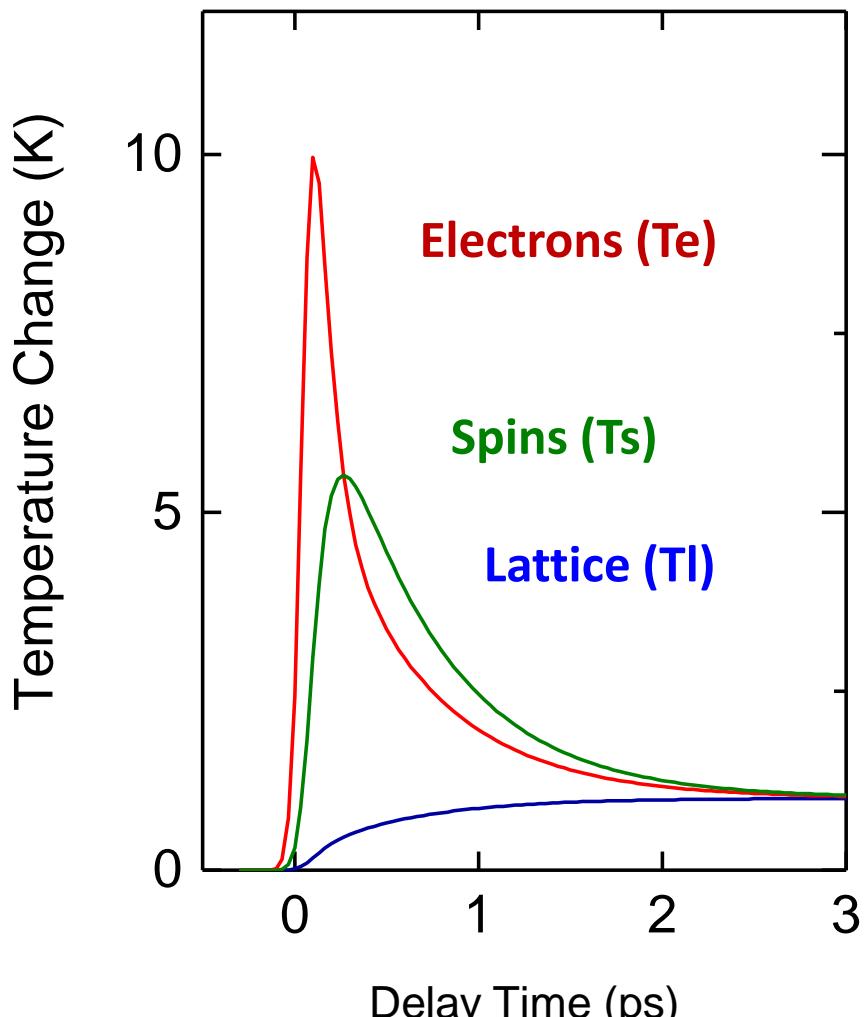
Demagnetization in few hundred femto-  
seconds

