

## Origin of Dzialoshinsky-Moriya interactions

$$H_{\text{DM}} = \mathbf{D}_{12} \cdot (\mathbf{S}_1 \times \mathbf{S}_2)$$

-Antisymmetric interaction:  $\neq 0$  only if there is no center of symmetry between sites 1 and 2

$$-\vec{D}_{12} = -\vec{D}_{21}$$

- Produces a canting of the spins



$$J = 0$$



$$J_{AF} \neq 0$$

## Moriya's rules for $\vec{D}_{12}$



1



M



2

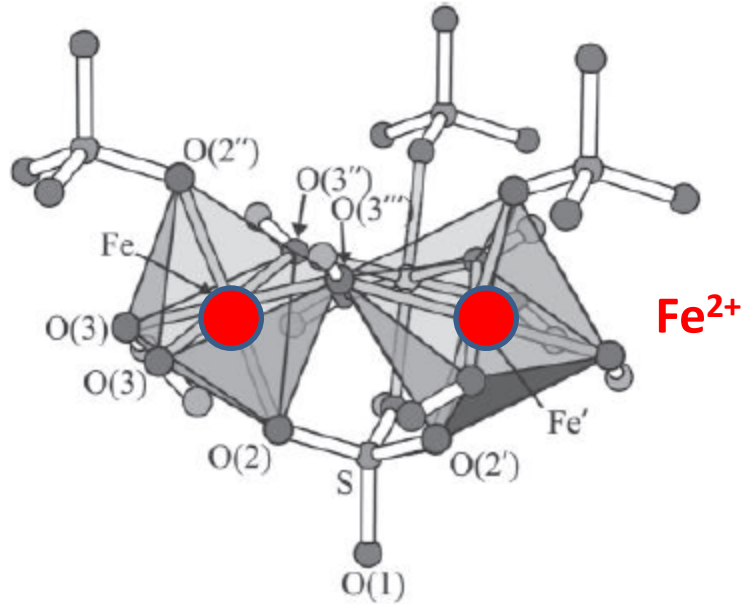
**M** is the middle point between 1 and 2

-If M is an inversion center:  $D=0$

-If the plane perpendicular to 1-2 (containing M) is a mirror plane, D is in this plane

-If 1 plane containing 1 and 2 is a mirror plane, D is perpendicular to this plane

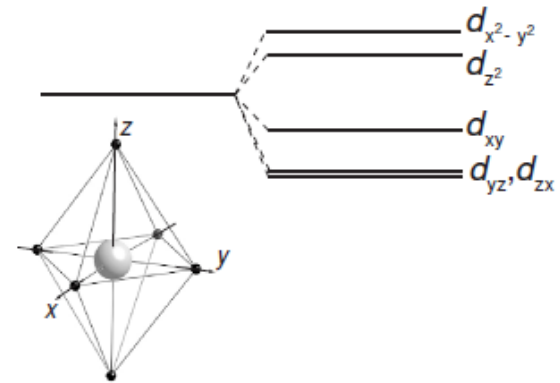
(+ other rules related to the existence of rotation axis)



Example: Fe jarosite

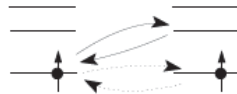
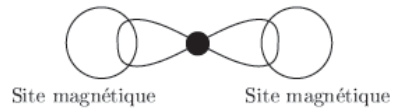
2<sup>nd</sup> rule can be applied

Crystal field splitting around each Fe ion:

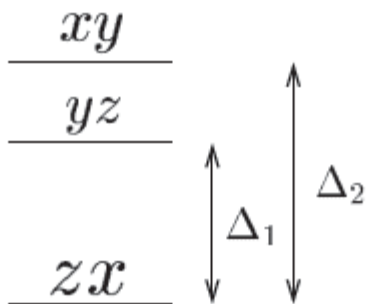


But x, y and z axis are different for each Fe site!