E. Beaurepaire, J.-C. Merle, A. Daunois and  
J.-Y. Bigot, Phys. Rev. Lett. 76, 4250 (1997)

# Ultrafast demagnetization by laser pulses



- Energy Transfer
- Angular Momentum Transfer
- Charge / Spin Current Transfer

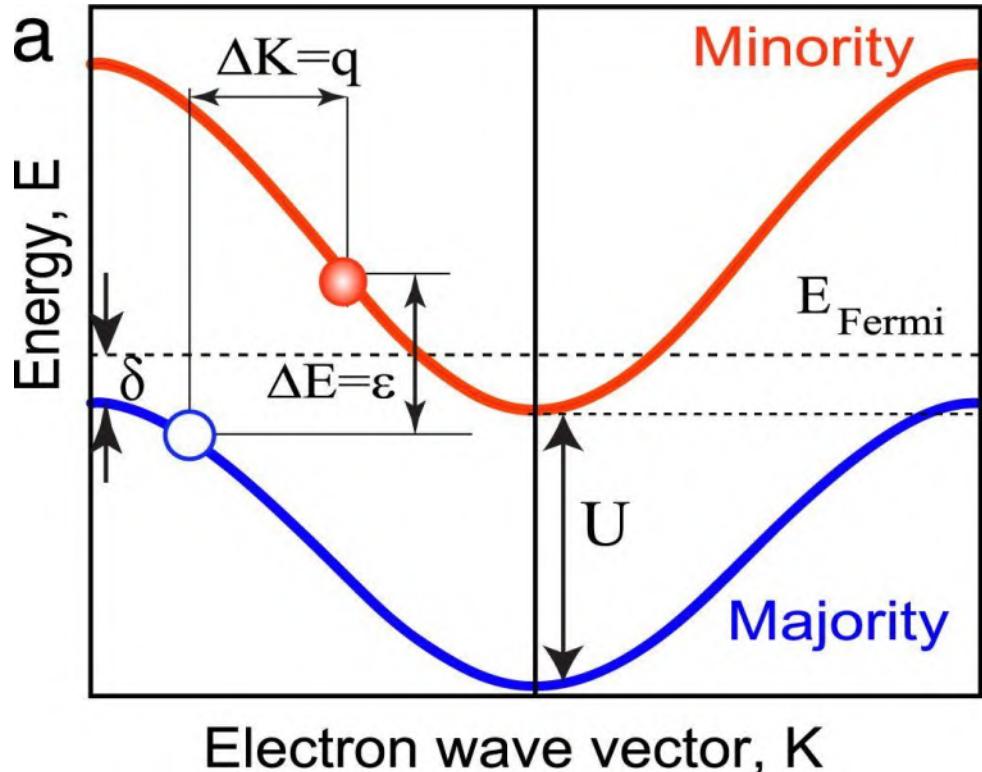


©S. Mangin

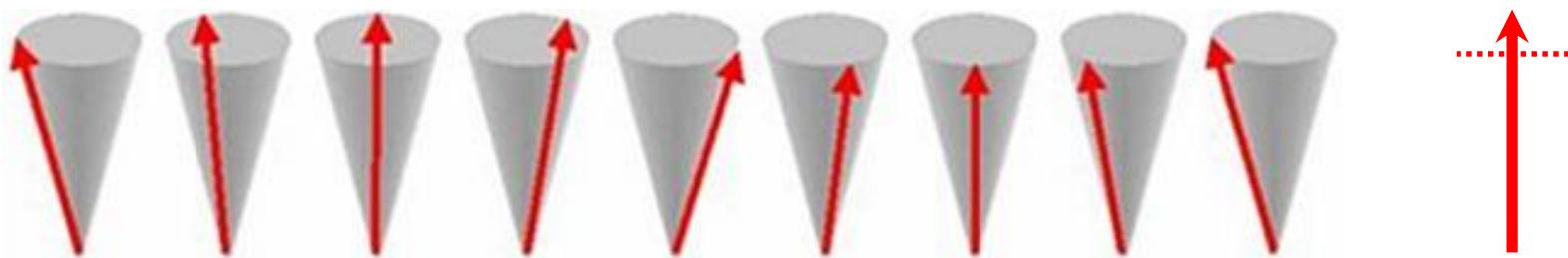
- Reduction of the length of the magnetic moment carried by each atom in the matter (creation of Stoner excitations resulting from a decrease of the exchange splitting).
- a tilt of these vectors resulting in a lower magnetization in average (creation of magnons).
- the propagation of spin-polarized hot electrons.

**Review :**

P. Scheid, Q. Remy, S. Lebègue, G. Malinowski, S. Mangin  
*J. Magn. Magn. Mater.* 560, 169596 (2022).

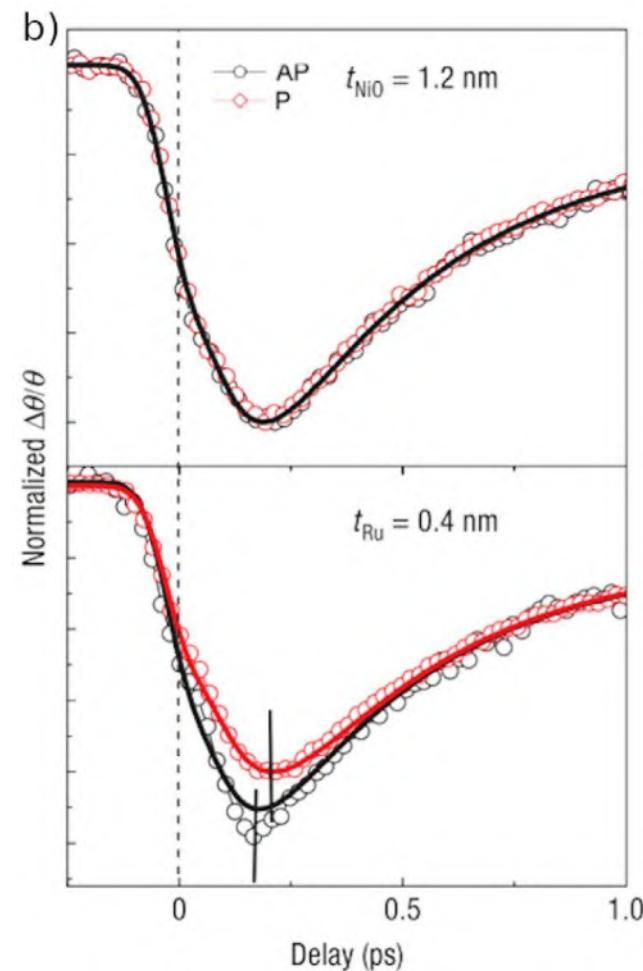
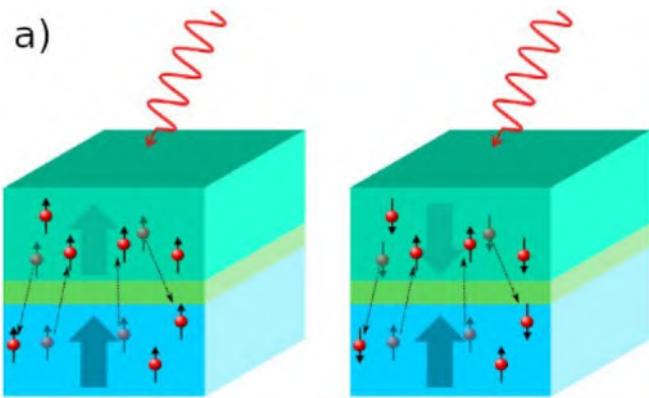


Stoner excitations :  
electron–hole pair  
excitations with electrons  
and holes in the bands of  
opposite spins



Spin-waves : magnetic « disorder » → magnetization decrease

# Ultrafast demagnetization by laser pulses



Comparison (Co/Pt)/Ru/(Co/Pt) and (Co/Pt)/NiO/(Co/Pt) trilayers

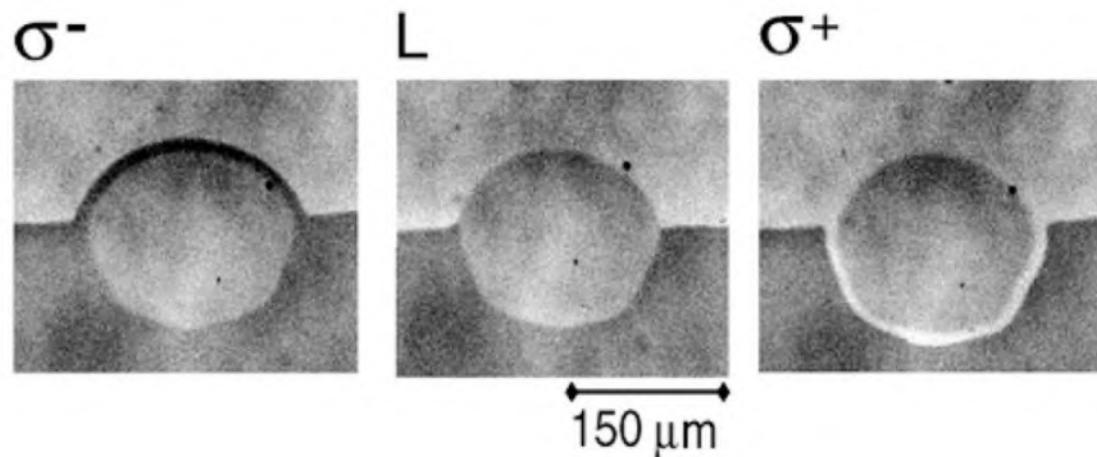
G. Malinowski et al., Nat. Phys. 4, 855 (2008).

« Superdiffusive » spin currents : hot spin-polarized electrons leaving the irradiated zone

## Where does the angular moment go ?

- Spin → orbital magnetic moment ?
- Lattice (phonons) ?
- Superdiffusive spin currents ?
- THz emission ?