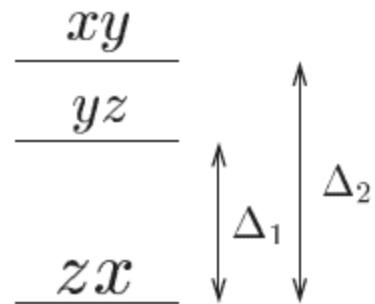
## Superexchange:



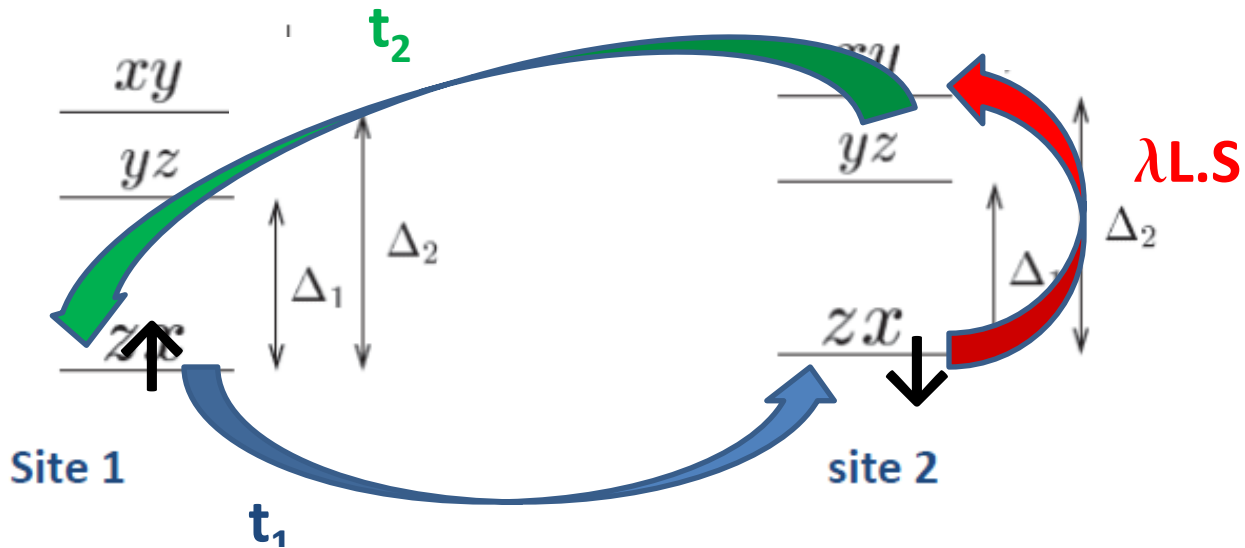
## D.M. interactions: involve spin-orbit interactions



Site 1



site 2



$$d_{xy} = -(i/\sqrt{2})(\psi_2 - \psi_{-2})$$

$$d_{yz} = (i/\sqrt{2})(\psi_1 + \psi_{-1})$$

**Matrix elements of  $\lambda L.S \neq 0$**

$$d_{zx} = -(1/\sqrt{2})(\psi_1 - \psi_{-1})$$

Contribution to the energy: perturbation in  $t$  and  $\lambda \Rightarrow \lambda \frac{t_1 t_2}{\Delta U}$

Superexchange :  $\frac{t_1^2}{U}$

$$D \sim \lambda \frac{t_1 t_2}{\Delta U} \quad J \sim \frac{t_1^2}{U}$$

To next order in  $\lambda$ , anisotropic exchange:  $\sum_{\alpha\beta} \Gamma_{\alpha\beta} (\mathbf{S}_1^\alpha \mathbf{S}_2^\beta + \mathbf{S}_2^\alpha \mathbf{S}_1^\beta)$

$$\Gamma \sim \lambda^2 \frac{t_1 t_2}{\Delta^2 U}$$

$$D/J \approx \lambda/\Delta \quad \Gamma/J \approx (\lambda/\Delta)^2$$

Small interactions, play an important role in multiferroics

Observed in many oxides of TM ions: Fe<sub>2</sub>O<sub>3</sub>, CoCO<sub>3</sub>, ....

Weak ferromagnetism, canting angle  $\approx D/J$