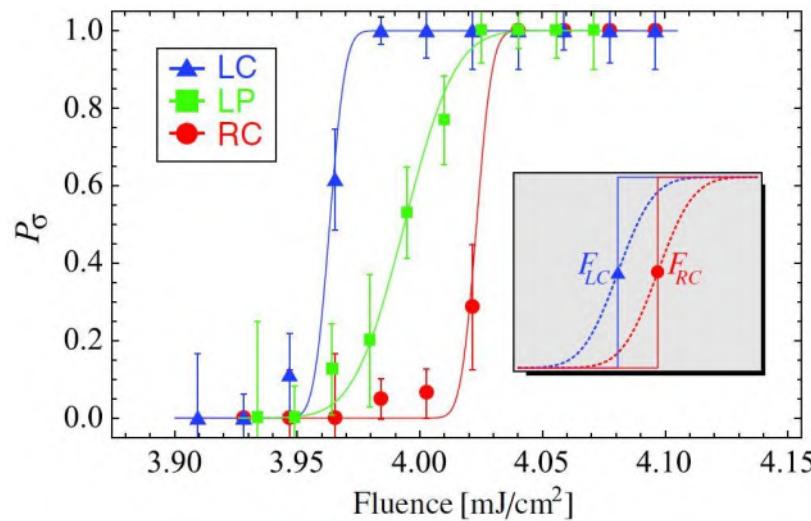
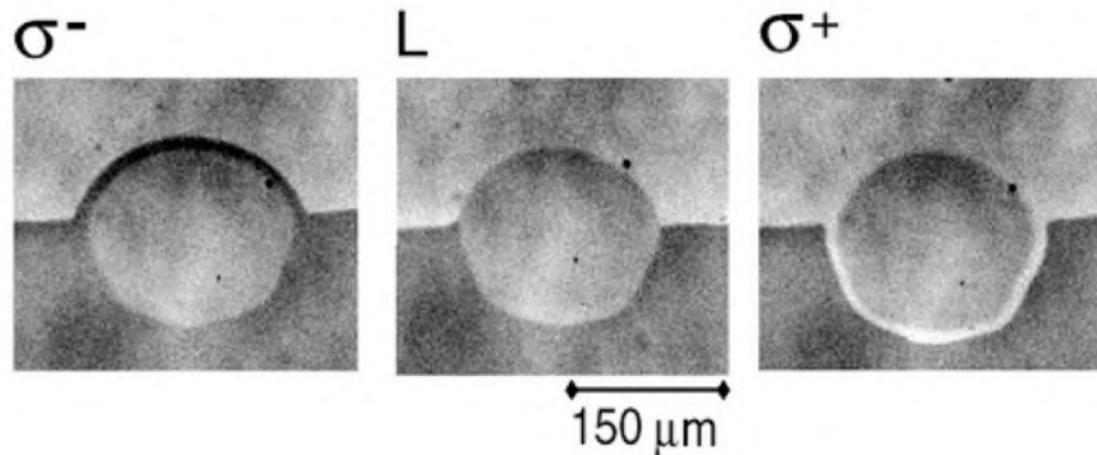
# All-optical magnetic switching



C.D. Stanciu, F. Hansteen, A.V. Kimel, A. Kirilyuk, A. Tsukamoto, A. Itoh, T. Rasing,  
Phys. Rev. Lett. 99 (4) (2007) 47601

Sample :  $\text{Gd}_{22}\text{Fe}_{74.6}\text{Co}_{3.4}$

# All-optical magnetic switching

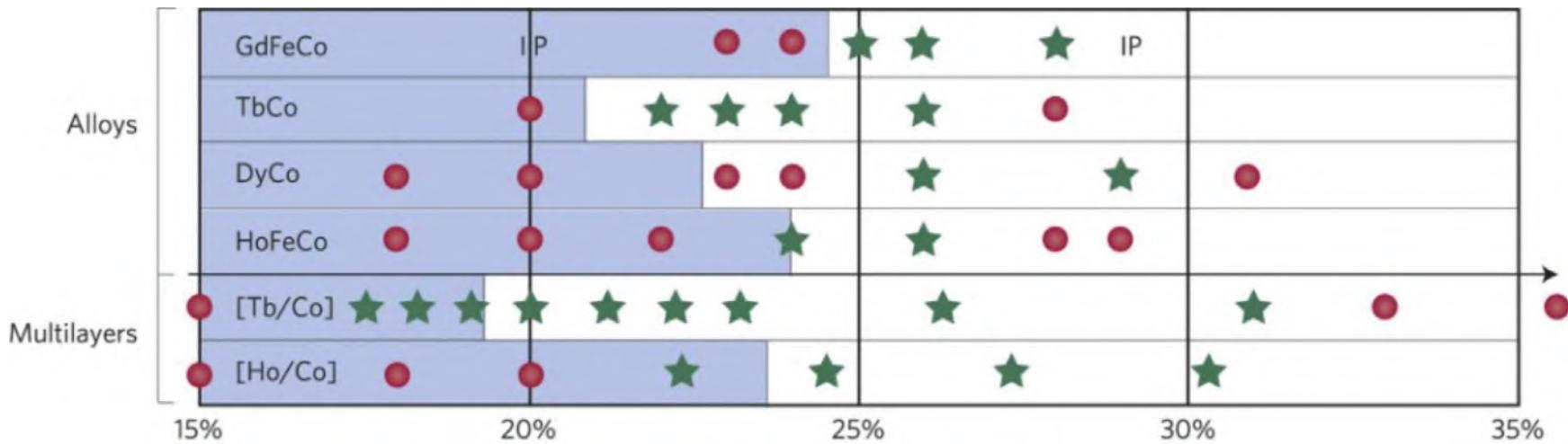


Magnetic circular dichroism in photon absorption

A.R. Khorsand, M. Savoini, A. Kirilyuk,  
 A.V. Kimel, A. Tsukamoto, A. Itoh, T.  
 Rasing, Phys. Rev. Lett. 108 (12)  
 (2012)

# All-optical magnetic switching

## AO - HDS



AO – HDS observed also in Rare-Earth-free synthetic antiferromagnets

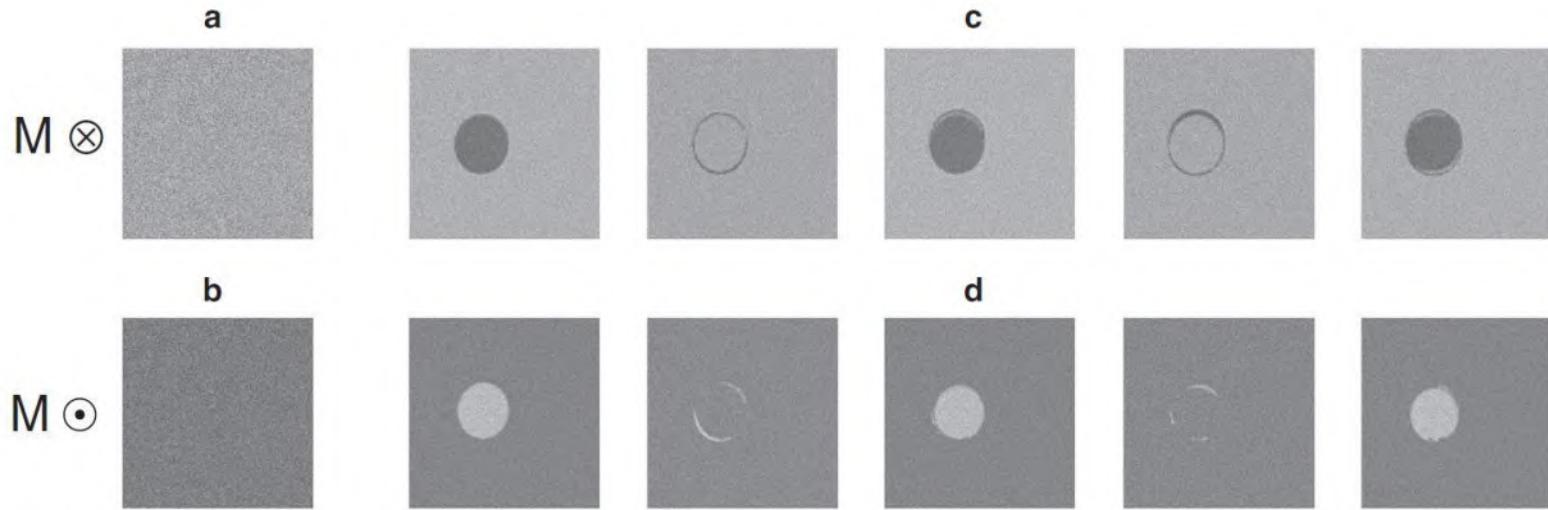
### Review :

P. Scheid, Q. Remy, S. Lebègue, G. Malinowski, S. Mangin  
 J. Magn. Magn. Mater. 560, 169596 (2022).

AO – HDS :

- 1) Magnetic circular dichroism : different absorption of left- and right-circularly polarized → different thermal effects
- 2) Inverse Faraday effect (effective magnetic field induced by the light, induced magnetization ?)
- 3) Demagnetization induced by optical induced transitions that, in the presence of spin-orbit coupling, do not conserve the magnetization. Magnitude depends on both the helicity of the light and the direction of the magnetization.
- 4) Direct spin-light coupling (in presence of SOC).

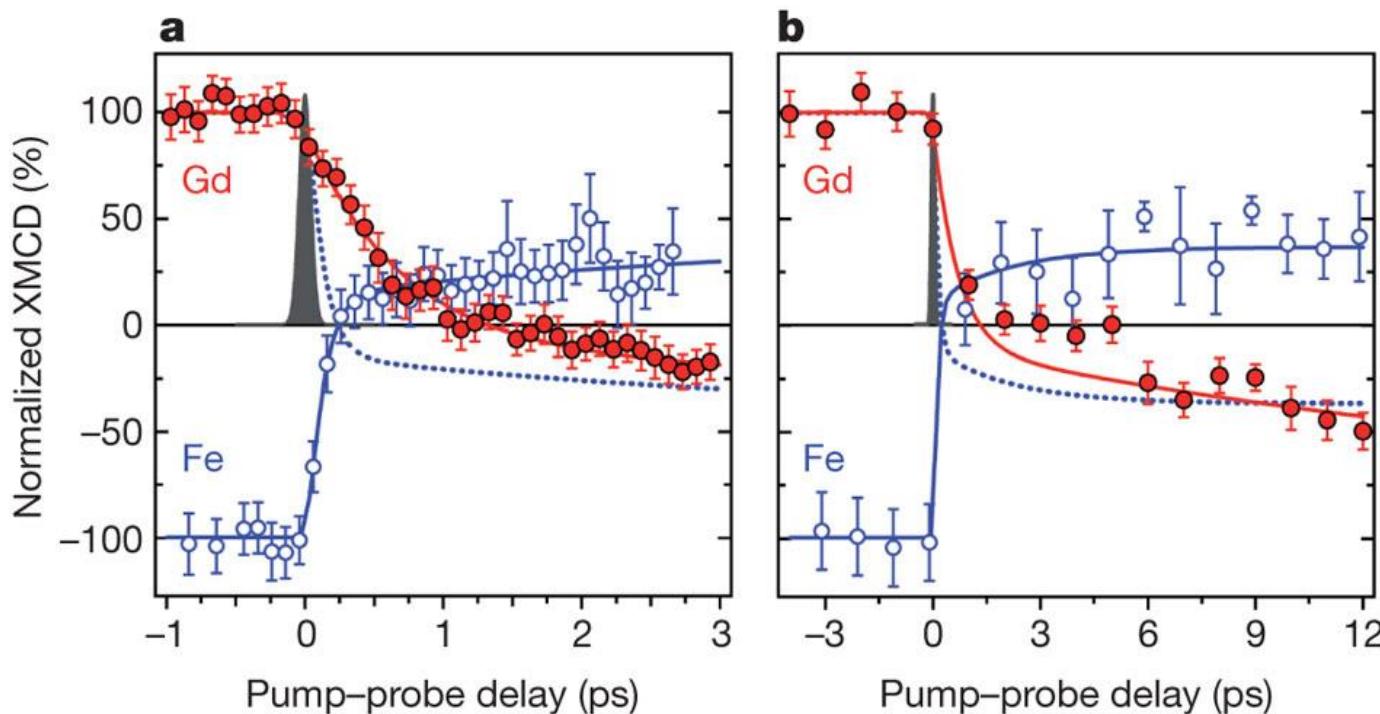
## All-optical helicity-independent switching (AO HIS)



T.A. Ostler et al., Ultrafast heating as a sufficient stimulus for magnetization reversal in a ferrimagnet, Nature Commun. 3 (2012) 666



## All-optical helicity-independent switching (AO HIS)



I. Radu, K. Vahaplar, C. Stamm, T. Kachel, N. Pontius, H.A. Dürr, T.A. Ostler,  
J. Barker, R.F.L. Evans, R.W. Chantrell, A. Tsukamoto, A. Itoh, A. Kirilyuk,  
T. Rasing, A.V. Kimel, Nature 472, 205 (2011).

- X-ray photons can be used to study element selective magnetic properties (spin, orbital moment) with good spatial resolution (<20nm) and temporal resolution BUT needs a synchrotron
- Optical photons can be used to study magnetic properties (hysteresis, reversal modes, imaging) with  $<\mu\text{m}$  spatial resolution (Kerr, Faraday)
- Brillouin Light Scattering spectroscopy is a versatile tool to study (thin film) properties like A, K, DMI
- High-intensity femtosecond laser pulses can strongly influence magnetic properties of thin films (demagnetization, all-optical switching)